Highlights on hadron physics at CLAS

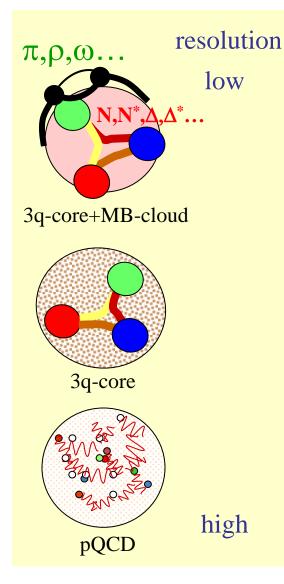
K. Hicks (Ohio U.) Hadron 2011 Conference June 16, 2011

Outline

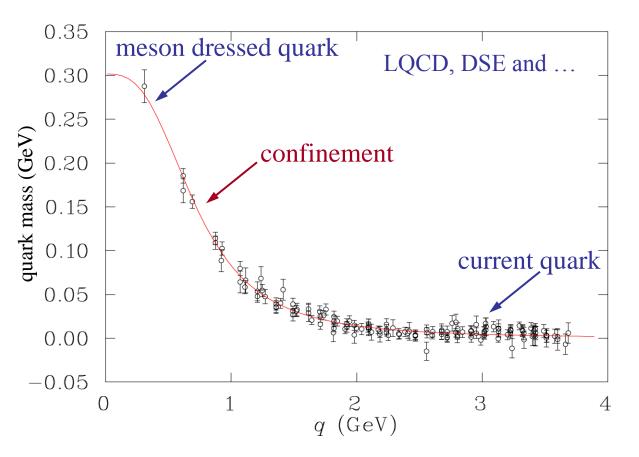
- Meson-Baryon Cloud (MBC) Effects
 - New results on baryon photocouplings
 - Need for coupled-channels analysis
- Spectroscopy of baryons and mesons
 - New and future analysis of ${\rm K}\Lambda$
 - New data: 2π , K*Y, KY*, etc.
- Future upgrade to CLAS12

- Probe transition from MBC to quark core.

Hadron Structure with Electromagnetic Probes



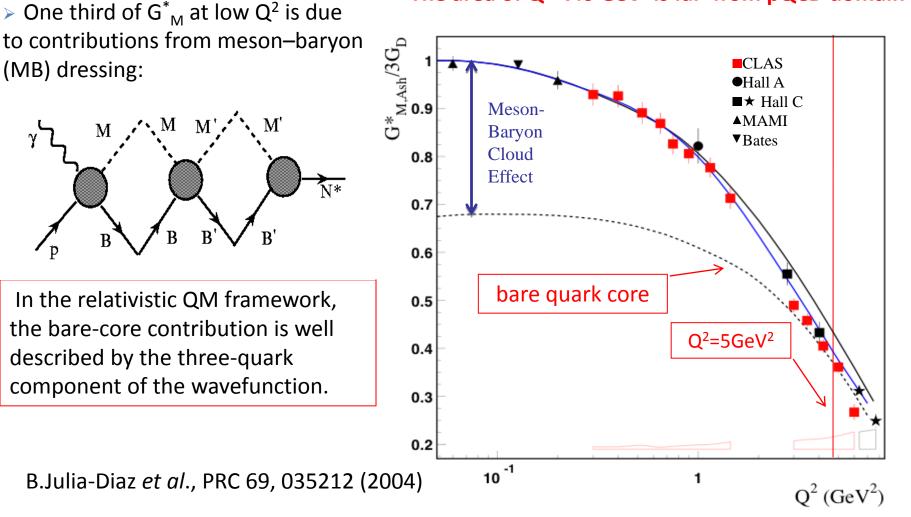
Quark mass extrapolated to the chiral limit, where *q* is the momentum variable of the tree-level quark propagator (curve=DSE, data=LQCD).



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N Δ Transition Form Factor (G_M) from EBAC analysis



The area of Q²<7.0 GeV² is far from pQCD domain

16 June 2011

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Dynamical coupled-channels model of EBAC

For details see Matsuyama, Sato, Lee, Phys. Rep. 439,193 (2007)

+

✓ Partial wave (LSJ) amplitude of a \rightarrow b reaction:

$$T_{a,b}^{(LSJ)}(p_a, p_b; E) = V_{a,b}^{(LSJ)}(p_a, p_b) + \sum_c \int_0^\infty q^2 dq V_{a,c}^{(LSJ)}(p_a, q) G_c(q; E) T_{c,b}^{(LSJ)}(q, p_b; E)$$

coupled-channels effect

✓ Reaction channels:

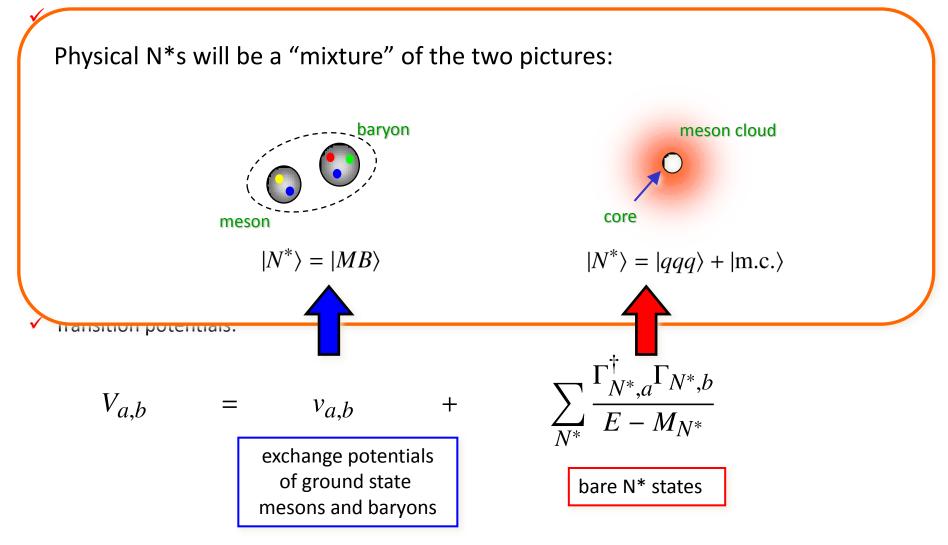
$$a, b, c = (\gamma^{(*)}N, \pi N, \eta N, \pi \Delta, \sigma N, \rho N, K \Lambda, K \Sigma, \cdots) \frac{\pi \pi N}{\pi n N}$$

