

The role of final-state interactions in Dalitz plot studies

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The role of final-state interactions in Dalitz plot studies

Introduction

- Dalitz plots and CP violation
- the usefulness of hadronic input

What do (low-energy) hadron physicists have on offer?

- scattering consistent with analyticity and unitarity: Roy equations
- decays linked to scattering: form factors and Omnès solution
- low-energy constraints: amplitudes consistent with chiral symmetry (only mentioned in passing)
- many-particle dynamics for the example of $\eta \rightarrow 3\pi$

Les Nabis group

input from C. Hanhart gratefully acknowledged

CP violation in three-body decays

Advantage of 3-body decays:

- resonance-rich environment
- larger branching fractions

here: $B^\pm \rightarrow K^\pm \pi^\mp \pi^\pm$

e.g. 3.7σ signal in $K\rho$

BELLE 2006, BABAR 2008

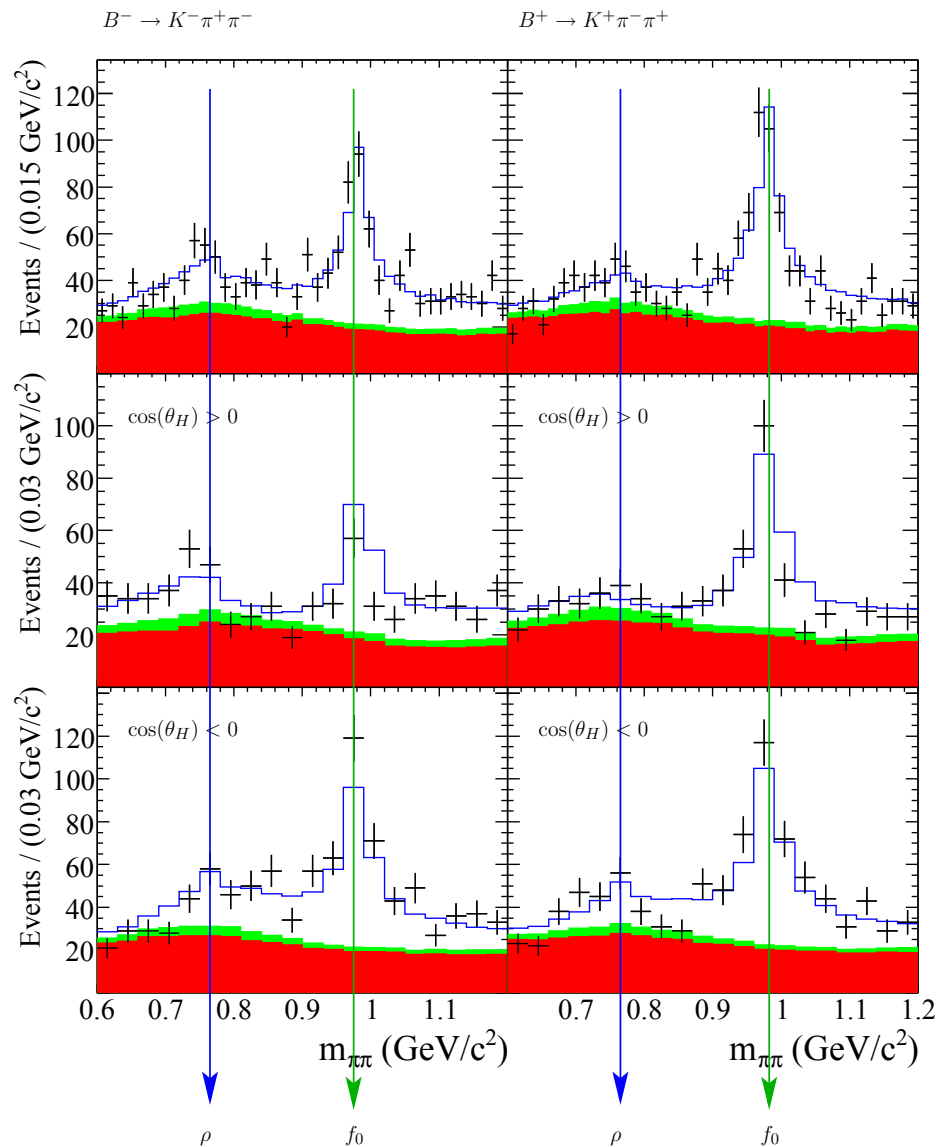
How to analyse CP violation in Dalitz plots?

