

and Dalitz plot analysis

Charm 2007 Cornell – August 5 - 8

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Outline

• Introduction:

- Dalitz plot and K-matrix formalism ... the issue

• Analysis:

- Implementation of K-matrix formalism in D-decays
- Examples from FOCUS
 - $\mathbf{D}^+ \rightarrow \pi^+ \pi^- \pi^+$ (2003) 1500 evts
 - $D + \rightarrow K^- \pi^+ \pi^+$ (2007) 53000 evts
- Results and Conclusions
 - What we have learnt so far
 - How we should proceed ...
 - prospects for the future

Dalitz plot: the revenge • SPIRES search for " title Dalitz and date after 1999" 137 entries after 2005 42 entries

• **Experiments:**

FOCUS, E791, CLEO (-c), BaBar-Belle, BES

- From D to B decays
- From decay dynamics to CPV to New Physics $\mathbf{B} \rightarrow \rho \pi$ α angle $\mathbf{B} \rightarrow \mathbf{D}^{(*)}\mathbf{K}^{(*)}$ γ angle

The issue

• to go from $B \rightarrow \pi \pi \pi$ to $B \rightarrow \rho \pi$

means selecting and filtering the desired states among the possible contributions, e.g. $\sigma\pi$, $f_0(980)\pi$, $\sigma\rho$, $\sigma\sigma$, $\rho\pi\pi$...

• a model for **D**⁰ decay is needed

 $-(K\pi)\pi, K(\pi\pi)$

...and a question

- In the era of precision measurements
 - How to deal with the underlying strong dynamics effects?
 - The $\pi\pi$, $K\pi$ S-wave are characterized by broad, overlapping states: **unitarity** is **not** explicitly guaranteed by a simple **sum of Breit -Wigner (BW)** functions
 - Independently of the nature of σ,κ (genuine resonances or strong dynamics structures), they are not simple BW's
 - $f_0(980)$ is a Flatté-like function, coupling to KK and $\pi\pi$

.. a possible answer a *bridge* of knowledge and terminology

• Many problems are already well known in nuclear and intermediate energy physics

K-matrix

– A cultural bridge towards the high energy community

– A common jargon

- An effort has been made in FOCUS to apply it to the Heavy Flavor sector
 - interesting for future B-studies

What is K-matrix? E.P.Wigner, Phys. Rev. 70 (1946) 15

S.U. Chung et al. Ann. Physik 4 (1995) 404

• It follows from S-matrix and, because of S-matrix unitarity, it is real

$$S = I + 2i\rho^{1/2}T\rho^{1/2}$$

$$K^{-1} = T^{-1} + i\rho$$
 $T = (I - iK \cdot \rho)^{-1}K$

- Viceversa, any real K-matrix would generate an unitary S-matrix
- This is the real advantage of the K-matrix approach:
 - It (heavily) simplifies the formalization of any scattering problem since the unitarity of S is automatically respected.

- For a single-pole problem, far away from any threshold, a K-matrix amplitude reduces to the standard BW formula
 - The two descriptions are equivalent
- In all the other cases, the BW representation is no longer valid
 - The most severe problem is that it does not respect unitarity



Adding BWs *a la* "traditional Isobar Model"

- Breaks Unitarity
- Heavily modify the phase motion!

From Scattering to Production

- Thanks to I.J.R. Aitchison (Nucl. Phys. A189 (1972) 514), the K-matrix approach can be extended to production processes
- In technical language,

- From
$$T = (I - iK \cdot \rho)^{-1} K$$

$$F = (I - iK \cdot \rho)^{-1}P$$

• The P-vector describes the coupling at the production with each channel involved in the process

- In our case the production is the D decay

First FOCUS study: D⁺, D_s⁺ $\rightarrow \pi^+\pi^-\pi^+$



$$F = (I - iK \cdot \rho)^{-1} P$$

Describes coupling of resonances to D

Comes from scattering data

Beside restoring the proper dynamical features of the resonances, K-matrix allows for the inclusion of all the knowledge coming from scattering experiments: **enormous amount of results and science!**

$\pi\pi$ S-wave scattering parametrization

"K-matrix analysis of the 00++-wave in the mass region below 1900 MeV" V.V Anisovich and A.V.Sarantsev Eur.Phys.J.A16 (2003) 229

