

The Theory of X,Y, and Z States

XIV International Conference on Hadron Spectroscopy

R.P. Springer, Munich, June 2011

Observed spectrum of states (above DD)

Examples from theory landscape

Effective field theory for the X(3872)

Selected X(3872) predictions

Reviews: Brambilla et al.1010.5827; Voloshin 0711.4556;
(hadron11: 9 talks on 3872; 6 on other X,Y,Z)

$\underline{\underline{Y(4660)}}$
 $\underline{\underline{X(4630)}}$

X,Y,Z states from Table 9,
Brambilla et al. 1010.5827

$\underline{\underline{Y(4360)}}$
 $\underline{\underline{Y(4260)}}$

$\underline{\underline{X(4350)}}$

$\underline{\underline{Z(4430)^+}}$

$\underline{\underline{Y(4274)}}$ $\underline{\underline{Z_2(4250)^+}}$

$\underline{\underline{Y(4008)}}$

$\underline{\underline{X(4160)}}$
 $\underline{\underline{Y(4140)}}$ $\underline{\underline{Z_1(4050)^+}}$

$\underline{\underline{X(3872)}}$

$\underline{\underline{X(3915)...}}$ $\underline{\underline{X(3940)}}$

$D\bar{D}(3730)$

J^{PC} 1^{--} (1^{++}) $0/2^{++}$ $0/2^{?+}$ $?^{?+}$?

Techniques/Descriptions/Strategies

QCD Sum Rules

Non-relativistic QCD

Heavy Quark Effective Theory

Heavy Hadron Chiral Perturbation Theory

X-EFT

Lattice

Potential Models

Molecule

Mixtures

Baryonium

Tetraquark

Hybrids

Coupled channels

Hadrocharmonium

Hadron 2011 X,Y,Z Theory talks/posters (to see or to have seen)

F. De Fazio (Plenary II) EFT radiative decays
X(3872)

C.Zanetti (Heavy Hadrons I) QCD Sum Rules
X(3872)

H. Noya (Quarkonia I) Diquark Cluster model (4-q- states)
X(3940),X(4260),X(4350),X(4430),X(4660)

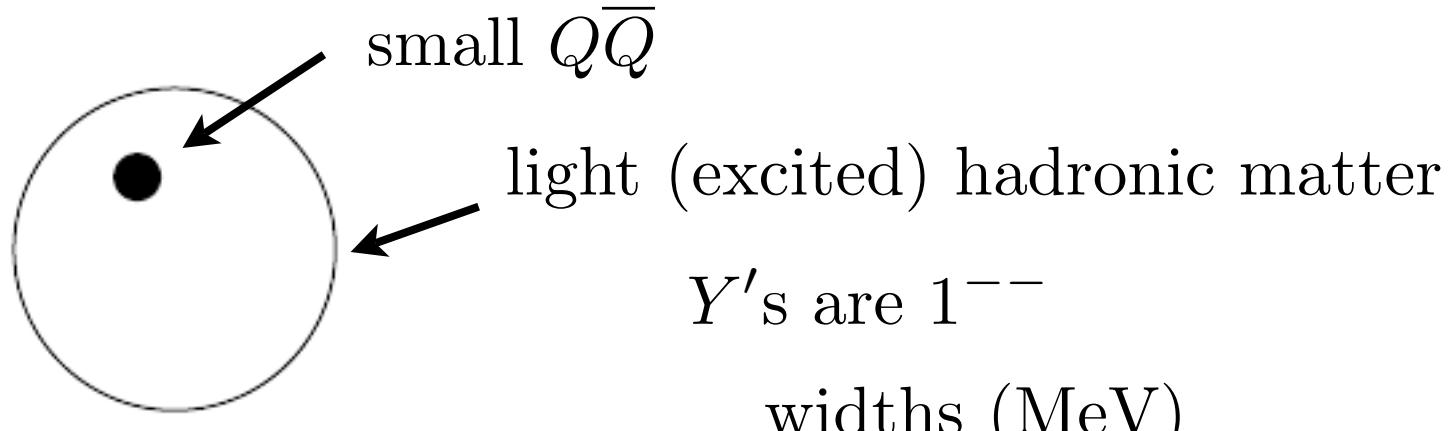
R. Molina (Poster) dynamically generated
Y(3940),Z(3930),X(4160)

D. Rodriguez Entem (Quarkonia 4) Molecules
X(3872),X(3915),X/Y(3940),Y(4008)

M. Karliner (Heavy Hadrons 5) Tetraquarks
X(3872), etc.

Hadrocharmonium

$J/\psi, \psi(2S), \dots$ even Υ ? affinity for light hadronic matter



$$Z_1(4050)^+ \rightarrow \pi^+ \chi_{c1}(1P) \quad 82^{+51}_{-55}$$

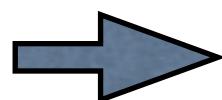
$$Z_2(4250)^+ \rightarrow \pi^+ \chi_{c1}(1P) \quad 177^{+321}_{-72}$$

$$Y(4260) \rightarrow \pi\pi J/\psi \quad 108 \pm 14$$

$$Y(4360) \rightarrow \pi^+ \pi^- \psi(2S) \quad 96 \pm 42$$

$$Z(4430)^+ \rightarrow \pi^+ \psi(2S) \quad 107^{+113}_{-71}$$

$$Y(4660) \rightarrow \pi^+ \pi^- \psi(2S) \quad 48 \pm 15$$



look for J/ψ with baryons; b analogs

Y(4260): BaBar/Belle/Cleo; no D's

- Charmonium Hybrid -- gluonic excitations
Lattice; heavy quark symmetry; NRQCD
Potential yields tower of excitations¹
- Bound state or molecule (next slide)
- Tetraquark^{2,3} $[cs][\bar{c}\bar{s}]$
- Hadrocharmonium (previous slide)
- QCDSR³ $[c\bar{q}]_1[\bar{c}q]_1$; $(S + V), (P + A)$