1. strictly model-independent extraction from data directly

Gardner et al. 2003, 2004

Bediaga et al. 2009

2. theoretical information on strong amplitudes as input!



Direct data analysis: significance

Bediaga et al. 2009

Significance in Dalitz plot distributions:

$$D_{\text{p}} S_{\text{CP}}(i) \doteq \frac{N(i) - \bar{N}(i)}{\sqrt{N(i) + \bar{N}(i)}}$$

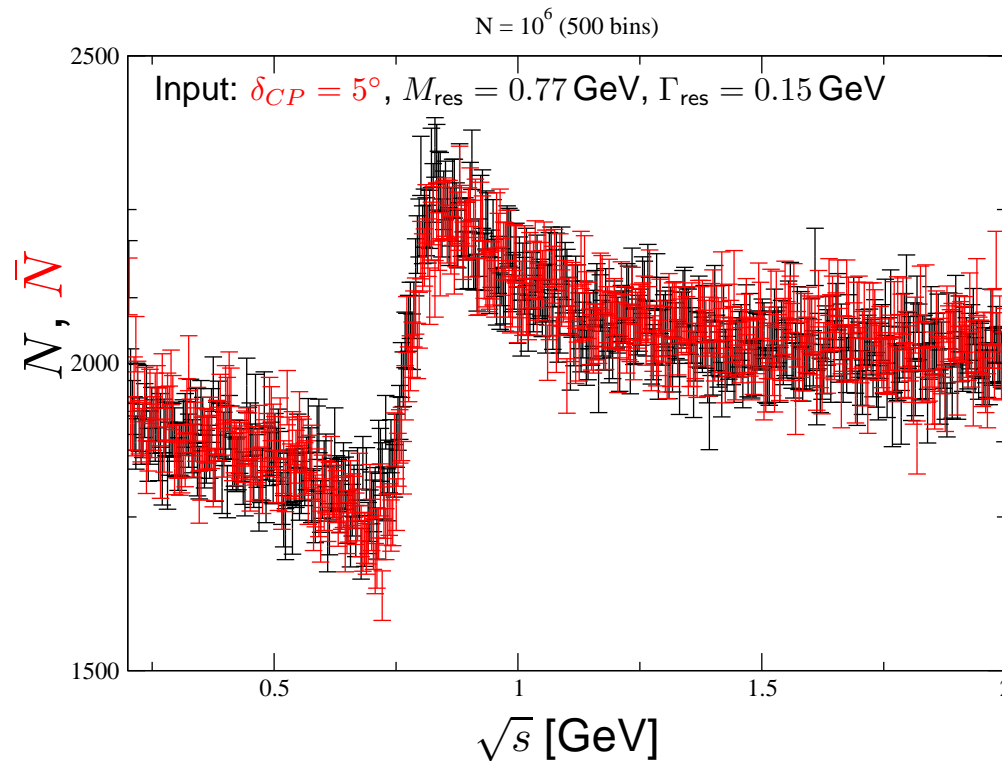
N, \bar{N} : CP-conjugate decays; i : label of a specific Dalitz plot bin

- allows to study **local** asymmetries
- no theoretical input required at all — strictly model-independent
- **B decays**: clear evidence, in particular in $B^{\pm} \rightarrow K^{\pm} \rho^0 (\rightarrow \pi^{\pm} \pi^{\mp})$
consistent with Standard Model **BELLE 2006, BABAR 2008**
- **D decays**: only upper limits (at few-percent level)
Standard Model prediction tiny **BABAR 2008**

Illustration: the use of hadronic amplitudes (1)

