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Dark Energy and CMB
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28.1 Cosmology and New Physics

Maps of the Universe when it was 400,000 years old from observations of the cosmic microwave background and over the last ten billion years from galaxy surveys point to a compelling cosmological model. This model requires a very early epoch of accelerated expansion, inflation, during which the seeds of structure were planted via quantum mechanical fluctuations. These seeds began to grow via gravitational instability during the epoch in which dark matter dominated the energy density of the universe, transforming small perturbations laid down during inflation into nonlinear structures such as million light-year sized clusters, galaxies, stars, planets, and people. Over the past few billion years, we have entered a new phase, during which the expansion of the Universe is accelerating presumably driven by yet another substance, dark energy.

Cosmologists have historically turned to fundamental physics to understand the early Universe, successfully explaining phenomena as diverse as the formation of the light elements, the process of electron-positron annihilation, and the production of cosmic neutrinos. However, the Standard Model of particle physics has no obvious candidates for inflation, dark matter, and dark energy. The amplitude of the perturbations suggest that the natural scale for inflation is at ultra-high energies\(^1\), so understanding the physics driving inflation could lead to information about the UV completions of our current theories. There are arguments that naturally link the dominant dark matter component to new physics hovering above the electroweak scale, and the powerful suite of experiments aiming to find this component(s) were the focus of several separate groups in the Snowmass process. Apart from the dominant component, neutrino oscillation experiments already inform us that neutrinos constitute a non-negligible fraction of the dark matter, and an important message is that experiments usually associated with dark energy and inflation are ideally suited to pin down the sum of the masses of the neutrinos and the cosmic existence of any additional (sterile) species. The situation with dark energy is more complex. A cosmological constant (\(\Lambda\)) has effective pressure equal to minus its energy density (equation of state \(w = -1\)) consistent with preliminary measurements, but for example in supersymmetric theories the most natural scale for \(\Lambda\) is at least as large as 100 GeV. A cosmological constant with this value would produce a universe accelerating so rapidly that the tips of our noses would be expanding away from our faces at a tenth the speed of light. If \(\Lambda\) is responsible for the current epoch of acceleration, its value is many orders of magnitude smaller than this but curiously just large enough that it began dominating the energy density of the universe only recently. So the mechanism driving the current accelerated expansion of the Universe remains a profound mystery.

The quest to understand dark energy, dark matter, and inflation then is driven by a fundamental tension between the extraordinary success of the model that explains our Universe and the failure of the Standard

\(^1\)Roughly, the measured amplitude of the density perturbations \(\delta \rho / \rho \approx 10^{-5} \sim (E_{\text{inf}} / m_{\text{Planck}})^2 / \sqrt{\epsilon}\), where \(\epsilon \approx 0.01\) is a small parameter in slow roll inflation.
Model of particle physics to provide suitable candidates for the dark sector that is so essential to our current view of the Universe. Experiments on the cosmic frontier have demonstrated that the Standard Model is incomplete; the next generation of experiments can provide the clues that will help identify the new physics required. Complementing the efforts on the Intensity and Energy Frontier, physicists working on the Cosmic Frontier are poised to unravel the mysteries of the cosmos.

### 28.2 Dark Energy

Physicists have proposed a number of different mechanisms that could be responsible for the accelerated expansion of the Universe. None is compelling, but some of them have been predictive enough to fail, while others have led to a deeper understanding of field theory and gravity. Apart from the cosmological constant solution itself, all are predicated on the assumption that there is some (unknown) mechanism that sets \( \Lambda \) to zero.

One possibility is that there is a previously undiscovered substance that contributes to the energy density of the Universe in such a way that the expansion accelerates. In the context of general relativity, acceleration (the positive second derivative of the scale factor \( a \)) is governed by Einstein’s Equations, which reduce to

\[
\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} [\rho + 3P] \tag{28.1}
\]

where \( G \) is Newton’s constant, \( \rho \) the energy density, and \( P \) the pressure. The substance that drives acceleration, therefore, must have negative pressure, or equation of state \( w = P/\rho < -1/3 \). A nearly homogeneous scalar field whose potential energy dominates over its kinetic energy satisfies this requirement, so many *quintessence* models emerged with potentials designed to fit the data. In almost all viable models, however, the mass of the field must be less than the Hubble scale today, of order \( 10^{-33} \) eV. Embedding such a field in some extension of the Standard Model therefore is challenging as one would expect the scalar mass to get loop corrections many orders of magnitude larger than this. This mass stability problem is an indication of just how hard the problem is: the mass that is protected is some 44 orders of magnitude smaller than the Higgs mass that lies at the center of the electroweak hierarchy problem. It seems clear that the new physics is an infrared phenomenon as opposed to all prior hints of new physics, which have entered from the ultraviolet.