✓ Transition potentials:

$$V_{a,b} = v_{a,b} + \sum_{N^*} \frac{\Gamma_{N^*,a}^! \Gamma_{N^*,b}}{E - M_{N^*}}$$
Meson-exchange potentials
(Derived from Lagrangians) bare N* states

Dynamical coupled-channels model of EBAC

For details see Matsuyama, Sato, Lee, Phys. Rep. 439,193 (2007)



Jefferson Lab Today

Jefferson Lab CLAS Detector

Hall B

Large acceptance spectrometer electron/photon beams

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Radiative Decay of Strange Baryons

General Motivation

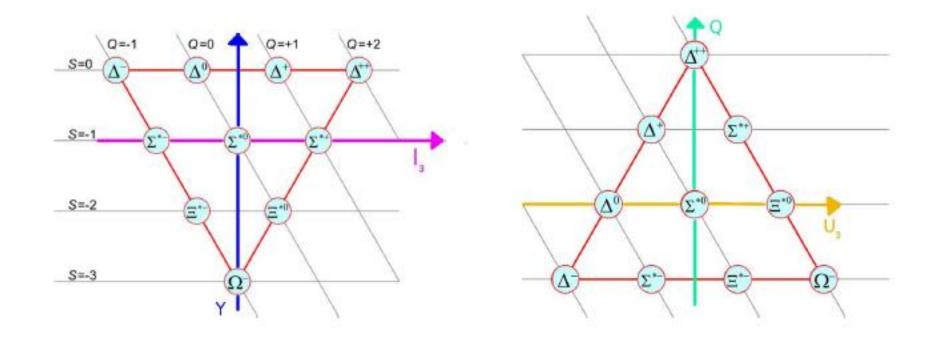
- Electromagnetic interactions are the cleanest way to access information on wavefunctions.
 - For <u>strange baryons</u>, the EM coupling can only be studied by measuring the decay.
- SU(6) wavefunctions provide a zeroth-order prediction for M1 transitions.
 - SU(6) symmetry provides a link between N* and
 Y* radiative decays: constrain SU(6) symmetries.

U-spin: connects Δ and Σ^* decays

- SU(3) has three equal symmetries:
 - I-spin: exchange of u and d quarks
 - U-spin: exchange of d and s quarks
 - V-spin: exchange of u and s quarks
- With respect to known symmetries:
 - I-spin conserves chiral symmetry (mass)
 - U-spin conserves EM symmetry (charge)
 - V-spin conserves neither chiral nor EM symmetry

Group structure of U-spin:

$SU(3)_f \supset SU(2)_I \times U(1)_Y \supseteq SU(2)_U \times U(1)_Q$



How is this useful?

It is possible using the U-spin SU(3) multiplet representation to obtain a prediction for the ratio of the $\Delta^0 \to n\gamma$ partial width to the $\Sigma^{*0} \to \Lambda\gamma$ partial width.

$$\begin{split} \left\langle \Delta^{0} | n \gamma \right\rangle &= \left\langle 1 - 1 | 1 - 1 \ 0 \ 0 \right\rangle = 1 \\ \left\langle \Sigma^{*0} | \Lambda \gamma \right\rangle &= -\frac{\sqrt{3}}{2} \left\langle 1 \ 0 | 1 \ 0 \ 0 \ 0 \right\rangle = -\frac{\sqrt{3}}{2}, \end{split}$$

leading to a ratio;

$$\frac{|\langle \Delta^0 | n\gamma \rangle|^2}{|\langle \Sigma^{*0} | \Lambda\gamma \rangle|^2} = \frac{4}{3},$$

including phase space factors (PRD 28 (1983) 1125)

$$\frac{\Gamma(\Delta^0 \to n\gamma)}{\Gamma(\Sigma^{*0} \to \Lambda\gamma)} = \left(\frac{M_n}{M_\Delta}\right) \left(\frac{M_\Lambda}{M_{\Sigma^{*0}}}\right)^{-1} \left(\frac{q_n}{q_\Lambda}\right)^3 \frac{4}{3} = 1.56.$$

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CLAS result

~10% uncert.

The final calculated ratio with all uncertainties is

$$R^{\Lambda\gamma}_{\Lambda\pi} = \frac{\Gamma[\Sigma^0(1385) \to \Lambda\gamma]}{\Gamma[\Sigma^0(1385) \to \Lambda\pi^0]} \neq 1.42 \pm 0.12(stat)^{+0.11}_{-0.07}(sys)$$

The electromagnetic decay partial width using the full width $\Gamma=36\pm5~{\rm MeV},$

$$\Gamma = 445 \pm 73(stat)^{+71}_{-66}(sys) \text{keV}$$

compared to the previous experimental result by Taylor,

$$\Gamma = 479 \pm 120(stat)^{+81}_{-100}(sys)$$
keV.

PhD thesis, Dustin Keller

Predictions from other models

Mode1	$\Delta \rightarrow N \gamma$	$\Sigma \to \Lambda \gamma$	$\Sigma^{*0} \to \Lambda \gamma$	$\Sigma^{*+} ightarrow \Sigma^+ \gamma$
NRQM[17, 23, 26]	360	8.6	273	104
RCQM[24]		4.1	267	
χCQM[25]	350		265	105
MIT Bag[26]		4.6	152	117
Soliton[27]			243	91
Skyrme[28, 29]	309-326		157-209	47
Algebraic model[30]	341.5	8.6	221.3	140.7
Experiment:	660 +/- 60)	479 +/- 73	3

Interpretation

 Meson-Baryon Cloud effects are substantial for the Y* resonances also.

– U-spin relation works better than QM, etc.

- We can now make a prediction for Σ^{*+} decay.
 - The Wigner-Eckart theorem requires that the branching ratios $\Delta^+ \rightarrow p\gamma$ and $\Delta^0 \rightarrow n\gamma$ are equal.
 - U-spin predicts a ratio:

$$\frac{\Gamma(\Delta^+ \to p\gamma)}{\Gamma(\Sigma^{*+} \to \Sigma^+ \gamma)} = \left(\frac{M_p}{M_\Delta}\right) \left(\frac{M_{\Sigma^+}}{M_{\Sigma^{*+}}}\right)^{-1} \left(\frac{q_p}{q_{\Sigma^+}}\right)^3 = 2.638.$$

New CLAS result

The final results for the ratio of the $\Sigma^{*+} \to \Sigma^+ \gamma$ to $\Sigma^{*+} \to \Sigma^+ \pi^0$ with systematic uncertainties is

 $R_{\Sigma^{*+}\to\Sigma^{+}\pi^{0}}^{\Sigma^{*+}\to\Sigma^{+}\gamma} = \frac{n_{\gamma}A_{\pi}(\Sigma\pi) - n_{\pi}A_{\gamma}(\Sigma\pi)}{n_{\pi}A_{\gamma}(\Sigma\gamma) - n_{\gamma}A_{\pi}(\Sigma\gamma)} = 11.95 \pm 2.21(stat)^{+0.52}_{-1.21}(sys)\%, (27)$ Preliminary!