¹Horn/Mandula; Hasenfratz/Horgan/Kuti/Richard; Juge/Kuti/
Morningstar; Bali/Pineda; Zhu; Kou/Pene; Close/Page

²Maiani/Riquer/Piccinini/Polosa

³Nielsen/Navarra/Lee

Molecules: do the constituents retain their identify as hadrons? (details in the X(3872) section)

$X(3872)$

$$\overline{D}^0 D^{*0}$$

$X(3915)$

$$\overline{D}^{*0} D^{*0} + D^{*+} D^{*-}$$

BGL

$Y(4140)$

$$D_s^{*+} D_s^{*-}$$

BGL

$Y(4260)$

$$D_0 \overline{D}^*, \psi(2S) f_0(980)$$

AN,TKGO

$Z(4430)^+$

$$D^{*+} \overline{D}_1^0$$

LMNN/BGL

$X(4630)$

$$\psi(2S) f_0(980)$$

GHHM

$Y(4660)$

$$\psi(2S) f_0(980)$$

GHM

BGL=Branz,Gutsche,Lyubovitskij

LMNN=Lee,Miharo,Navarro,Nielsen

TKGO=Torres,Kehmchandari,Gamermann,Oset

AN=Albuquerque,Nielsen

GH(H)M=Guo,(Haidenbauer),Hanhart,Meissner

b Exotics above threshold - Belle

MeV	$Z_b(10610)$	$Z_b(10650)$	$Y_b(10888)$
mass	10608.1 ± 1.7	10653.3 ± 1.5	10888.4 ± 3.0
width	15.5 ± 2.4	14.0 ± 2.8	$30.7^{+8.0}_{-7.7}$
mode	$\Upsilon(nS)\pi^+\pi^-$ $n = 1, 2, 3$ $h_b(mP)\pi^+\pi^-$ $m = 1, 2$	$\Upsilon(nS)\pi^+\pi^-$ $n = 1, 2, 3$ $h_b(mP)\pi^+\pi^-$ $m = 1, 2$	$\Upsilon(nS)\pi^+\pi^-$ $n = 1, 2, 3$

Explanations:

$Y_b(1^{--})$ could be analog of $Y(4260)$

hybrid $b\bar{b}g$

disturbed $\Upsilon(5S)$

$Z_b : I^G = 1^+ \ J^P = 1^+$ This conference : A.Kuzmin

charged : cannot be $\bar{b}b$ alone

tetraquarks: Karliner/Lipkin prediction

$$\Upsilon(nS) \rightarrow \pi^\pm T_{bb}^\mp \rightarrow \Upsilon(mS)\pi^-\pi^+$$

↗
isovector charged tetraquark $\bar{b}\bar{b}\bar{d}u$ $\bar{b}\bar{b}\bar{u}d$

prediction : look for subthreshold $I = 0$ state

Molecular overlap: Bugg

$\overline{B}B^*$ threshold = 10604.3 ± 0.6 MeV

$Z_b(10610)$ analogous to $X(3872)$?

\overline{B}^*B^* threshold = $\overline{B}B^* + 46$ MeV

$Z_b(10650)$ analogous to $X(3915)$?

hypothesis : each is $\bar{b}b + \overline{B}B^{(*)}$

Voloshin: more molecules from heavy quark spin symmetry; also predicts decay ratios

$$1^-(0^+) = \frac{1}{2} 0_{b\bar{b}}^- \times 0_{l\bar{t}}^- - \frac{\sqrt{3}}{2} \left(1_{b\bar{b}}^- \otimes 1_{l\bar{t}}^- \right)_{J=0}$$

$$Z_b = 1^+(1^+) = \frac{1}{\sqrt{2}} \left(0_{b\bar{b}}^- \times 1_{l\bar{t}}^- + 1_{b\bar{b}}^- \otimes 0_{l\bar{t}}^- \right) \quad \searrow$$

$$Z'_b = 1^+(1^+) = \frac{1}{\sqrt{2}} \left(0_{b\bar{b}}^- \times 1_{l\bar{t}}^- - 1_{b\bar{b}}^- \otimes 0_{l\bar{t}}^- \right) \quad \nearrow \Upsilon\pi, h_b\pi, \eta_b\rho$$

$$1^-(0^+) = \frac{\sqrt{3}}{2} 0_{b\bar{b}}^- \times 0_{l\bar{t}}^- + \frac{1}{2} \left(1_{b\bar{b}}^- \otimes 1_{l\bar{t}}^- \right)_{J=0}$$

$$\rightarrow \eta_b\pi, \chi_b\pi, \Upsilon\rho$$

X(3872): molecular state?