With no real guidance from theory, quintessence models are nevertheless appealing because they open up the parameter space: most models have \( w \neq -1 \), and many have evolving equations of state so that \( dw/da \) is also non-zero. This class of models therefore offers a clear, and arguably more appealing, alternative to the cosmological constant that would be favored if surveys find deviation from \( w = -1 \). Although the quintessence field does not clump on scales smaller than its (very large) Compton wavelength, on the largest scales dark energy should clump, yet another difference from the cosmological constant model. Finally, some models allow for episodic dark energy domination, so that the present accelerating era is not particularly special. Indeed, an early epoch of inflation is one such epoch, but there may have been others, for example during phase transitions. Measuring the effects of dark energy in a series of redshift bins is therefore necessary to distinguish among the many possibilities. This is an area where the cosmic microwave background, whose lensing maps are sensitive to structure from \( z = 10^3 \) until today, can be profitably combined with galaxy surveys, whose lensing maps probe structure at a sequence of lower redshifts.

Quintessence models explicitly introduce an extra degree of freedom in the form of a new scalar field. Early attempts to modify gravity (the “left-hand side of Einstein’s Equations”) implicitly introduced an extra degree of freedom; the scalar field could provide the mechanism to explain the observed acceleration of the Universe.

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2Redshift \( z \equiv a^{-1} - 1 \); high redshift corresponds to early cosmic time.
degree of freedom \[2\], but in a slightly different way. Technically, these “scalar-tensor” models differ from quintessence in that the coefficient of the Ricci scalar $R$ in the Jordan frame in the action depends on the new field. Different scalar-tensor models can therefore produce the variety of predictions found in quintessence models, but they extend the possibilities in a new direction: the non-canonical coupling to the Ricci scalar propagates to the equations that govern the evolution of perturbations. This leads to the general conclusion that modified gravity models typically predict that structure grows at a different rate (e.g., \[3\]) than in models based on general relativity. Differentiating between modified gravity and General Relativity-based dark energy therefore will require both measurements of cosmological distances to pin down the background expansion and then measurements of the growth of structure to distinguish between them.

Although there have been many proposals for how to modify gravity, the most interesting development recently traces back to an idea first proposed by Fierz and Pauli \[4\], that the graviton has non-zero mass. Qualitatively a massive graviton seems like an appropriate way to decrease the strength of the gravitational force on very large scales and hence to explain the acceleration of the Universe. In practice, the theory runs into two problems, both related to the fact that a massive spin-2 particle carries degrees of freedom beyond those of the massless graviton. These extra degrees of freedom typically lead to modifications to general relativity in the Solar System, modifications that are excluded by the tight limits on post-Newtonian parameters. The second challenge for massive graviton models is to avoid Boulware-Deser ghosts \[5\], the instability of one of the extra degrees of freedom.

To satisfy the Solar System constraints, the extra degrees of freedom need to be **screened**, i.e. heavily suppressed by limiting their range of interaction or effective coupling to matter in environments like the Solar System. Any successful modified gravity model needs a screening mechanism. One possibility is Vainshtein screening \[6\], which arises due to non-linearities in the Fierz-Pauli potential. In general, these non-linear terms do not avoid the ghost problem, but recently \[7\] the set of terms in the potential that are safe from ghosts have been identified. These **Galileon** models offer a potentially attractive way of addressing the acceleration of the Universe within a consistent framework. Beyond this theoretical breakthrough, the Vainshtein screening intrinsic to this model (and other proposed mechanisms) open up yet another axis of tests: how and where do modified gravity theories transition to normal Newtonian gravity in the Solar System? The full suite of ways to test modified gravity models is still under development.