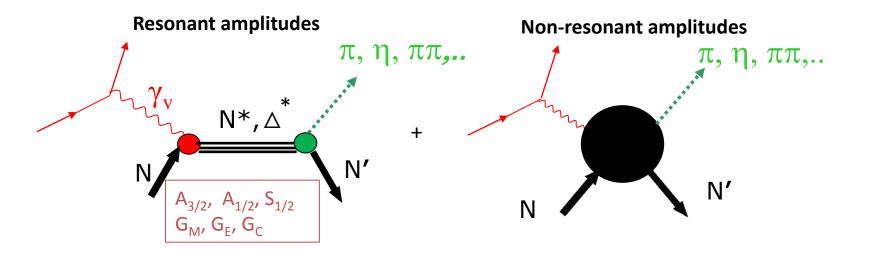
$$\Gamma_{\Sigma^{*+}\to\Sigma^{+}\gamma} = R^{\Sigma^{*+}\to\Sigma^{+}\gamma}_{\Sigma^{*+}\to\Sigma^{+}\pi^{0}} R(\Sigma^{*}\to\Sigma^{+}\pi^{0}) \Gamma_{Full} = 250 \pm 56.9(stat)^{+34.3}_{-41.2}(sys) \text{keV}.$$

This implies that the U-spin prediction for the partial width of the electromagnetic decay using the width of the $\Delta^+ \to p\gamma$ decay is,

$$2.638^{-1} \times \Gamma(\Delta^+ \to p\gamma) = 2.638^{-1} \times 660 \pm 60 = \frac{250 \pm 23 \text{keV.}}{23 \text{keV.}}$$
(29)

The MBC in electroproduction

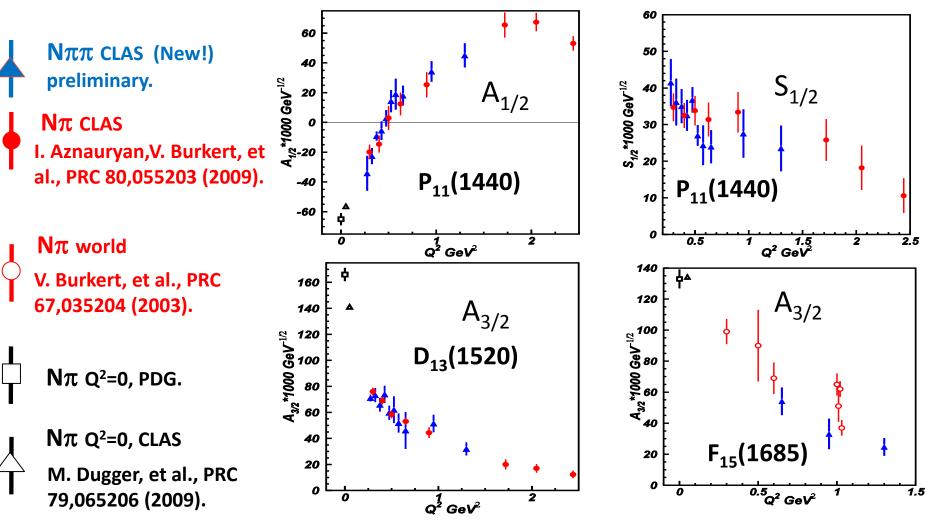
N* electrocouplings from analyses of exclusive channels



- Separation of resonant/non-resonant contributions
- N* 's can couple to various exclusive channels with entirely different nonresonant amplitudes, while their electrocouplings should remain the same.
- Consistent results from the analyses of N π and N $\pi\pi$ electroproduction channels show that model uncertainties are under control.

See the afternoon talk today by Victor Mokeev.

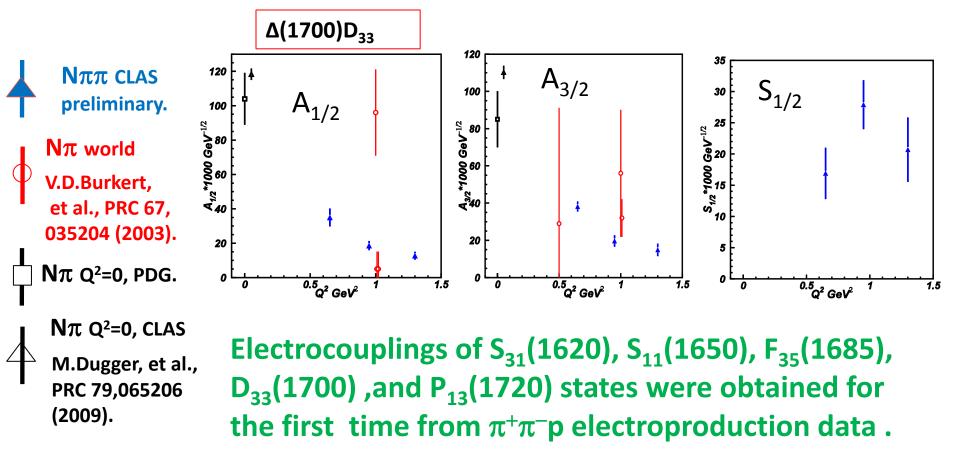
$\gamma_v NN^*$ electrocouplings from $N\pi/N\pi\pi$ production



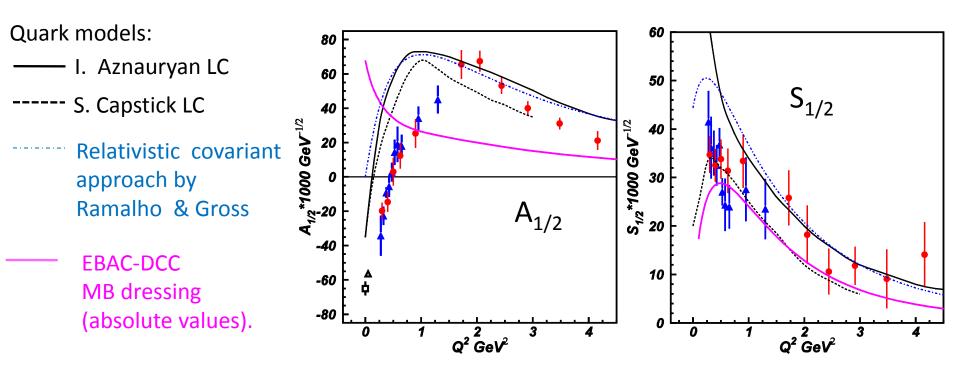
<u>Good agreement</u> between the <u> $N\pi$ and $N\pi\pi$ </u> channels.