- model: resonance plus CP-violating **phase** provided by C. Hanhart

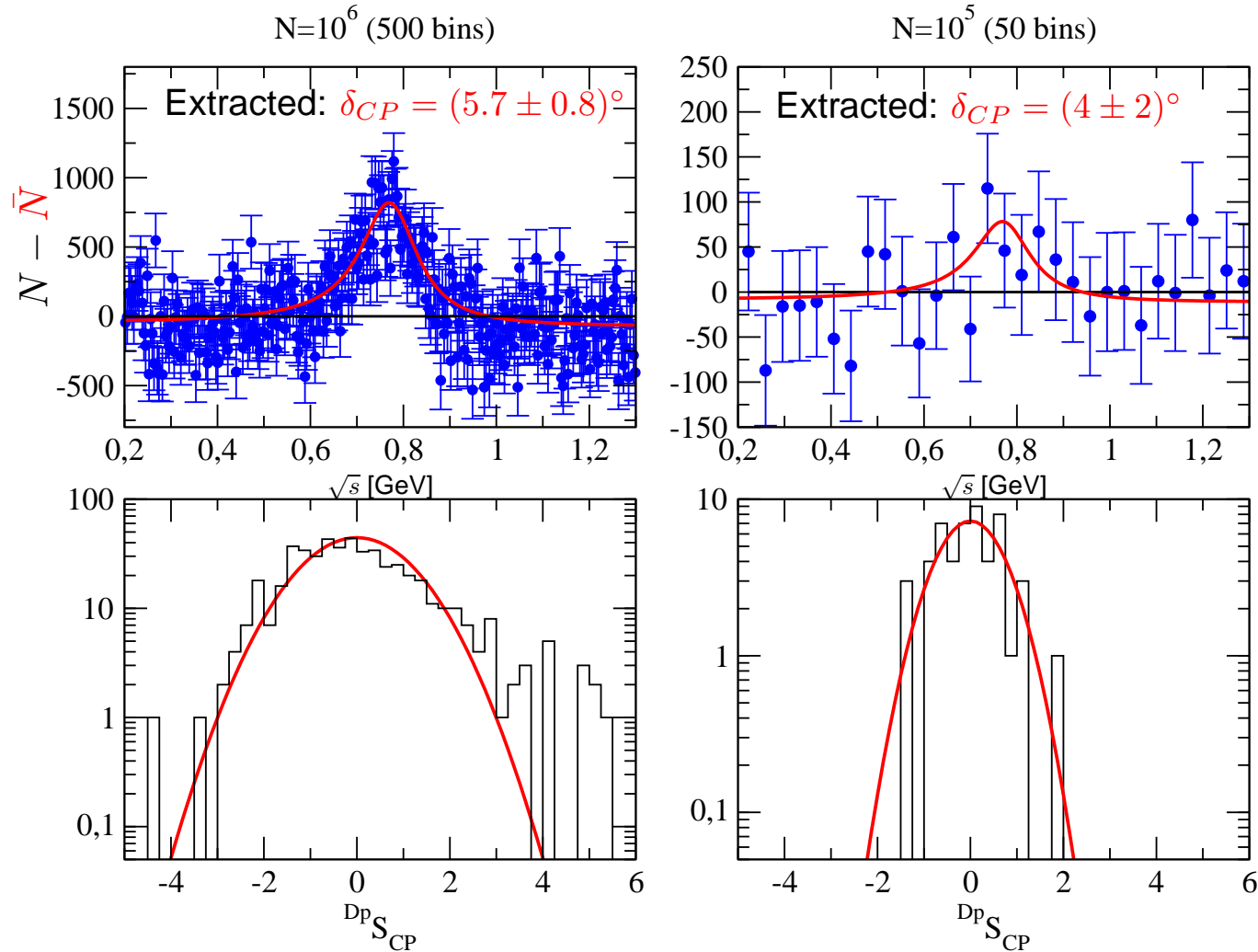
$$N, \bar{N} = \alpha + \beta \operatorname{Re} \left\{ \frac{\exp(\pm i\delta_{CP})}{s - M_{\text{res}}^2 + iM_{\text{res}}\Gamma_{\text{res}}} \right\}$$



- asymmetry: $N - \bar{N} = \sin \delta_{CP} \times \frac{2\beta M_{\text{res}}\Gamma_{\text{res}}}{(s - M_{\text{res}}^2)^2 + (M_{\text{res}}\Gamma_{\text{res}})^2}$

Illustration: the use of hadronic amplitudes (2)

Input: $\delta_{CP} = 5^\circ$, $M_{res} = 0.77 \text{ GeV}$, $\Gamma_{res} = 0.15 \text{ GeV}$



- no signal in significance
hadronic amplitudes still allow to extract phase δ_{CP}

$\pi\pi$ scattering constrained by analyticity and unitarity

compare also talk by M. Hoferichter on Tuesday

Roy equations = coupled system of partial-wave dispersion relations
+ crossing symmetry + unitarity