The modern view of the formation of structure in the Universe has been confronted with a growing array of precise tests, including the anisotropy spectrum (both temperature and polarization) of the cosmic microwave background; light curves of distant Supernovae; abundances of galaxy clusters; clustering of galaxies, quasars, and Lyman alpha systems; gravitational lensing; and cross-correlations between different pairs of these observations. The basic framework, with inflation, dark matter, and dark energy at its core, has been confirmed repeatedly over the past decade. We need to keep pushing: either the basic picture will break or the agreement will become even more remarkable. Thinking in parameter space, the next decade will enable us to reduce the uncertainty in the equation of state by a factor ten. Historically in physics, precision measurements of this sort have been pursued for these ends: will the simple theory hold up or will it need to be replaced by something more profound? In the case of the cosmic acceleration, where there is no appealing fiducial model, the push for greater precision takes on an even greater importance.

Since scientists in the United States discovered evidence for cosmic acceleration over a decade ago, the US has been the leader in the field of dark energy studies. The 2006 Dark Energy Task Force report \[8\] provided a systematic discussion of experimental approaches to dark energy and identified a sequence of “Stage III” and “Stage IV” dark energy experiments to build on those then in progress. Stage III surveys were designed to address systematics with the goal of statistics-limited constraints from four independent

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\[3\] Theories can be written either in the Jordan frame, in which matter couples to the metric only through $\sqrt{g}L_m$, or in the Einstein frame, in which the gravitational part of the action remains Einstein-Hilbert and there are additional couplings of the new degrees of freedom to matter.
probes. Each of these probes will be developed over the ensuing decade to the point at which Stage IV surveys will enable sub-percent level consistency checks for all probes, some of which will have reached their cosmic variance limit (i.e., that we have only a single observable universe to probe). This staged categorization was reiterated in the more recent Community Dark Energy Task Force Report [9], which particularly emphasized the importance of complementing planned imaging experiments with spectroscopic experiments. Fig. 1 illustrates the timelines for several of the major dark energy experiments in which US scientists are playing an important or leading role with further details provided in a separate document [10].

Figure 1. A timeline of Stage III and Stage IV dark energy experiments – photometric and spectroscopic – in which US scientists are playing an important or leading role. Most of the projects are ground-based with either US leadership (BOSS, DES, HETDEX, eBOSS, DESI, LSST) or active participation (HSC, PFS). The two space missions are Euclid, led by the ESA with a NASA-sponsored team of US participants, and WFIRST, led by NASA.

Major gains beyond the current road map of dark energy projects will require advances on one of a number of fronts. One potential avenue is to develop new techniques or probes of cosmology and new physics that have not yet been developed; data from many cosmological surveys have been used for tests not envisioned when those experiments were first designed, and we anticipate that that trend will continue [11]. Another path is to obtain complementary information that will enhance planned experiments; for instance, a targeted program of spectroscopic redshift measurements can enhance the dark energy constraints from LSST compared to its baseline capabilities [12]. Finally, new instrumental capabilities – e.g., methods that suppress flux from emission lines from the night sky, which are much brighter than distant galaxies at infrared wavelengths, or new detector technologies that would measure photon energy rather than just quantity – could enable powerful new dark energy experiments to be conducted at reduced costs. Modest investments in these avenues may yield large potential payoffs in the future.

The following sub-sections summarize the physical probes and measurements that speak to dark energy science. Each of these approaches is detailed further in a separate document [13, 14, 11, 15]. Taken together, these measurements and the projects that enable them constitute an all-out attack on the problem of cosmic acceleration. The program planned over the coming decade guarantees continuing US leadership in this field.
28.2 Dark Energy

28.2.1 Distances

The relationship between the redshift and distance of an object is one of the primary tests of the expansion history of the Universe, and therefore played a key role in the discovery of the accelerating Universe. The simple graph of the distance scale of the Universe as a function of redshift, indicating the evidence for cosmic acceleration, has become an iconic plot in the physical sciences. The data for this plot so far come from measurements of Type Ia supernovae and baryon acoustic oscillations (BAO) and these will be the sources of streams of data in coming years. DES and LSST will provide an essentially limitless supply of supernova, thousands, then hundreds of thousands. The DES collaboration and LSST-DESC will coordinate the spectroscopic classification of a fraction of these objects. The challenge is to make measurements thoroughly enough to mitigate systematic problems, especially those that depend on redshift. Detailed studies of nearby supernovae are beginning to provide clues for how to do this. Much would be gained if observations could be made from space, but a substantial gain will be also achieved if we make ground-based observations that avoid the atmospheric lines in the near infrared.