N* electrocouplings are <u>measurable</u> and <u>model independent.</u>

High lying resonance electrocouplings from $\pi^+\pi^-p$



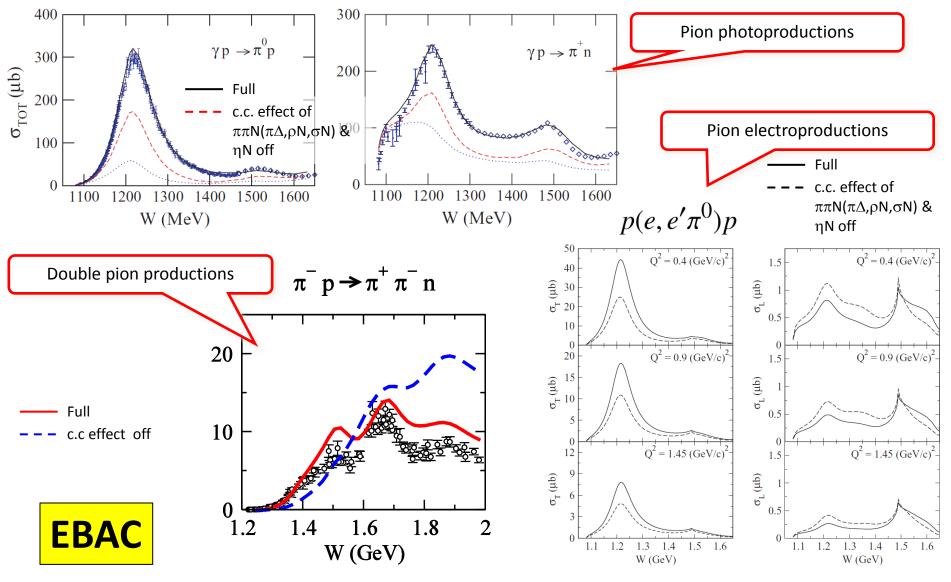
Mystery of P₁₁(1440) structure is solved



The electrocouplings are consistent with <u>P₁₁(1440)</u> structure as combination:
 a) quark core as a first <u>radial excitation of the nucleon</u>, and
 b) meson-baryon dressing.

MBC effects could explain the data at low Q².

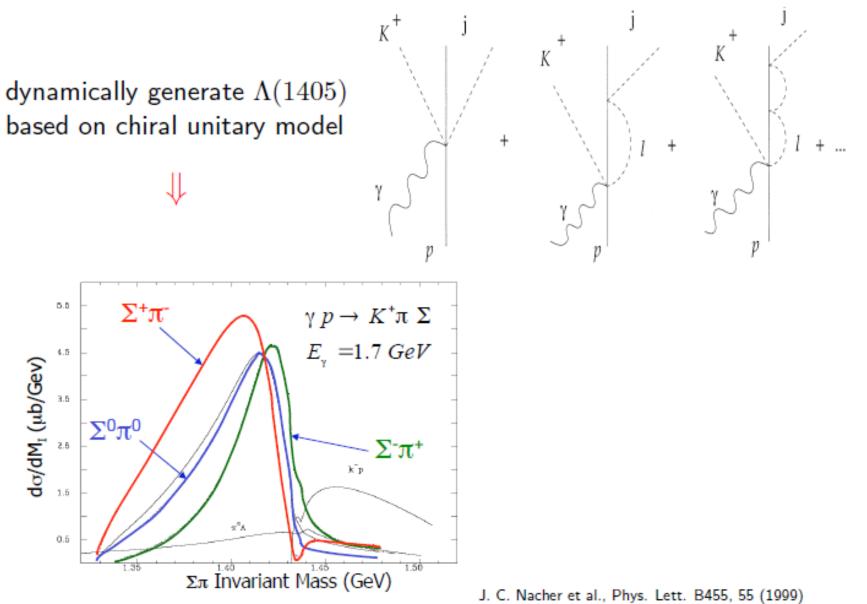
Coupled-channels effect in various reactions



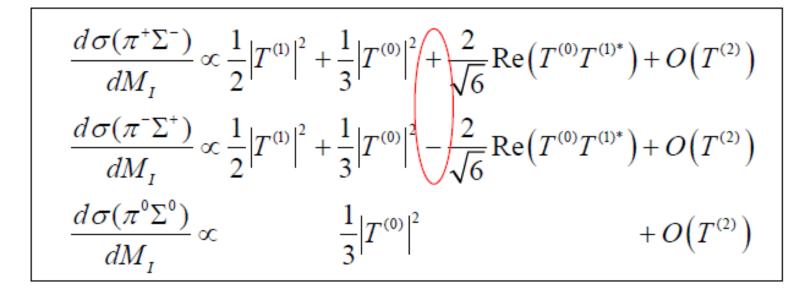
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Lineshape of the $\Lambda(1405)$

Chiral Unitary Coupled Channel Approach



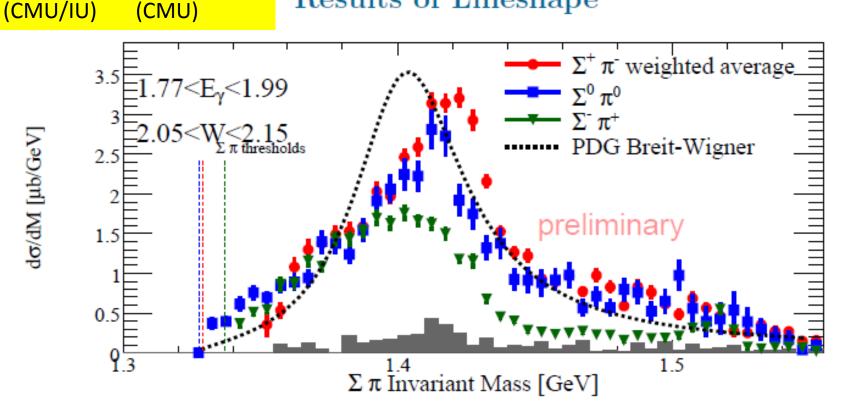
Difference in Lineshape



J. C. Nacher et al., Nucl. Phys. B455, 55

- Difference in lineshapes is due to interference of isospin terms in calculation (T^(I) represents amplitude of isospin I term)
- Distortion of the lineshape is connected to underlying QCD amplitudes that generate the $\Lambda(1405)$
- This analysis will measure all three $\Sigma\pi$ channels