$$\frac{1}{\sqrt{2}} \left(D^0 \bar{D}^{0*} + \bar{D}^0 D^{0*} \right)$$

Isospin issue:

$$\frac{\Gamma[\rightarrow J/\psi \pi^+ \pi^- \pi^0]}{\Gamma[\rightarrow J/\psi \pi^+ \pi^-]} = 1.0 \pm 0.4 \pm 0.3 \quad \text{Belle 2005}$$

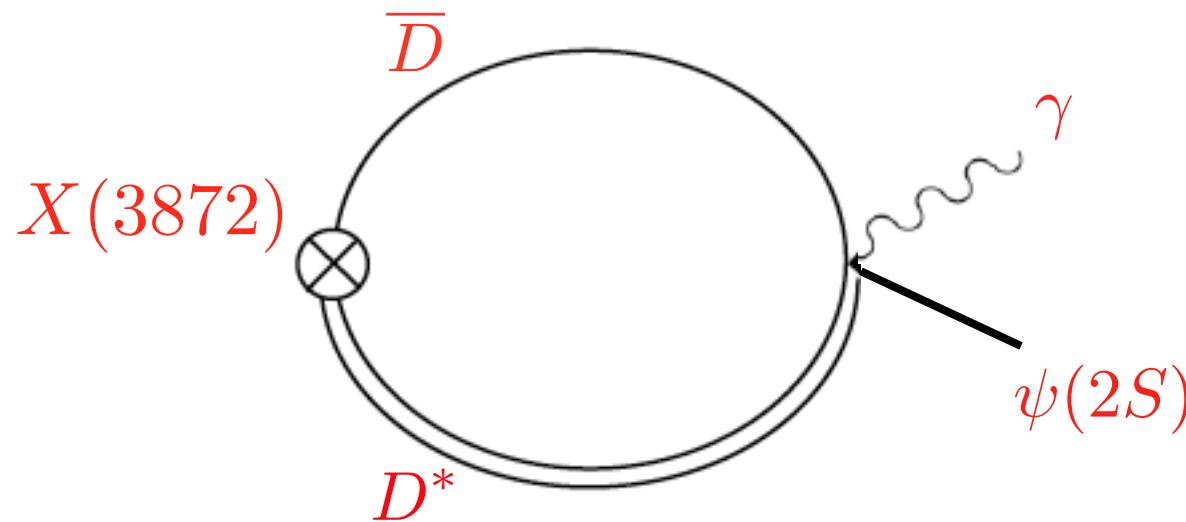
$$\frac{\Gamma[\rightarrow J/\psi \omega]}{\Gamma[\rightarrow J/\psi \pi^+ \pi^-]} = 0.8 \pm 0.3 \quad \text{BaBar 2010}$$

$$J^{PC} = 1^{++} \text{ or } 2^{-+} \quad \text{multipole question}$$

$$m_{D^0 \bar{D}^{0*}} - m_{X(3872)} = 0.42 \pm 0.39 \text{ MeV}$$

S-wave

X-Effective Field Theory: Fleming, Kusunoki, Mehen, van Kolck



Factorization theorems: Braaten/Kusunoki/Lu

$$\text{Rate} = \frac{1}{3} \sum_{\lambda} |\langle 0 | \frac{1}{\sqrt{2}} \epsilon_i(\lambda) (V^i \bar{P} + \bar{V}^i P) | X(3872, \lambda) \rangle|^2 \\ \times (\text{phase space}) \times |\mathcal{C}(\bar{D} D^* \rightarrow f)|^2$$

Universal shallow-bound-state properties from effective range theory: Braaten/Voloshin...

$$\psi_{DD^*}(r) \propto \frac{e^{-\gamma r}}{r} \quad B = \frac{1}{2\mu_{D^*D} a^2} \quad \begin{aligned} \gamma &\sim 20 \text{ MeV} \\ a &\sim 7 \text{ fm} \\ \langle r \rangle &\sim 5 \text{ fm} \end{aligned}$$

Compare Systematic NN treatment: NN-EFT (no pions)

$$A = \mathcal{C} + i\mathcal{C}^2 \frac{Mp}{4\pi} + \mathcal{C}^3 \left(\frac{Mp}{4\pi} \right)^2 + \dots$$

$$A = \frac{4\pi}{M} [-a + ia^2 p + \frac{1}{2}(a^3 - a^2 r_0)p^2 + \dots]$$

$$a^{(1S_0)} \sim -\frac{1}{8 \text{ MeV}}$$

$$a^{(3S_1)} \sim \frac{1}{36 \text{ MeV}}$$

does not converge

Both S-wave scattering lengths anomalously large => momentum expansion fails => reorganize to treat C's nonperturbatively

$$A = -\frac{4\pi}{M} \frac{1}{1/a + ip} + \dots$$

with effective range:

$$A = -\frac{4\pi}{M} \frac{1}{1/a - \frac{1}{2}rp^2 + ip} + \dots$$

EM effects easily included

IF

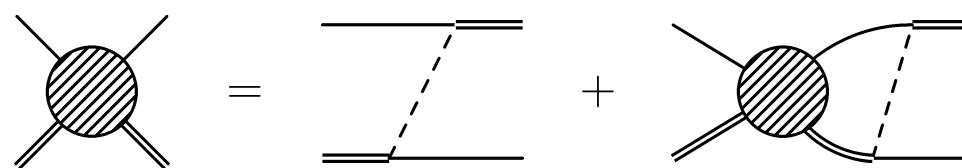
$$X(3872) \sim \frac{1}{\sqrt{2}}(D^0 \bar{D}^{*0} + D^{*0} \bar{D}^0)$$