• A global fit to (all) the available data has been performed

*	GAMS	πp→π ⁰ π ⁰ n,ηηn, ηη'n, t <0.2 (GeV/c ²)		
*	GAMS	$\pi p \rightarrow \pi^0 \pi^0 n, 0.30 < t < 1.0 (GeV/c^2)$		
*	BNL	$\pi p^{-} \rightarrow K K n$		
*	CERN-Munich	$\pi^+\pi^- ightarrow \pi^+\pi^-$		
*	Crystal Barrel	$p\overline{p} ightarrow \pi^0 \pi^0 \pi^0$, $\pi^0 \pi^0$ η , π^0 ηη	At rest, from liquid	H_{2}
*	Crystal Barrel	$p\overline{p} \rightarrow \pi^0 \pi^0 \pi^0, \pi^0 \pi^0 \eta$	At rest, from gaseous	H_2
*	Crystal Barrel	$p\overline{p} \rightarrow \pi^+\pi^-\pi^0$, $K^+K^-\pi^0$, $K_sK_s\pi^0$, $K^+K_s\pi^-$	At rest, from liquid	H_{2}
*	Crystal Barrel	$n\overline{p} \rightarrow \pi^0\pi^0\pi^-, \pi^-\pi^-\pi^+, K_sK^-\pi^0, K_sK_s\pi^-$	At rest, from liquid	D_2
*	E852	$\pi p \rightarrow \pi^0 \pi^0 n, 0 < t < 1.5 (GeV/c^2)$		

• It provided the K-matrix input to our three-pion D analysis

$D^+ \rightarrow \pi^+ \pi^- \pi^+ K$ -matrix fit results

PLB 585 (2004) 200



decay channel	fit fractions (%)	phase (deg)
$(S - wave)\pi^+$	$56.00 \pm 3.24 \pm 2.08$	0(fixed)
$f_2(1275)\pi^+$	$11.74 \pm 1.90 \pm 0.23$	-47.5±18.7±11.7
$ ho^0(770)\pi^+$	$30.82 \pm 3.14 \pm 2.29$	-139.4±16.5±9.9

Reasonable fit with <u>no retuning</u> of the A&S K-matrix. No new ingredient (resonance) required not present in the scattering!



$D_s^+ \rightarrow \pi^+ \pi^- \pi^+ K$ -matrix fit results



Yield D⁺ =1527 \pm 51 evts Yield D_s =1475 \pm 50 evts

The high statistics test

- Three-pion analysis suggested:
 - two-body dominance
 - consistency with scattering data
- It was important (mandatory) to **test** the formalism **(a)** high statistics

- the D⁺ \rightarrow K⁻ $\pi^+\pi^+$ channel, i.e. my latest nightmare





e-Print: arXiv:0705.2248 [hep-ex] (to appear in Phys. Lett. B)

The Kπ S-wave scattering parametrization (Mike Pennington)

- two isospin states (I=1/2 and I=3/2) \iff two K-matrices
 - fit S-wave K⁻ $\pi^+ \rightarrow$ K⁻ π^+ LASS data above 825 MeV Nucl. Phys., B 296 (1988) 493

and $K^- \pi^- \rightarrow K^- \pi^-$ scattering from Estabrooks *et al*

Nucl. Phys, B 133 (1978) 490

• extrapolate down to $K\pi$ threshold according to dispersive analysis consistent with ChPT (Buttiker et al, Eur.Phys.J C33 (2004) 409).



		I=	=1/2 K-ma	atrix	
		1 pole	-2 channels	(Κπ -Κη')	
$K_{11} = K_{22} =$	$\left(\frac{s - s_{01/2}}{s_{norm}}\right)$ $\left(\frac{s - s_{01/2}}{s_{norm}}\right)$	$\left(\frac{g_1 \cdot g_1}{s_1 - s} + C_{110}\right) \left(\frac{g_2 \cdot g_2}{s_1 - s} + C_2\right)$	$+C_{111}\tilde{s} + C_{112}\tilde{s}^{2}$ $+C_{221}\tilde{s} + C_{222}\tilde{s}^{2}$	$\mathbf{g_1, g_2}$: real coupling to the first and sec $\mathbf{s_{01/2}} = 0.23 \text{ GeV}^2$ is position in the I=1 scattering amplitud	gs of the s ₁ pole ond channel s the Adler zero 1/2 ChPT elastic de
$K_{12} =$	$\left(\frac{s - s_{01/2}}{s_{norm}}\right)$	$\left(\frac{g_1 \cdot g_2}{s_1 - s} + C_{12}\right)$	$_{0} + C_{121}\widetilde{s} + C_{122}\widetilde{s}^{2}$	$s = m^2 (F)$ $s_{norm} = r$ $\widetilde{s} = s/s$	$K\pi$) $m_{K}^{2}+m_{\pi}^{2}$
Values of	parameters for the I	= 1/2 K -mairix	5	3 5/5	norm 1
Pole (GeV	⁽²)	Coupling (GeV)	C_{1U}	C_{12l}	C22/
z ₁ = 1.79	19	$g_1 = 0.31072$ $g_2 = -0.02323$	Z		
		1000	$C_{110} = 0.79299$	$C_{120} = 0.15040$	$C_{220} = 0.17054$
			$C_{111} = -0.15099$ $C_{112} = 0.00811$	$C_{121} = -0.038266$ $C_{122} = 0.0022596$	$C_{221} = -0.0219$ $C_{222} = 0.00085655$