Results of Lineshape



- lineshapes do appear different for each $\Sigma\pi$ decay mode
- $\Sigma^+\pi^-$ decay mode has peak at highest mass, narrow than $\Sigma^-\pi^+$
- lineshapes are summed over acceptance region of CLAS
- difference is less prominent at higher energies

K. Moriya & R. Schumacher

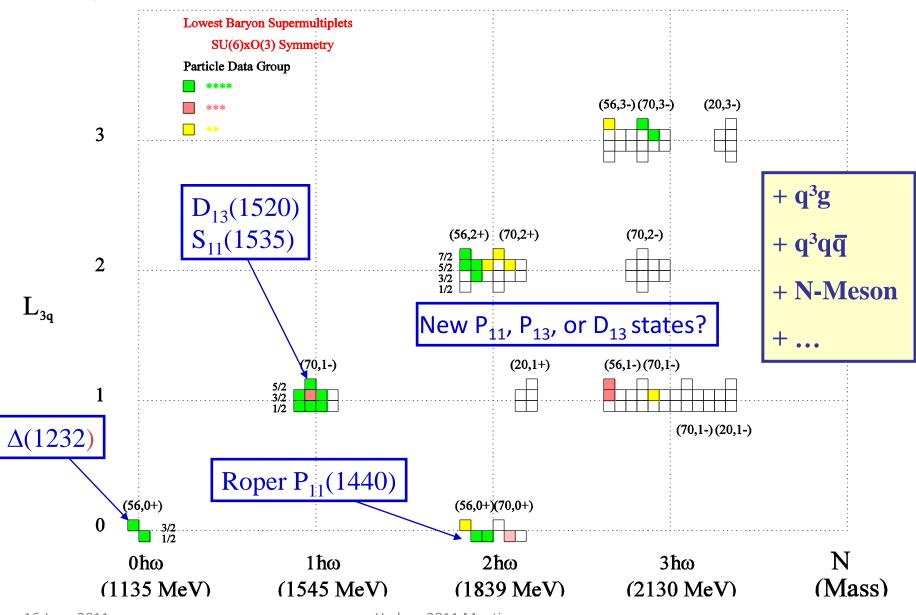
Interpretation

- For a baryon resonance with a single pole, isospin symmetry $\rightarrow \Sigma^+\pi^- = \Sigma^-\pi^+$.
- The data favor a dynamically-generated resonance (two-pole solution $\Sigma\pi$ and NK).

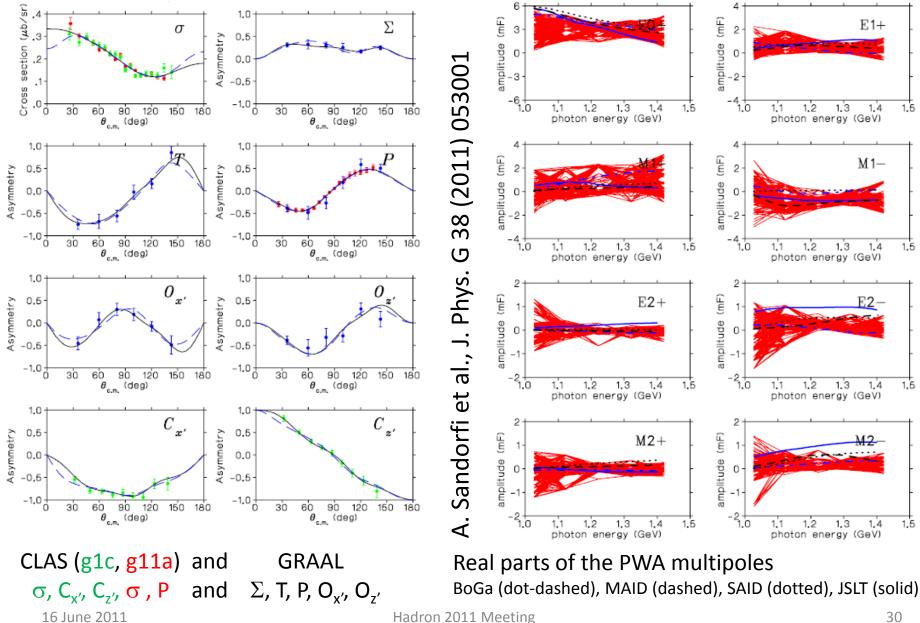
– Evidence of MB coupled-channels effects?

Spectroscopy with KY, K*Y and KY*

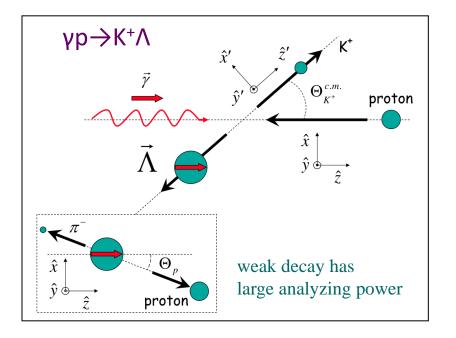
Quark Model Classification of N*



Amplitude Uncertainty in $\dot{\gamma} p \rightarrow K^+ \vec{\Lambda}$



FROST/HD $\vec{\gamma} \vec{N} \rightarrow \pi N$, ηN , $K \vec{\Lambda}$, $K \vec{\Sigma}$, $N \pi \pi$



Discussed Monday, talk by V. Cede

- Process is described by 4 complex, parity conserving amplitudes
- ➢ 8 well-chosen measurements are needed to determine amplitude
- ➢ For hyperon finals state 16 observables will be measured in CLAS ➡ large redundancy in determining the photoproduction amplitudes ➡ allows many cross checks
- ➤ 8 observables measured in reactions without recoil polarization