$$m_X = (3871.55 \pm 0.20) \text{ Mev} \quad E_X = (-0.26 \pm 0.41) \text{ MeV}$$

$$a^{-1} \sim \sqrt{2\mu_X B_X}$$

$$\begin{aligned} \mathcal{L} = & \sum_{j=D^0, D^{*0}, \bar{D}^0, \bar{D}^{*0}} \psi_j^\dagger \left(i\partial_t + \frac{\nabla^2}{2m_j} \right) \psi_j + \Delta X^\dagger X \\ & - \frac{g}{\sqrt{2}} (X^\dagger (\psi_{D^0} \psi_{\bar{D}^{*0}} + \psi_{D^{*0}} \psi_{\bar{D}^0}) + \text{h.c.}) + \dots \end{aligned}$$

Integral equation:

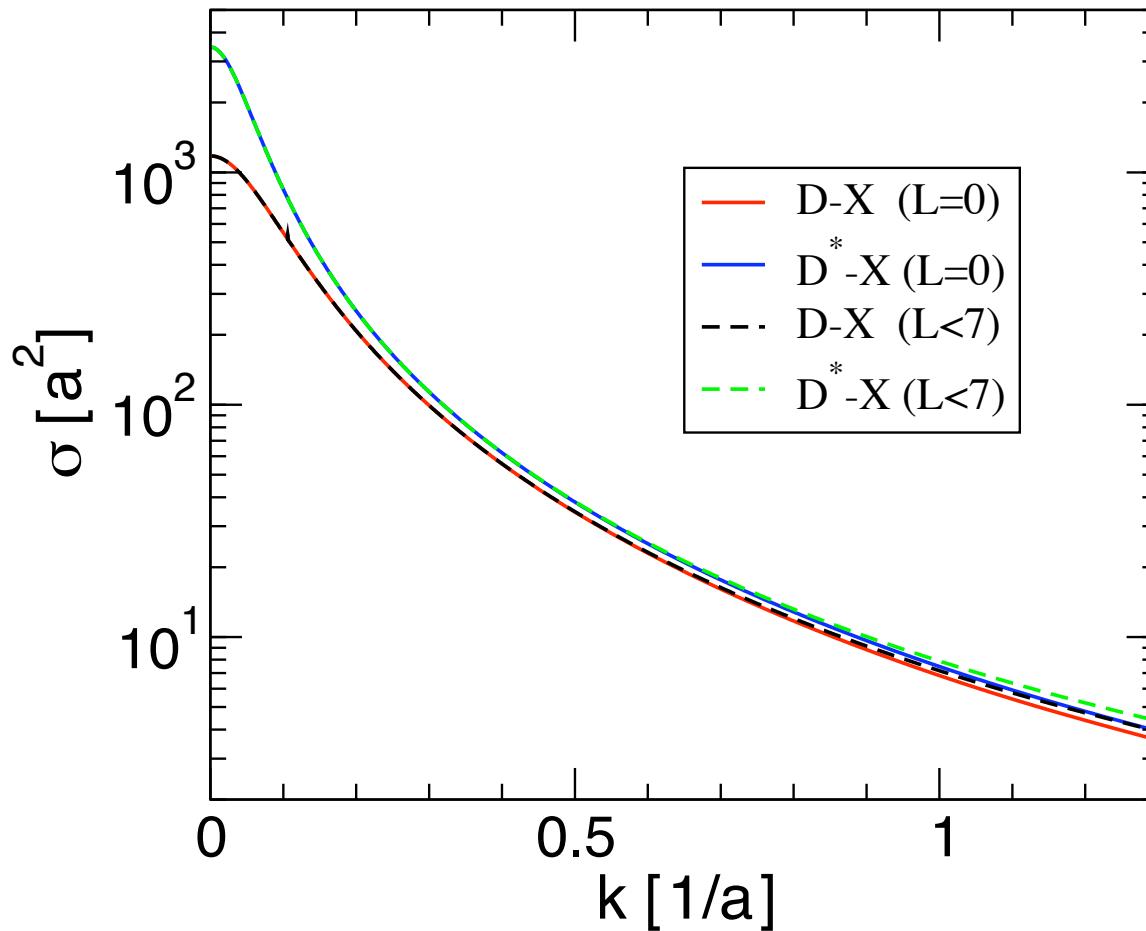


Results depend only on scattering length

$$a_{D^0 X} = -9.7a$$

$$a_{D^{*0} X} = -16.6a$$

Three body cross section vs scattering length