The subtle pattern of anisotropy in the cosmic microwave background, just one part in $10^5$, is mapped at the two dimensional boundary of a three-dimensional feature, the fluctuations in matter density throughout space. The counterpart of the oscillations in the CMB power spectrum is a peak in the correlation between the densities at points separated by 153 Mpc measured relative to the scale size of the Universe, left behind by baryon acoustic oscillations in the early Universe. This very large meter stick can be observed as far out as redshifts $z = 1.6$ using galaxies as traces of matter density, and even out to $z = 3$ using light from quasars. The current Baryon Oscillation Spectroscopic Survey (BOSS) is likely to report a distance measurement soon with 1% accuracy. The eBOSS survey is designed to extend the reach of BAO measurements to higher redshifts. The Stage-IV BAO experiment, DESI, should provide more than 30 similarly accurate independent distance measurements.

Fig. 2 shows the current and projected future constraints on cosmic distances using these two techniques. If our basic understanding is correct, the supernova and BAO measurements should be in absolute agreement. The distance-versus-redshift curve of the Universe is fundamental and exploring it with completely different techniques is essential. These stunning measurements will allow for percent level determination of the equation of state and be extremely sensitive to evolution of the dark energy at earlier times. By pinning down the distance-redshift relation, they will also allow for apples-to-apples comparisons of modified gravity vs. dark energy models using the growth of structure.

28.2.2 Growth of Structure

The quantity and quality of cosmic structure observations have greatly accelerated in recent years, and further leaps forward will be facilitated by imminent projects. These will enable us to map the evolution of dark and baryonic matter density fluctuations over cosmic history. The way that these fluctuations vary over space and time is sensitive to several pieces of fundamental physics: the primordial perturbations generated by GUT-scale physics; neutrino masses and interactions; the nature of dark matter; and dark energy. We focus on the last of these here: the ways that combining probes of growth with those of cosmic distances will pin down the mechanism driving the acceleration of the Universe.

If the acceleration is driven by dark energy, then distance measurements provide one set of constraints on $w$, but dark energy also affects how rapidly structure grows. Upcoming surveys are therefore designed to probe $w$ in two distinct ways: direct observations of the distance scale and the growth of structure, each
Figure 2. Current and projection for future uncertainties on cosmic distance as a function of redshift.
complementing the other on both systematic errors and dark energy constraints. A consistent set of results will greatly increase the reliability of the final answer.

**Figure 3.** Constraints on the growth of density fluctuations in the universe with errors projected from DESI. The curves show the derivative of the logarithmic growth with respect to logarithmic scale factor — a quantity readily measured from the clustering of galaxies in redshift space — as a function of redshift. We show theory predictions for the standard \( \Lambda \)CDM model, as well as for two modified-gravity models, the Dvali-Gabadadze-Porratti (DGP) model [16], and for the \( f(R) \) modification to Einstein action [17]. Because growth in the \( f(R) \) models is generically scale-dependent, we show predictions at wave numbers, \( k = 0.02 \) h Mpc\(^{-1}\) and \( k = 0.1 \) h Mpc\(^{-1}\). LSST projects to impose constraints of similar excellent quality on the growth function \( D(a) \).

If cosmic acceleration is driven by modified gravity, then probes of structure become even more important. Generically, modified gravity models are able to reproduce any expansion history that can be attained in dark energy models, but at the cost of altering the growth of structure. How rapidly structure grows is quantified by the dimensionless growth function \( D(a) \). Figure 3 illustrates how different models make predictions that differ from those of general relativity even though the distance-redshift relation in all these models is identical. The growth of structure then will be able to distinguish modified gravity from dark energy as an explanation for the cosmic acceleration.

Fig. 3 projects constraints from a spectroscopic survey that measures the local velocities of galaxies by observing redshift space distortions. Similarly powerful constraints are projected from photometric surveys that are dedicated to measuring the shapes of galaxies and therefore are sensitive to the signal from weak gravitational lensing.