- twice-subtracted fixed- t dispersion relation:

$$T(s, t) = c(t) + \frac{1}{\pi} \int_{4M_\pi^2}^{\infty} ds' \left\{ \frac{s^2}{s'^2(s' - s)} + \frac{u^2}{s'^2(s' - u)} \right\} \text{Im}T(s', t)$$

- subtraction function $c(t)$ determined from crossing symmetry

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- subtraction function $c(t)$ determined from **crossing symmetry**
- project onto partial waves $t_J^I(s)$ (angular momentum J , isospin I)
 \Rightarrow **coupled system of partial-wave integral equations**

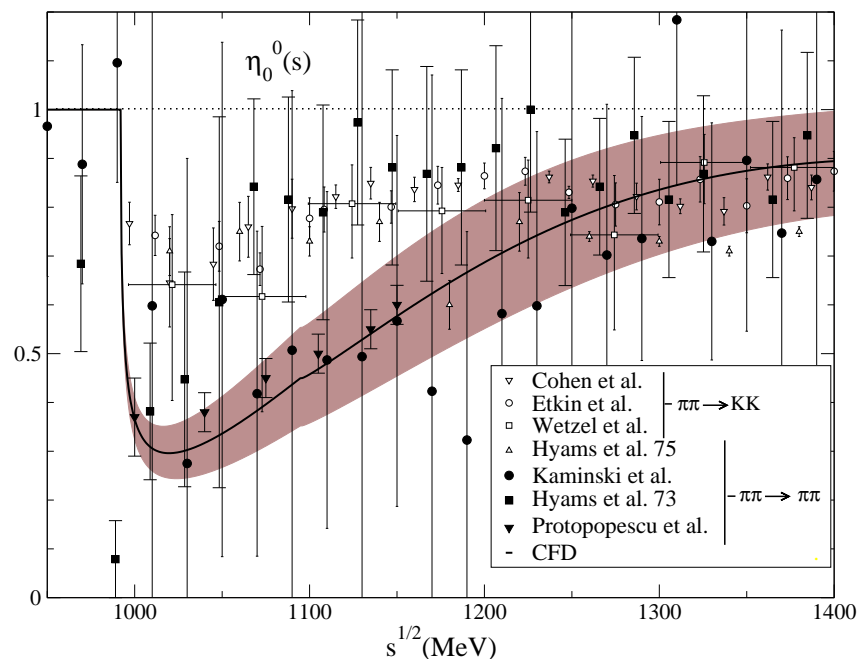
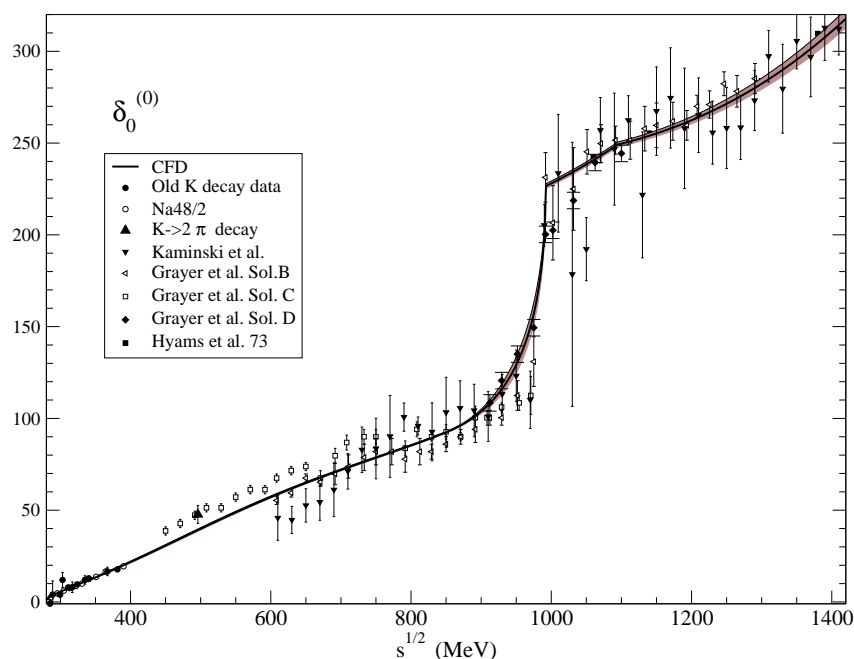
$$t_J^I(s) = k_J^I(s) + \sum_{I'=0}^2 \sum_{J'=0}^{\infty} \int_{4M_\pi^2}^{\infty} ds' K_{JJ'}^{II'}(s, s') \text{Im}t_{J'}^{I'}(s')$$