LHC possibilities: $B_c \sim 10^7$ per week

$B\bar{B}$ final state interactions

$$\sigma(b\bar{b}) \sim 0.4 \text{ mb}$$

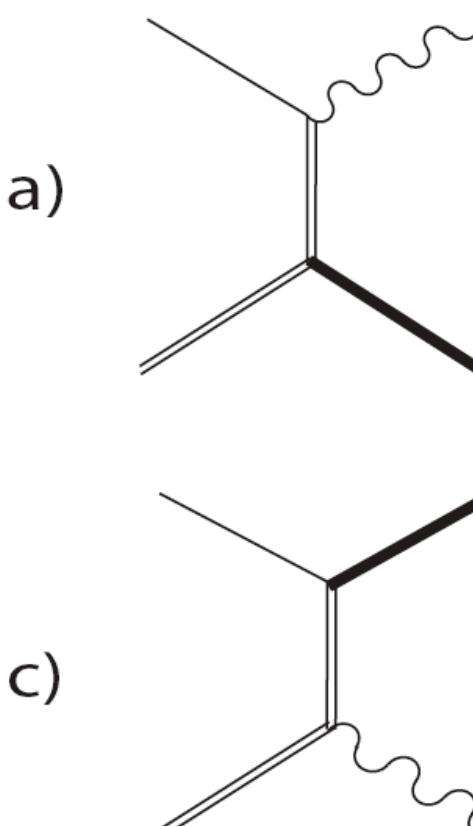
$$\sigma(b\bar{b}b\bar{b}) \sim 5 \text{ fb}$$

$X(3872) \rightarrow \psi(2S)\gamma$
Mehen/RPS

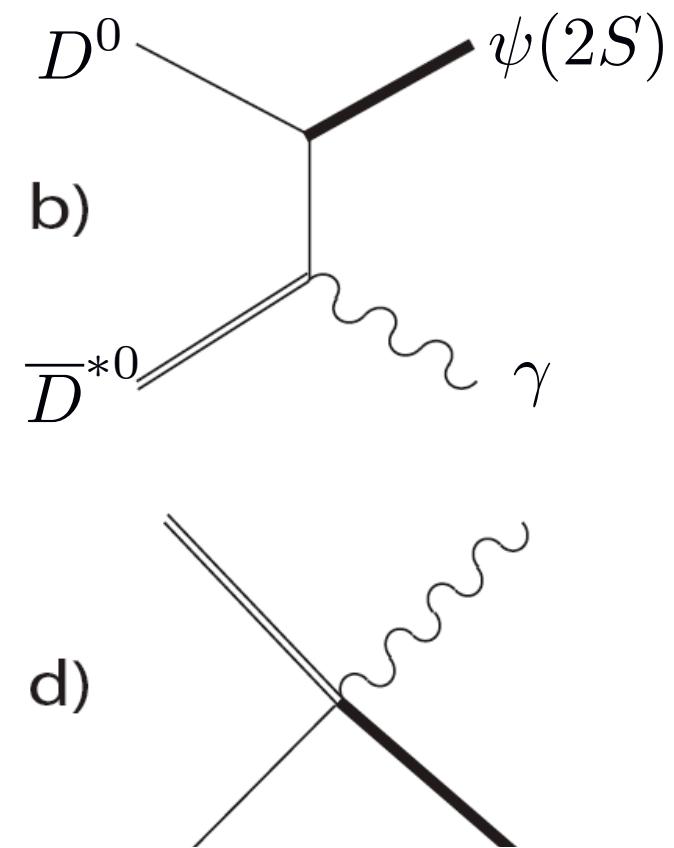
$\beta^{-1} \sim 356$ MeV
Hu, Mehen

$g_2 \sim 0.81$ $\text{GeV}^{-3/2}$
Guo et al., 0907.0521

factorization



XEFT + HBChPT



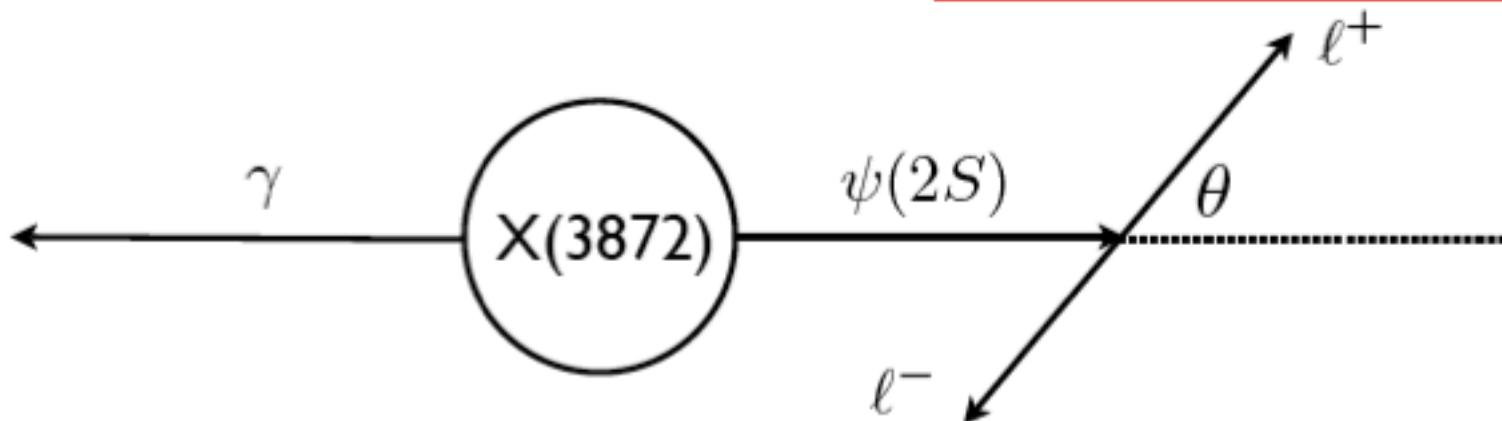
$$\mathcal{L} = \frac{e\beta}{2} \text{Tr}[H_1^\dagger H_1 \vec{\sigma} \cdot \vec{B} Q_{11}] + c.c. + i\frac{g_2}{2} \text{Tr}[J^\dagger H_1 \vec{\sigma} \cdot \vec{\partial} \bar{H}_1] + h.c.$$