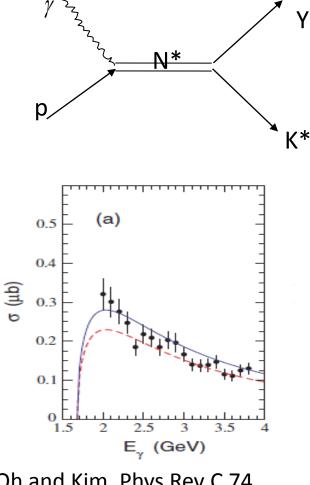
Photon beam			Target		Recoil		Target - Recoil									
					<i>x'</i>	у'	<i>Z</i> '	<i>x'</i>	<i>x'</i>	<i>x'</i>	у'	у'	у'	Z'	Ζ'	<i>z</i> ′
		x	У	Z		1	1	x	У	Z	x	У	Z	x	У	Z
unpolarized	σ ₀		Т	997-197-197-197-197-197-197-197-197-197-		Р		<i>T_x</i> ,	(1877-1887-1887-1887-1887-1887-1887-188	$L_{x'}$	That had had had had had had ha	zanananan D	911011001001001001001001	Tz'	97/197/197/197/197/197/197/197	$L_{z'}$
linearly P_{γ}	Σ	H	Р	G	<i>O</i> _{x'}	T	O z'	<i>L</i> _z ,	<i>Cz</i> '	T z'	E		F	$L_{x'}$	$C_{x'}$	T _{x'}
circular P_{γ}		F		E	$C_{x'}$		<i>C</i> _{z'}		0 z'		G		H		0 _{x'}	

K*+Y Photoproduction

Motivation:

1. Search for higher-mass N* resonances

- 2. Compare with KY and K*⁰Y
 - a. K^*Y coupling sensitive to $K_0(800)$.
 - b. $K_0(800)$ is part of the scalar nonet.
 - c. $K_0(800)$ has not been directly observed.

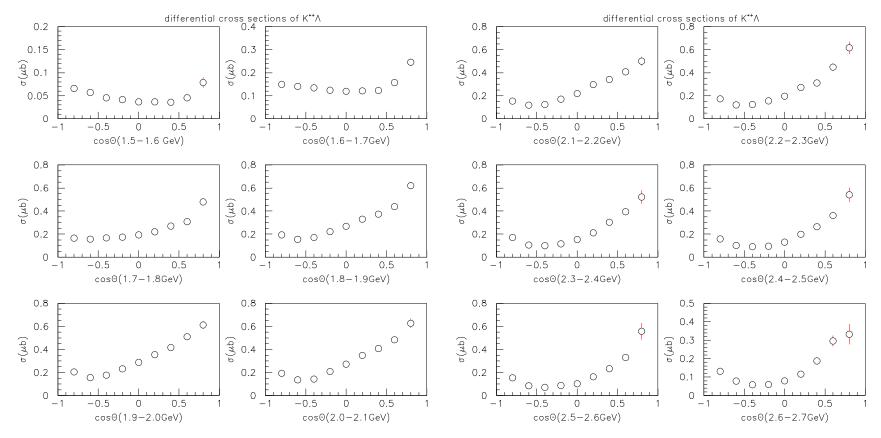


Oh and Kim, Phys Rev C 74, 015208(2006).

Differential cross sections of $K^{*+}\Lambda$:

1.5 --- 2.1 GeV

2.1 ---2.7 GeV

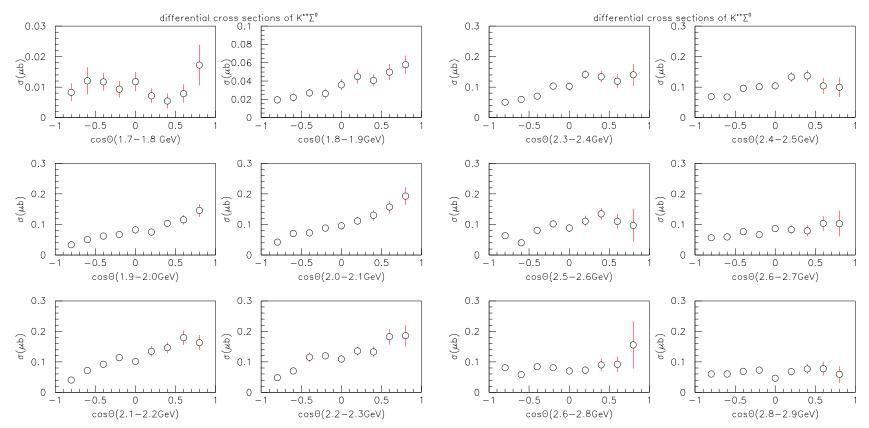


Preliminary result

Differential cross sections of $K^{*+} \Sigma^{0}$:

1.7 --- 2.3 GeV

2.3 --- 2.9 GeV

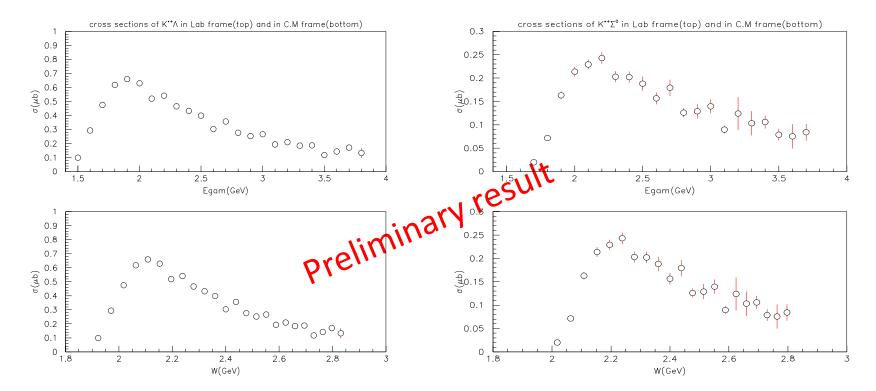


Preliminary result

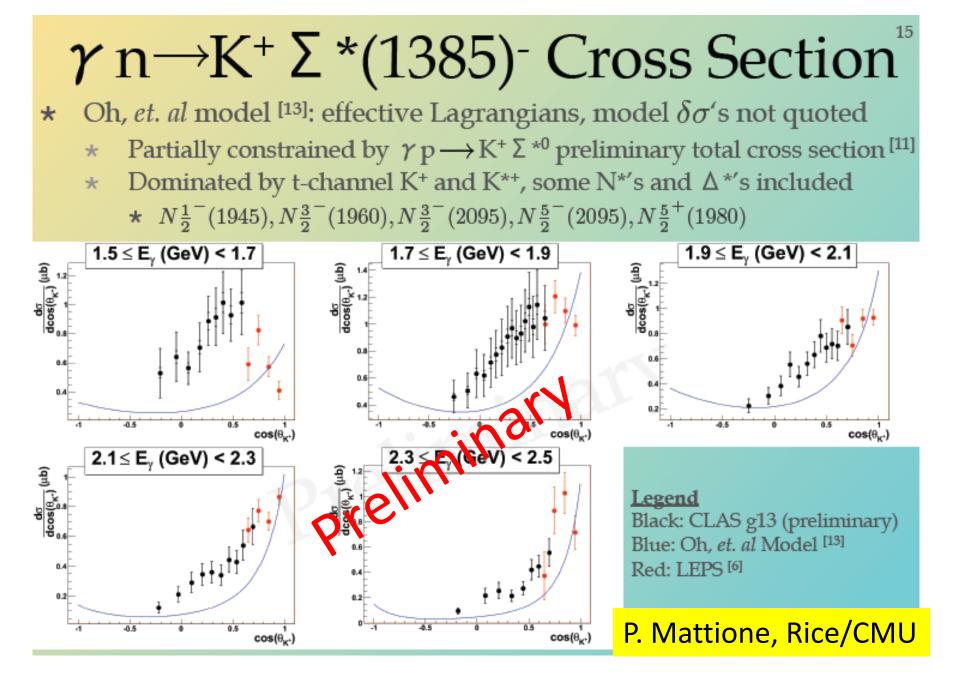
Total Cross Sections

 $\gamma + p \rightarrow K^{*+} + \Lambda$

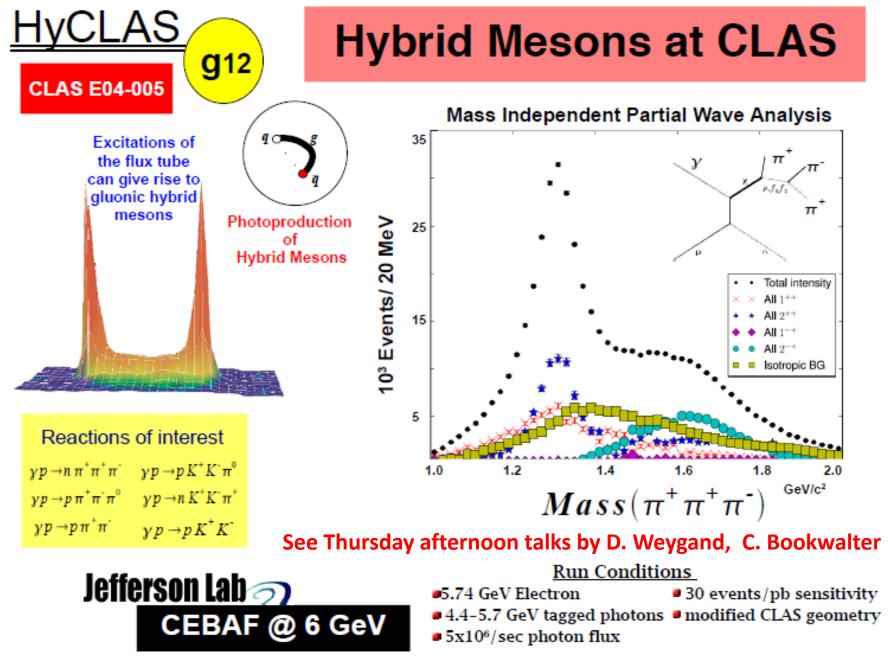
 $\gamma + p \longrightarrow K^{*+} + \Sigma^0$



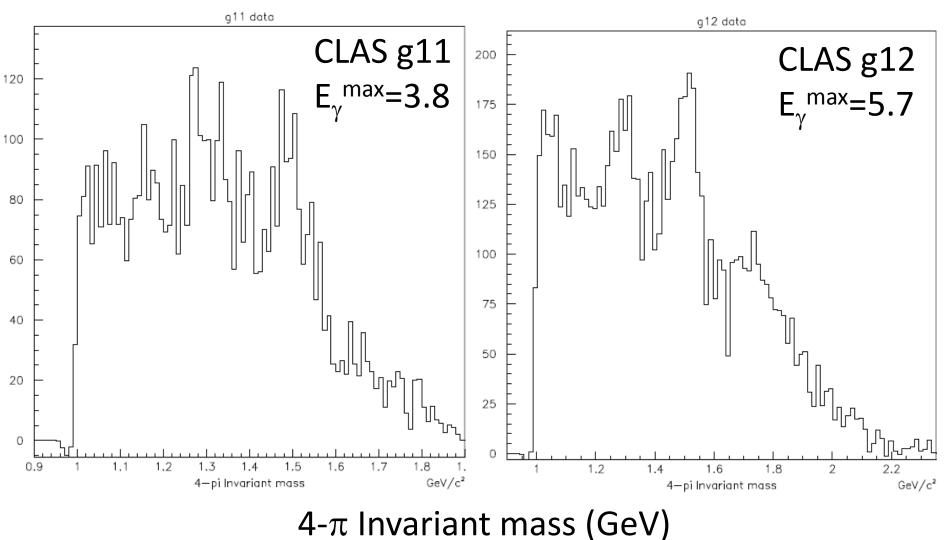
PhD thesis of Wei Tang



Meson Spectroscopy



New: Scalar Mesons from $\gamma p \rightarrow K_s K_s p$



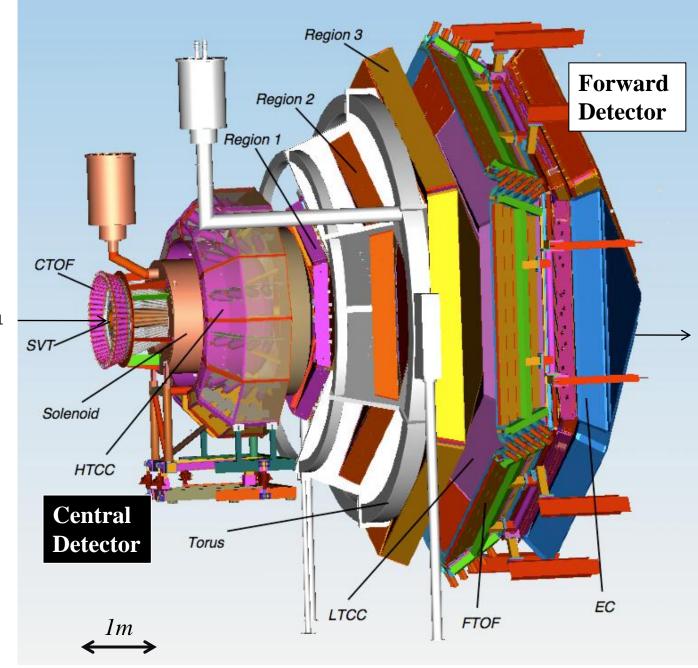
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The future: CLAS12