Achieving these powerful constraints will require both wide field imaging and spectroscopic redshift surveys, as depicted in Fig. 1. The results will pin down far more than the equation of state to percent level accuracy, although this in itself will be an important clue as to whether the cosmological constant or an alternative is driving the acceleration of the Universe. We will learn also whether the equation of state is varying with time and whether dark energy was relevant at high redshift. The surveys will probe cosmological perturbations as a function of both length and time, opening up dozens of possible failure modes for GR-based dark energy. If any of these is reliably detected to differ from the GR prediction, we will have a revolution on our hands.
28.2.3 Novel Probes

Surveys enabling the twin probes of distances and structure can distinguish between modified gravity and dark energy on cosmological scales. It has become apparent over the last few years that non-cosmological tests can also play an important role in determining the mechanism driving the acceleration of the Universe. The basic idea is that gravity is known to reduce to general relativity (GR) on Solar System scales, so any modified gravity model must have a screening mechanism, wherein the additional forces operative on large scales are suppressed in the solar system. Indeed, many of the models have screening built into them, so the solar system constraints can be naturally satisfied. The key issue is the nature of the transition from large (modified gravity) to small (GR) scales, and how that transition can be detected observationally.

Screening mechanisms typically utilize some measure of the mass distribution of halos, such as the density or the Newtonian potential, to recover general relativity (GR) well within the Milky Way. This leaves open the possibility that smaller halos, the outer parts of halos, or some components of the mass distribution, are unscreened and therefore experience enhanced forces. For a given mass distribution, unscreened halos will then have internal velocities and center of mass velocity larger than predicted by GR. Deviations from GR are typically at the ten percent level, with distinct variations between different mechanisms in the size of the effect and the way the transition to GR occurs. It is important to note that observable effects are typically larger on galaxy scales than on large (cosmological) scales or at high redshift. The comparison of dynamical and lensing masses provides a powerful test that is being implemented on a wide range of scales: from individual galaxies to large-scale cross-correlations that are also discussed in the Growth of Structure and Cross-Correlations sections.

![Figure 4](image-url)

*Figure 4.* Astrophysical [18, 19, 20] and cosmological [21, 22, 23] limits on chameleon theories, in particular \( f(R) \) models. The spatial scale on the x-axis gives the range of length scales probed by particular experiments. The parameter on the y-axis is the value of the field that mediates the additional force in units of the Planck mass, or equivalently the range of the additional force. The rectangles show excluded regions; the two rectangles with dots are meant to indicate preliminary results from ongoing work.

Fig. 4 shows a range of tests that have been implemented in one particular model, f(R) gravity, with Chameleon screening. The screening in this model, as in many others, depends on the value of a field in the region of interest. The y-axis in Fig. 4 depicts the value of this field in Planck units divided by a coupling constant, while the x-axis shows that the model has been probed on scales ranging from the Solar System...
all the way out to cosmological scales, with all tests to date verifying GR. In the next decade essentially the entire accessible parameter space of Chameleon theories can be probed using the tests shown in the figure. Tests for the other important screening mechanism, Vainshtein screening, are at early stages, with potential for rapid progress.

The program of testing gravity theories via these novel probes is in its infancy, but it is becoming increasingly clear that even modest investments in non-cosmological observations have enormous potential to contribute to the cosmic acceleration problem.

### 28.3 Inflation

Cosmic inflation is the leading theory for the earliest history of the universe and for the origin of structure in the universe. Current observations of the large-scale distributions of dark matter and galaxies in the universe and measurements of the Cosmic Microwave Background (CMB) are in stunning agreement with the predictions of inflation. The next generations of experiments in observational cosmology are poised to explore the detailed phenomenology of the earliest moments of the universe.
Figure 5. Map of a galaxy cluster [24] using three probes: (i) weak gravitational lensing (blue contours with labels showing the projected density $\kappa$); (ii) hot gas as measured by the Sunyaev-Zel’dovich distortion of the CMB (white contours with labels giving signal to noise); and (iii) galaxies as observed in three optical bands (background).