Roy 1971

- subtraction polynomial $k_J^I(s)$: $\pi\pi$ **scattering lengths**
can be matched to chiral perturbation theory **Colangelo et al. 2001**
- kernel functions $K_{JJ'}^{II'}(s, s')$ known analytically

$\pi\pi$ scattering constrained by analyticity and unitarity

- elastic unitarity \Rightarrow coupled integral equations for **phase shifts**
- modern precision analyses:
 - \triangleright $\pi\pi$ scattering Ananthanarayan et al. 2001, García-Martín et al. 2011
 - \triangleright πK scattering Büttiker et al. 2004
- example: $\pi\pi$ $I = 0$ S-wave phase shift & inelasticity

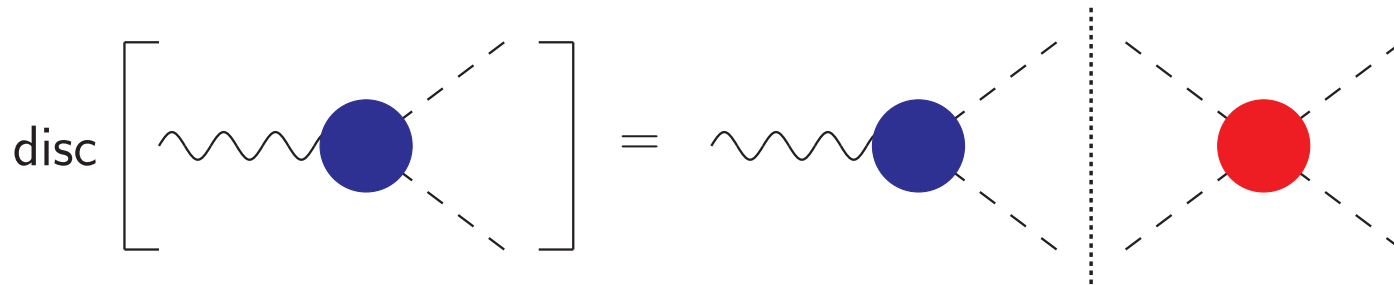


García-Martín et al. 2011

- strong constraints on data from analyticity and unitarity!

Analyticity and unitarity: form factor

- just two particles in final state (form factor); from unitarity:



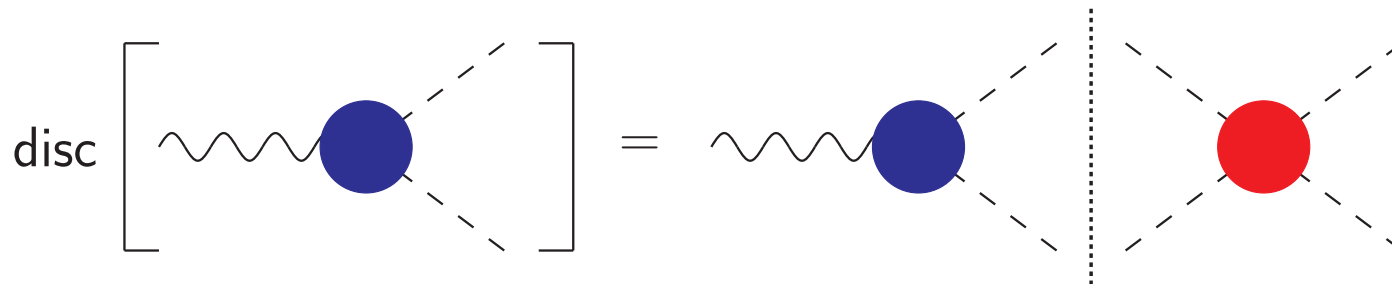
$$\text{disc } F_I(s) = F_I(s) \times \theta(s - 4M_\pi^2) \times \sin \delta_I(s) e^{i\delta_I(s)}$$

\Rightarrow **Watson's final-state theorem**

Watson 1954

Analyticity and unitarity: form factor

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⇒ **Watson's final-state theorem**