$$+ i\frac{ec_1}{2} \text{Tr}[J^\dagger H_1 \vec{\sigma} \cdot \vec{E} \bar{H}_1] + h.c. \quad J = (\eta_c(2S), \psi(2S))$$

$$H_a \sim (D_a, D_a^*); \quad a = 1, 2, 3$$

$$\frac{\Gamma(X(3872) \rightarrow \psi(2S)\gamma)}{\Gamma_{tot}} > 0.03 \text{ (BaBar, PDG)}$$

- **Polarization** $\psi(2S) \rightarrow \ell^+ \ell^-$ $\frac{d\Gamma}{d\cos\theta} \propto 1 + \alpha \cos^2\theta$ $\alpha = \frac{1 - 3f_L}{1 + f_L}$



contact interaction

i) $g_2\beta \ll c_1$ **d) only**

$$f_L = \frac{1}{2}, \alpha = -\frac{1}{3}$$

$$\mathcal{M} \propto \vec{\epsilon}_X \cdot \vec{\epsilon}_\psi^* \times \vec{\epsilon}_\gamma^*$$

constituent decay

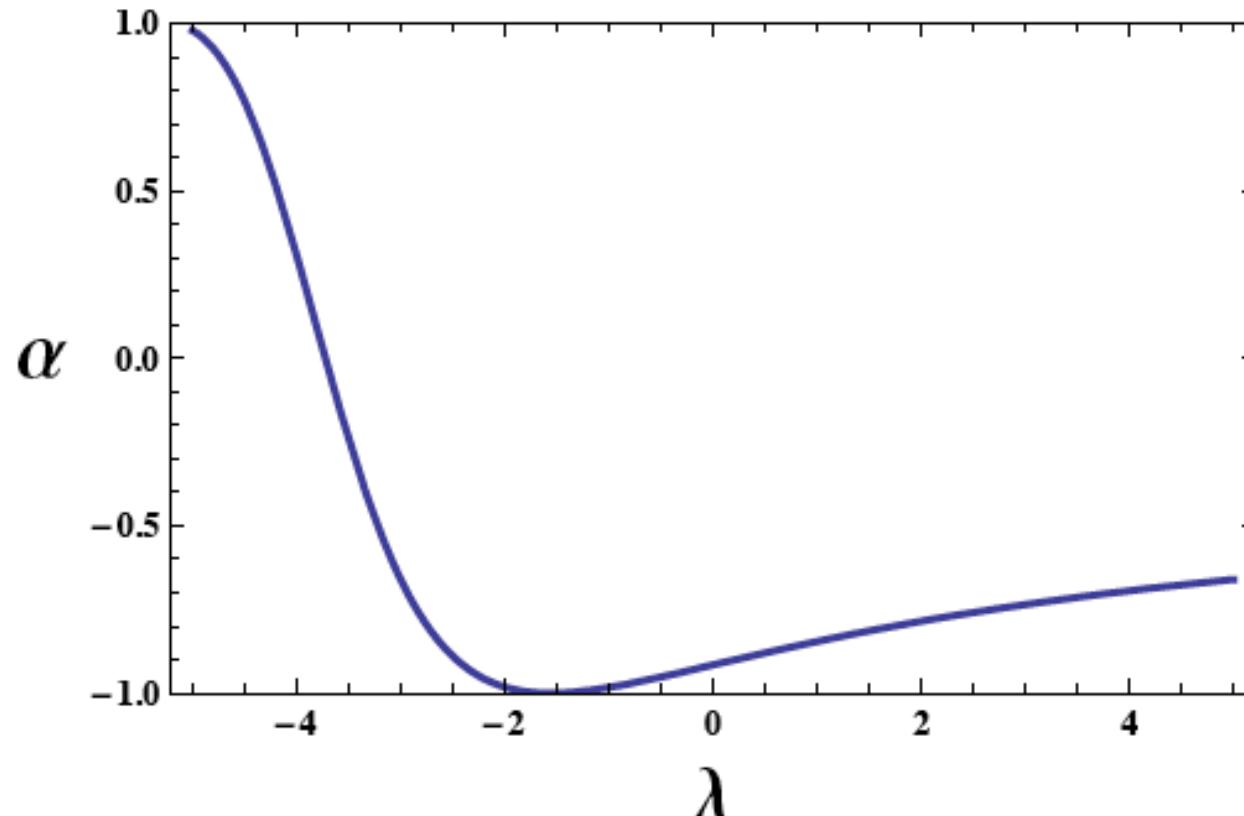
ii) $g_2\beta \gg c_1$ **a-c) only b) dominate**

$$f_L = \frac{4E_\gamma^4}{4E_\gamma^4 + (2E_\gamma + \Delta)^2(E_\gamma - \Delta)^2} = 0.92$$

$$\alpha = -0.91$$

- Polarization measurement would shed light on relative importance of decay mechanisms