Figure 6. Expected signal levels for the CMB Polarization $E$-mode (red, solid), inflationary gravity-wave $B$-mode (blue, solid), and lensing $B$-mode (green, solid) signals. The gravitational wave $B$-mode signals are shown for tensor-to-scalar ratios of $r = 0.001$ (the Stage IV goal) and $r = 0.01$ (the boundary between small-field and large-field inflation models). The lensing $B$-mode signal is shown as a band encompassing the predicted signal for values of the sum of neutrino masses $0 \leq \sum m_\nu \leq 0.1$eV. De-lensing by a factor of 4 in amplitude is shown schematically by the green arrow, with the residual signal at $\ell \leq 200$ (where the de-lensing is critical to the constraint on $r$) shown by the green, long-dashed line. The black, short-dashed line shows the level of current 95% upper limits on $B$ modes from WMAP, BICEP, QUIET and QUAD experiments. The brown, long-dashed lines show the expected polarized foreground contamination at 95 GHz for the cleanest 1% and 25% of the sky.
While the landscape of possible models for inflation is large, the theoretical underpinnings are well understood, and we are able to make concrete predictions for observable quantities. One key prediction is the existence of a background of gravitational waves from inflation that produce a distinct signature in the polarization of the CMB (the B-modes in Fig. 6). Under the so-called Lyth condition [25], all models in which the field driving inflation varies by an amount of order $m_{\text{Planck}}$ will produce gravitational wave (tensor) fluctuations that are at least 1% of the amplitude of density-fluctuation (scalar) power in the CMB ($r = 0.01$ in the figure). Definitive evidence one way or another as to the presence of tensor modes with amplitudes at or above this level would therefore serve as a lever to an infinite sequence of Planck suppressed operators. Such scales require a quantum gravity treatment and this will test string-theoretic mechanisms for large field inflation. If not detected it at least decides between two broad classes of models, since it reaches the Planck scale threshold. This motivates the design of a next-generation CMB experiment with the sensitivity and systematics control to detect such a signal with at least 5σ significance, thus ensuring either a detection of inflationary gravitational waves or the ability to rule out large classes of inflation models [26].

There are several other handles on the physics of inflation. Inflation generically predicts small deviations from a scale-invariant spectrum, and current measurements confirm this prediction at the 5-sigma level. DESI projects to obtain a 15-sigma detection thereby further reducing the range of allowed models. BOSS, eBOSS, and DESI will potentially constrain the running of the primordial spectrum (deviation from a pure power law) at the 0.2% level, a factor of five tighter than current constraints [27].

Non-gaussianity of the primordial perturbations can take many forms. The search for one – so-called local non-gaussianity – is particularly important because single field models of inflation generically predict negligible local non-gaussianity [26], so any detection will falsify a large class of models. Planck has placed strong upper limits on this and other forms of non-gaussianity consistent with these predictions. The upcoming surveys eBOSS, DESI, and LSST will constrain a variety of forms of primordial non-gaussianity on different spatial scales and be subject to different systematics than the CMB. They will therefore pave the way for even more stringent bounds on inflationary models.

### 28.4 Neutrinos

One of the most remarkable aspects of physical cosmology is that the study of the largest physical structures in the Universe can reveal the properties of particles with the smallest known cross-section, the neutrinos. At the simplest level, this cosmological sensitivity to neutrino properties is due to the fact that the neutrino cosmological number density is so large as to be second only to CMB photons. More specifically, the properties of neutrinos alter the effective energy density of cosmological radiation and therefore the amplitude, shape and evolution of matter perturbations, leading to changes in observables in the CMB anisotropies and in measures of large-scale structure.

The CMB and large scale structure (LSS) measured in galaxy surveys are sensitive to the sum of the neutrino masses and the number of species produced in the early Universe (dubbed $N_{\text{eff}}$). These observations are therefore complementary to laboratory probes of neutrinos, which measure mass differences and potentially CP-violation. CMB experiments are sensitive enough to the neutrino energy density to rule out $N_{\text{eff}} = 0$ at more than 10-sigma; that is, these experiments have already (indirectly) detected the cosmic neutrinos. Together with priors on the redshift-distance relation from galaxy surveys like BOSS, these experiments also place a stringent upper limit on the sum of the neutrino masses, currently around 0.23 eV [28].

A global fit to solar and atmospheric neutrino flavor oscillations in the standard 3-generation model determines two mass differences so a third parameter, which can be taken as either the sum of the masses or the lightest mass, is unknown. Due to a sign ambiguity in one of the mass differences, there are two
discrete possibilities (normal and inverted hierarchy as depicted in Fig. 7) for the relationship between these two parameters. As indicated in Fig. 7, upcoming CMB experiments and LSS surveys will unambiguously detect the sum of the masses if the hierarchy is inverted and will likely do so at greater than 3-sigma even if nature has chosen the normal hierarchy. Measures of the power spectrum from DESI with a Planck prior could improve the current constraint on $\sum m_\nu$ to 17 meV, and a Stage IV CMB survey combined with BAO measurement from DESI project to achieve similar sensitivity [29]. Fig. 7 highlights another aspect of the complementarity of laboratory experiments and cosmological probes: if the sum of the masses is found to be 0.15 eV by the latter, then it will take the former to determine the lightest neutrino mass by identifying whether the hierarchy is normal or inverted.