Watson 1954

- solution to this homogeneous integral equation known:

$$F_I(s) = P_I(s)\Omega_I(s), \quad \Omega_I(s) = \exp \left\{ \frac{s}{\pi} \int_{4M_\pi^2}^{\infty} ds' \frac{\delta_I(s')}{s'(s' - s)} \right\}$$

$P_I(s)$ polynomial, $\Omega_I(s)$ **Omnès function**

Omnès 1958

completely given in terms of phase shift $\delta_I(s)$

Pion vector form factor and a_μ^{hvp}

- more refined representation:

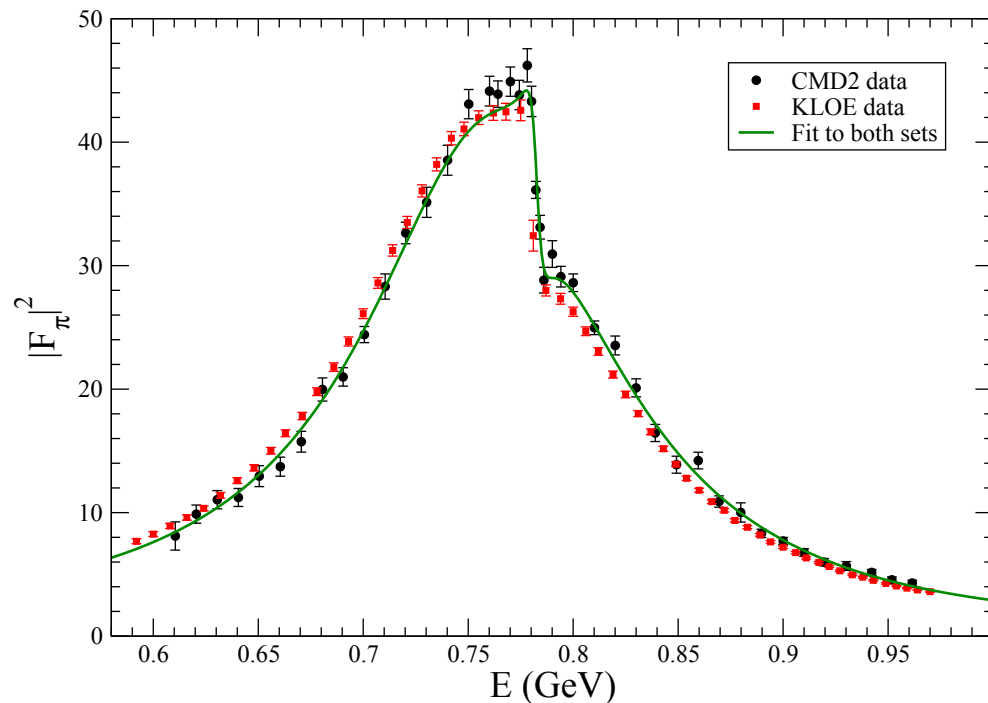
taken from talk by G. Colangelo 2008

$$F_V^\pi(s) = \Omega_1(s)\Omega_{\text{inel}}(s)G_\omega(s)$$

$\Omega_{\text{inel}}(s)$: inelastic for $\sqrt{s} \gtrsim (M_\pi + M_\omega)$, parametrized using conformal mapping techniques

Trocóniz, Ynduráin 2002

$G_\omega(s)$: $\rho - \omega$ mixing



- achieve amazing precision for hadronic contribution to a_μ below 1 GeV:

$$a_\mu^{\text{hvp}}(\sqrt{s} \leq 2M_K) = (493.7 \pm 1.0) \times 10^{-10}$$

Colangelo et al. (preliminary)

- check of data compatibility with analyticity / unitarity

Dispersion relations for three-body decays

compare also following talk by P. Magalhães

Example: $\eta \rightarrow 3\pi$

- interesting due to relation to light quark mass ratios
- $\mathcal{M}(s, t, u) \propto \mathcal{A}(\eta \rightarrow \pi^+ \pi^- \pi^0)$ can be decomposed according to

$$\mathcal{M}(s, t, u) = \mathcal{M}_0(s) + (s-t)\mathcal{M}_1(u) + (s-u)\mathcal{M}_1(t) + \mathcal{M}_2(t) + \mathcal{M}_2(u) - \frac{2}{3}\mathcal{M}_2(s)$$