- Polarization as function of $\lambda \equiv \frac{3c_1}{g_2\beta} \approx 1.3 \frac{c_1}{\text{GeV}^{-5/2}} \sim O(1)$

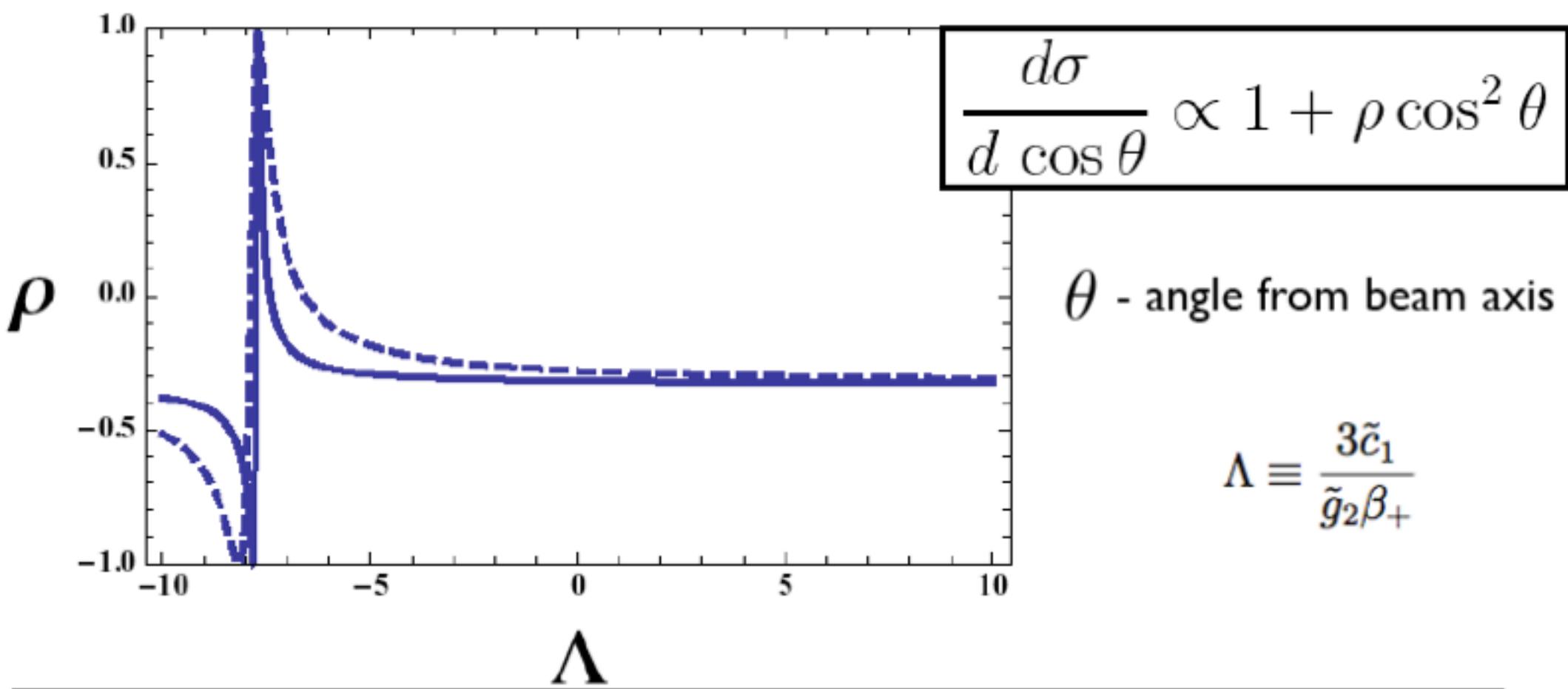


- Longitudinal Polarization ($\alpha < -0.5$) for $-3.5 \leq \lambda \leq 5$
- $X(3872)$ as 2^{-+} : $\alpha = 0.08$

- $e^+e^- \rightarrow \psi(4040) \rightarrow X(3872)\gamma$ (BES?)

$\psi(4040)$ produced with polarization transverse to beam axis (LO)

same (crossed) graphs as $X(3872) \rightarrow \psi(2S)\gamma$



- $J^{PC} = 2^{-+}$ predicts $\rho = 0.08$
molecule predicts $\rho \approx -1/3$ for most of parameter space

Summary

Many new and interesting states in the charmonium “sector”
that we do not understand

The exercise for theorists and experimentalists can only
improve our understanding of QCD/spectra

The new results from the b-side may help clarify things
-- or not

LHC has now seen the X(3872) (Y. Gao)

Better mass determination on D's coming from CLEO --
may clarify “binding” issue of X(3872)

Utilize polarization observables to probe X(3872)
quantum numbers and “wavefunction” questions.