![Figure 7](image.png)

**Figure 7.** Current constraints and forecast sensitivity of cosmology to the sum of neutrino masses. In the case of an “inverted hierarchy,” with an example case marked as a diamond in the upper curve, future combined cosmological constraints would have a very high-significance detection, with $1-\sigma$ error shown as a blue band. In the case of a normal neutrino mass hierarchy with an example case marked as diamond on the lower curve, future cosmology would still detect the lowest $\sum m_\nu$ at greater than 3-$\sigma$.

Short baseline neutrino oscillation results hint at a richer neutrino sector than three active neutrinos participating in flavor oscillations, with one or more sterile flavors also participating [30, 31, 32, 33]. These same future CMB experiments will achieve a 1-sigma error of $\Delta N_{\text{eff}} = 0.027$, which will complement future sterile neutrino searches and inform model building (since sterile neutrinos might be detected in the laboratory even though they were not produced in the early universe and vise versa).

### 28.5 Conclusions

Cosmological surveys are sensitive to fundamental physics. To date, this basic fact has led to the discovery of the accelerating universe, strong evidence for an epoch of early acceleration near the GUT scale, the indirect detection of the cosmic neutrino background, and the most compelling evidence for non-baryonic dark matter. However, surveys to date have measured only a fraction of all information available. If the Universe were contained in an area the size of the United States, galaxy maps so far would have surveyed the city of Birmingham, Alabama. CMB experiments have provided low-noise maps of the temperature down to angular scales of order a tenth of a degree. Strategic, valuable information remains unmined...
28.5 Conclusions

in higher resolution temperature maps and the virtually uncharted polarization field. We have outlined the projections for how this extra information will constrain dark energy, neutrinos, and inflation; these projections are extraordinary. But even they ignore the very real possibility that future experiments on the cosmic frontier will do just what their predecessors have done: discover something fundamentally new!

The community has rallied behind previous reports [8, 34, 9] which are consistent with the current consensus to support the following key steps:

- **REMAIN A LEADER IN DARK ENERGY RESEARCH**
  The U.S. played the leading role in discovering the acceleration of the Universe, as was recognized by the 2011 Nobel Prize. The acceleration remains a mystery, whose solution may usher in a revolution in either our theory of gravity or our understanding of particle physics. Different classes of theories make different predictions for the growth of structure given a redshift-distance relation. A combination of spectroscopic and photometric surveys can determine both distances and structure growth, so will help pinpoint the new physics driving the acceleration of the universe. The current suite of surveys, Stage III, will be the first to implement the vision of multiple probes and small systematics. This vision will be realized fully with the Stage IV surveys (LSST and DESI), as they reach the level where exquisite-precision dark energy constraints from different probes, in some cases approaching the cosmic variance limit, can be checked for consistency. Therefore, the community strongly supports continuing the program of Stage III and Stage IV dark energy experiments, and moving forward as quickly as possible with the construction of LSST and DESI.

- **BUILD A GENERATION IV CMB POLARIZATION EXPERIMENT**
  Cosmic microwave background experiments can measure the sum of the neutrino masses and the energy scale of inflation, as well as constrain exotic physics such as early dark energy and extra neutrino species. After the current generation of small scale experiments complete data taking near the end of the decade, the community understands that the next generation experiment – one that can pin down neutrino masses and the scale of large field inflationary models – requires a nationwide coherent effort. Moving from thousands to hundreds of thousands of detector elements will require the involvement of the national laboratories working together with the university community.

- **EXTEND THE REACH**
  With small additional investments the dark energy program can be augmented in three important ways. A targeted spectroscopic campaign designed to optimize and calibrate methods of redshift estimation from imaging surveys can enhance their science returns beyond their nominal capabilities [35]. Second, a continued investment in instrumentation R&D will allow the community to do more science for less money and to be ready to pounce on future discoveries. Finally, a suite of novel probes of gravity and dark energy can discover signatures of modified gravity and new physics in the dark sector. We have the opportunity to catalyze the next generation of tests by supporting work at the interface of theory, simulation and data analysis, and making small enhancements to the dark energy survey program.
References


Community Planning Study: Snowmass 2013