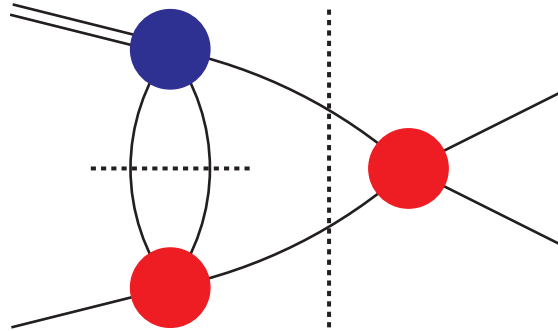
$\mathcal{M}_I(s)$ functions of **one variable** with only a **right-hand cut**

Stern, Sazdjian, Fuchs 1993; Anisovich, Leutwyler 1998

- I : isospin, i.e. $\mathcal{M}_{0,2}$ S-waves, \mathcal{M}_1 P-wave
- decomposition exact if discontinuities in D- and higher partial waves neglected

From unitarity to integral equations: inhomogeneities

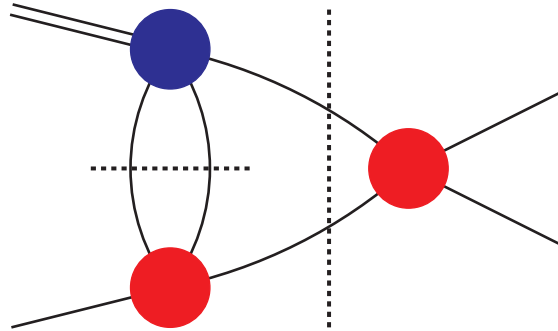
- more complicated unitarity relation for 4-point functions:



$$\text{disc } \mathcal{M}_I(s) = \{ \mathcal{M}_I(s) + \hat{\mathcal{M}}_I(s) \} \times \theta(s - 4 M_\pi^2) \times \sin \delta_I(s) e^{i\delta_I(s)}$$

From unitarity to integral equations: inhomogeneities

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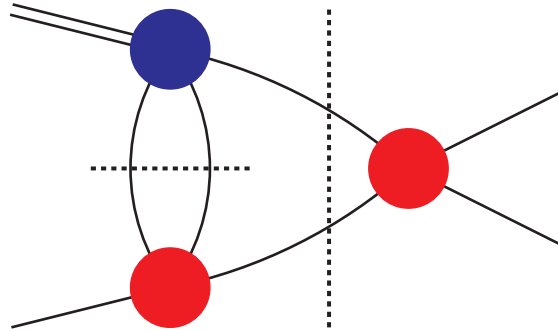
- inhomogeneities** $\hat{\mathcal{M}}_I(s)$: **angular averages** over the $\mathcal{M}_I(s)$: e.g.

$$\hat{\mathcal{M}}_0 = \frac{2}{3} \langle \mathcal{M}_0 \rangle + \frac{20}{9} \langle \mathcal{M}_2 \rangle + 2(s - s_0) \langle \mathcal{M}_1 \rangle + \frac{2}{3} \kappa \langle z \mathcal{M}_1 \rangle$$

$$\langle z^n f \rangle(s) = \frac{1}{2} \int_{-1}^1 dz z^n f(t(s, z))$$

From unitarity to integral equations: inhomogeneities

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$$\langle z^n f \rangle(s) = \frac{1}{2} \int_{-1}^1 dz z^n f(t(s, z))$$

- allows for **cross-channel scattering** between s -, t -, and u -channel
- "angular averaging" non-trivial
 \Rightarrow generates **complex analytic structure** (3-particle cuts)

From unitarity to integral equations: solution

- integral equations including the inhomogeneities $\hat{\mathcal{M}}_I$:

$$\mathcal{M}_0(s) = \Omega_0(s) \left\{ \alpha_0 + \beta_0 s + \gamma_0 s^2 + \frac{s^3}{\pi} \int_{4M_\pi^2}^{\infty} \frac{ds'}{s'^3} \frac{\sin \delta_0(s') \hat{\mathcal{M}}_0(s')}{|\Omega_0(s')|(s' - s - i\epsilon)} \right\}$$