CLAS12

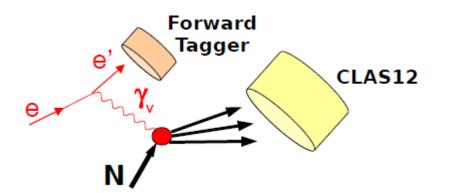
- \blacktriangleright Luminosity > 10³⁵ cm⁻²s⁻¹
- ➤ Hermeticity
- Polarization
- Baryon Spectroscopy
- Elastic Form Factors
- ➢ N to N* Form Factors
- GPDs and TMDs
- DIS and SIDIS
- Nucleon Spin Structure
- Color Transpareny

▶ ...



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Forward Photon Tagger for Spectroscopy



M. Battaglieri

$E_{scattered}$	0.5 - 4.5 GeV
θ	2.5^{o} - 4.5^{o}
ϕ	0° - 360°
ν	6.5 - 10.5 GeV
Q^2	$0.01 - 0.3 \text{ GeV}^2 \ (< Q^2 > 0.1 \text{ GeV}^2)$
W	3.6 - 4.5 GeV

Calorimeter + hodoscope + tracker

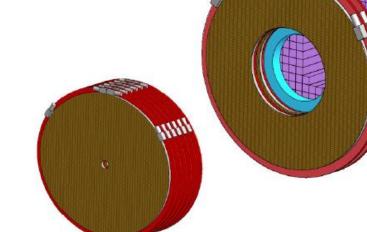
Electron energy/momentum

Photon energy (v=E-E') Polarization $\epsilon^{-1} \sim 1 + v^2/2EE'$

Veto for photons

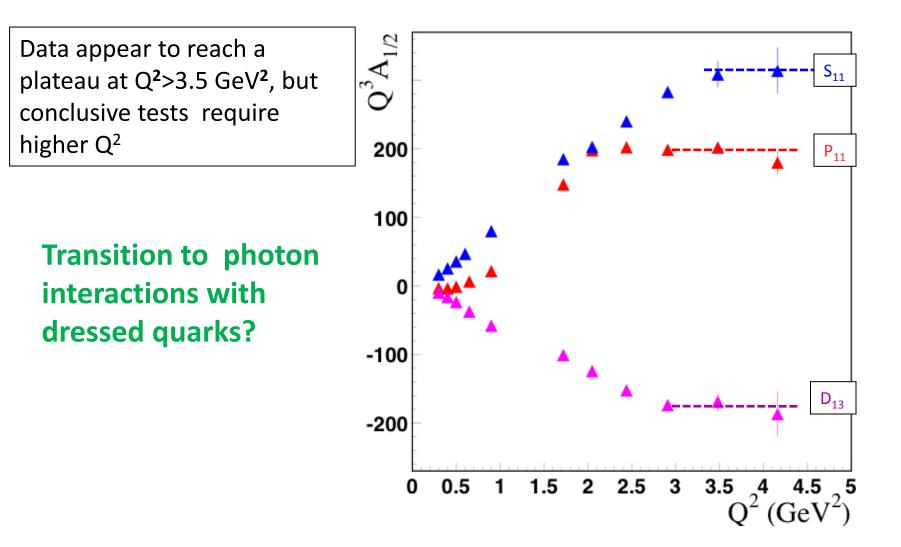
Electron angles Q²= 4 E E' sin² ϑ/2

Scattering plane



Rates in the forward tagger $L_e \sim 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ (N_y ~ 5 10⁸ y/s)

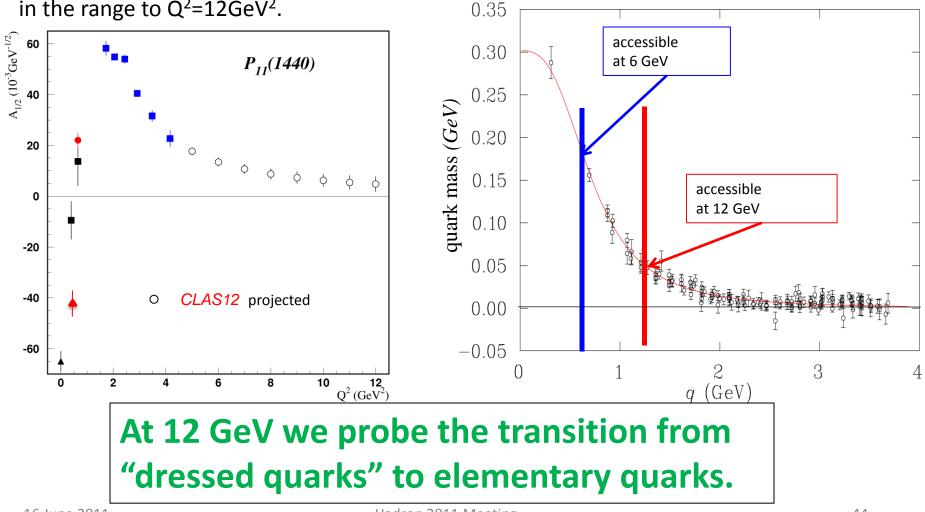
A new N* regime at $Q^2>3.5$ GeV²?



Resonance Transitions at 12 GeV

Experiment E12-09-003 will extend access to transition FF for many prominent states in the range to $Q^2=12GeV^2$.

Electromagnetic form factors are sensitive to the effective quark mass.



Summary

- The Meson-Baryon Cloud has significant effects on photon coupling observables
 We cannot ignore coupled-channel effects!
- There are precise new data on KY, K*Y, KY*
 - This will help the search for missing resonances
 - K^* data will determine the role of the K₀(800).
- PWA for mesons: exotics & scalar mixing
- Future: CLAS12: transition to current quarks.

Backup Slides

t-dependence of f(1500)