Additional Slides

Initial State Radiation $J^{PC} = 1^{--}$

(MeV)	Y(4660)	Y(4350)	Y(4260)
Mass	4664 ± 12	4361 ± 13	4263 ± 5
Width	48 ± 15	74 ± 18	108 ± 14
Mode	$\pi^+ \pi^- \psi(2S)$	$\pi^+ \pi^- \psi(2S)$	$\pi^+ \pi^- J/\psi$

4360: BaBar 0610057

4260: BaBar 050608, 0808.1543; Cleo
06011021; Belle 0707.2541

4360,4660: Belle 0707.3699; Liu 0805.3560

MeV	$Z(4430)^+$	$Z_2(4250)^+$	$Z_1(4050)^+$
mass	4443^{+24}_{-18}	4248^{+185}_{-45}	4051^{+24}_{-23}
width	107^{+113}_{-71}	177^{+321}_{-72}	82^{+51}_{-55}
mode	$\pi^+ \psi(2S)$	$\pi^+ \chi_{c1}(1P)$	$\pi^+ \chi_{c1}(1P)$

cannot be $c\bar{c}$

quark quantum numbers : $\bar{c}c\bar{d}u$

Hadroncharmonium? Dubynskiy/Voloshin

Charmonium

$c\bar{c}$

$\overline{\overline{\psi(4040)}}$

$\overline{Y(3940)} \ X(3940)$
 $Z(3930)$

$\overline{\overline{\psi(3770)}}$

$X(3872)$

$D\overline{D}(3730)$

$\overline{\psi(2^3S_1)}$
3690

$\eta_c^{'}(2^1S_0)$
3640

$\overline{h_c(1^1P_1)}$
3520

$\chi(1^3P_1)$ $\overline{\chi(1^3P_2)}$

$\overline{\chi(1^3P_0)}$

$\eta_c(1^1S_0)$
2980

$J/\psi(1^3S_0)$
3100

J^{PC}

0^{-+}

1^{+-}

1^{--}

0^{++}

1^{++}

2^{++}

X(3872) Properties (Braaten)

Belle (2003): $B \rightarrow K X(3872)$

$$X(3872) \rightarrow \pi^+ \pi^- J/\psi$$

angular distribution $\Rightarrow 1^{++}$

$$X(3872) \rightarrow J/\psi \gamma \Rightarrow C = +$$

$$M_X = 3871.55 \pm 0.20 \text{ MeV} \quad \Gamma < 2.3 \text{ MeV}$$

$$M(X) - [M(D^{*0}) + M(D^0)] = -0.26 \pm 0.41 \text{ MeV} \Rightarrow \text{S-wave}$$

Threshold resonance universality (Braaten/Hammer) $r \sim 5 \text{ fm}$

$$\frac{\text{Br}(X \rightarrow J/\psi \pi^+ \pi^- \pi^0)}{\text{Br}(X \rightarrow J/\psi \pi^+ \pi^-)} = 1.0 \pm 0.5$$

Amplitudes

$$\begin{aligned}
 a) &= -\frac{g_2 e \beta}{3} \frac{1}{E_\gamma + \Delta} (\vec{k} \cdot \vec{\epsilon}_\psi^* \vec{\epsilon}_{D^*} \cdot \vec{k} \times \vec{\epsilon}_\gamma^* - \vec{k} \cdot \vec{\epsilon}_{D^*} \vec{\epsilon}_\psi^* \cdot \vec{k} \times \vec{\epsilon}_\gamma^*) \\
 b) &= \frac{g_2 e \beta}{3} \frac{1}{\Delta - E_\gamma} \vec{k} \cdot \vec{\epsilon}_\psi^* \vec{\epsilon}_{D^*} \cdot \vec{k} \times \vec{\epsilon}_\gamma^* \\
 c) &= \frac{g_2 e \beta}{3} \frac{1}{E_\gamma} \vec{k} \cdot \vec{\epsilon}_{D^*} \vec{\epsilon}_\psi^* \cdot \vec{k} \times \vec{\epsilon}_\gamma^* \\
 d) &= -e c_1 E_\gamma \vec{\epsilon}_{D^*} \cdot \vec{\epsilon}_\psi^* \times \vec{\epsilon}_\gamma^*
 \end{aligned}$$

$$|\mathcal{M}|^2 = g_2^2 \beta^2 F_1(\Delta, E_\gamma) + g_2 \beta c_1 F_2(\Delta, E_\gamma) + c_1^2 F_3(E_\gamma)$$

$$\begin{aligned}
 |\mathcal{M}(\vec{\epsilon}_\psi)|^2 &= (2g_2^2 \beta^2 A^2 E_\gamma^4 + 4g_2 \beta c_1 A C E_\gamma^2 + 2c_1^2 C^2) |\hat{k} \cdot \vec{\epsilon}_\psi|^2 \\
 &\quad + (g_2^2 \beta^2 B^2 E_\gamma^4 - 2g_2 \beta c_1 B C E_\gamma^2 + c_1^2 C^2) |\hat{k} \times \vec{\epsilon}_\psi|^2
 \end{aligned}$$

$$\Delta \sim 142 \text{ MeV}; \quad E_\gamma \sim 181 \text{ MeV}$$

$$\psi(4040) \rightarrow X(3872)\gamma$$

$$E_\gamma \sim 165 \text{ MeV}$$

$$g_2 \rightarrow \tilde{g}_2; \quad c_1 \rightarrow \tilde{c}_1$$

$$(g'_2)^2 < 0.63 \text{ GeV}^{-3} \quad \text{from width of } \psi(4040)$$

$$\Gamma \text{ to this channel} \sim 10^{-5}$$

by using scattering length to get matrix element