+ 2 similar for $\mathcal{M}_{1,2}(s)$; **4 subtraction constants** to be fixed

Khuri, Treiman 1960; Aitchison 1977; Anisovich, Leutwyler 1998

- solve these equations **iteratively** by a numerical procedure

From unitarity to integral equations: solution

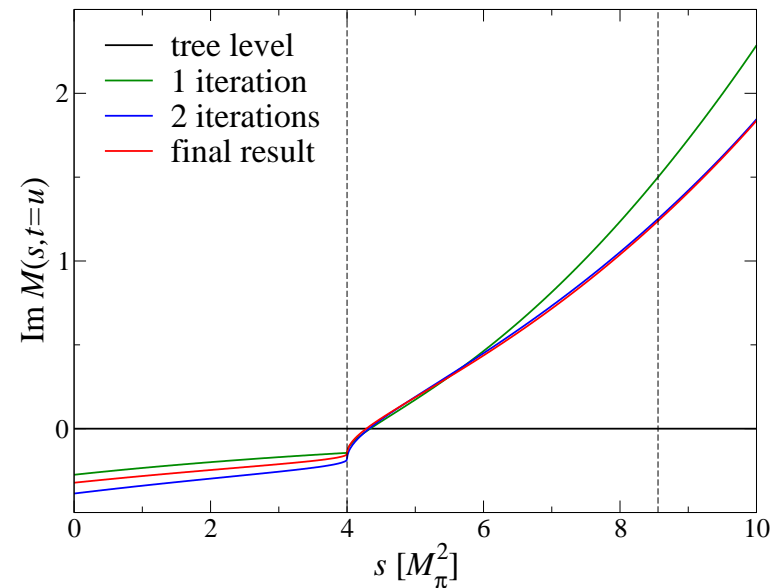
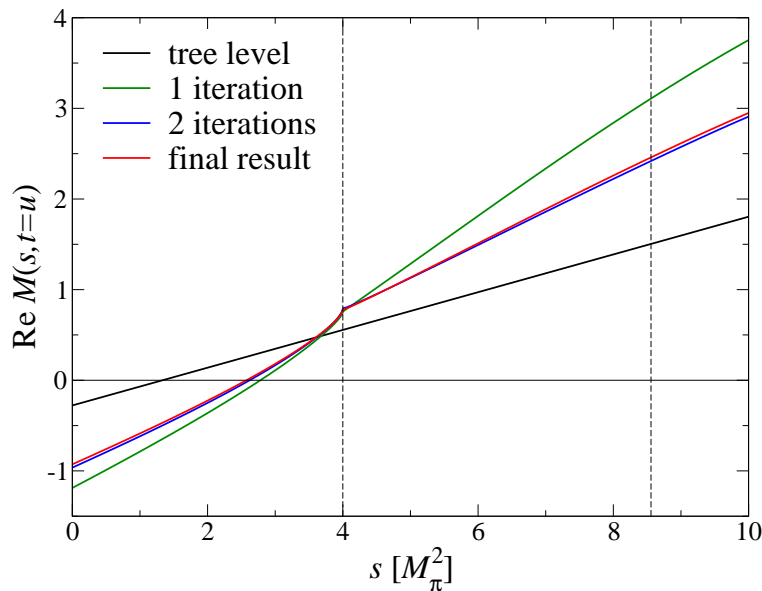
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- solve these equations **iteratively** by a numerical procedure



Schneider, Kubis; compare Colangelo et al. 2010

- fast convergence: close to final result after 2 iterations

Extensions to heavier decays

- currently extended to other decays: $\eta' \rightarrow \eta\pi\pi$, $\eta' \rightarrow 3\pi$, $\omega \rightarrow 3\pi$
Schneider, Nicknig, Kubis
- challenges for going to heavier meson decays:
 - ▷ at higher energies: **coupled-channel** integral equation
 - ▷ **inelasticities** certainly not negligible
 - ▷ perturbative treatment of **crossed-channel effects** reliable?
 - ▷ when are **higher partial waves** non-negligible?

Les Nabis



Paul Serusier, *The talisman* (1888)

Les Nabis

informal network to bring together

particle ("heavy-quark") and

hadron ("light-quark") physicists

from **theory**

I. Bigi, S. Gardner, C. Hanhart, B. Kubis, T. Mannel, U.-G. Meißner, J.R. Peláez, M.R. Pennington...

and **experiment**

I. Bediaga, A.E. Bondar, A. Denig, T.J. Gershon, W. Grandl, B.T. Meadows, K. Peters, U. Wiedner, G. Wilkinson...

to optimize future Dalitz plot CP-studies along these lines!