S-matrix pole : $E = M-i\Gamma/2 = 1.408 - i0.110$ GeV

I=3/2 K-matrix 1 channel scalar function

$$K_{3/2} = \left(\frac{s - s_{03/2}}{s_{norm}}\right) \left(D_{110} + D_{111}\tilde{s} + D_{112}\tilde{s}^{2}\right)$$

 $s_{03/2} = 0.27 \text{ GeV}^2$ is the Adler zero position in the I=3/2 ChPT elastic scattering amplitude

$$s = m^{2}(K\pi)$$

$$s_{norm} = m^{2}_{K} + m^{2}_{\pi}$$

$$\widetilde{s} = s/s_{norm} - 1$$

 $D_{110} = -0.22147$ $D_{111} = 0.026637$ $D_{112} = -0.00092057$

P and F-vectors

• P-vectors

$$(P_{1/2})_{1=K\pi} = \frac{\beta g_1 e^{i\theta}}{s_1 - s} + (c_{10} + c_{11}\hat{s} + c_{12}\hat{s}^2)e^{i\gamma_1}$$

$$(P_{1/2})_{2=K\eta'} = \frac{\beta g_2 e^{i\theta}}{s_1 - s} + (c_{20} + c_{21}\hat{s} + c_{22}\hat{s}^2)e^{i\gamma_2} \qquad \hat{s}_c = 2 \text{ GeV}^2$$

$$P_{3/2} = (c_{30} + c_{31}\hat{s} + c_{32}\hat{s}^2)e^{i\gamma_3}$$

...and F-vectors

$$F_{3/2} = (I - iK_{3/2}\rho)^{-1}P_{3/2}$$
$$(F_{1/2})_{1=K\pi} = (I - iK_{1/2}\rho)_{1j}^{-1}(P_{1/2})_j$$

 β , θ , c_{ij} , γ_i are the free parameters, all the others are fixed to scattering

... finally ready to fit $D^+ \rightarrow K^- \pi^+ \pi^+$ = 202.114



coefficient	phase (deg)			п	
$\beta = 3.389 \pm 0.152 \pm 0.002 \pm 0.068$	$\theta=286\pm4\pm0.3\pm3.0$	BW-like for J>0 states			
$c_{10} = 1.655 \pm 0.156 \pm 0.010 \pm 0.101$	$\gamma_1 = 304 \pm 6 \pm 0.4 \pm 5.8$	component	fit fraction (%)	phase δ_i (deg)	coefficient
$c_{11} = 0.780 \pm 0.096 \pm 0.003 \pm 0.090$		$K^{*}(892)\pi^{+}$	13.61 ± 0.98	0 (fixed)	1 (fixed)
$c_{12} = -0.954 \pm 0.058 \pm 0.0015 \pm 0.025$			\pm 0.01 \pm 0.30	ь	
$c_{20} = 17.182 \pm 1.036 \pm 0.023 \pm 0.362$	$\gamma_2 = 126 \pm 3 \pm 0.1 \pm 1.2$	$K^{*}(1680)\pi^{+}$	1.90 ± 0.63	1 ± 7	0.373 ± 0.067
$c_{30} = 0.734 \pm 0.080 \pm 0.005 \pm 0.030$	$\gamma_3 = 211 \pm 10 \pm 0.7 \pm 7.8$		$\pm~0.009\pm0.43$	$\pm \ 0.1 \pm 6$	\pm 0.009 \pm 0.04
Total S-wave fit fraction = $83.23 \pm 1.50 \pm 0.04 \pm 0.07$ %	S-wave fraction	$K_2^*(1430)\pi^+$	0.39 ± 0.09	296 ± 7	0.169 ± 0.017
Isospin 1/2 fraction $~=~207.25\pm25.45\pm1.81\pm12.23~\%$		11	$\pm~0.004\pm0.05$	$\pm~0.3\pm1$	\pm 0.010 \pm 0.01
Isospin 3/2 fraction $= 40.50 \pm 9.63 \pm 0.55 \pm 3.15$ %	83 ±1.5 %	$K^{*}(1410)\pi^{+}$	0.48 ± 0.21	293 ± 17	0.188 ± 0.041
S-L1- 2			$\pm 0.012 \pm 0.17$	$\pm 0.4 \pm 7$	$\pm 0.002 \pm 0.03$

 ± 0.047

 ± 0.012

 ± 0.030

Comparison with the isobar fit

Adaptive binning and χ^2 contributions for dp to kpp (data)



•serves as the standard for fit quality •requires two "ad hoc" scalars states with free masses and widths (BW) with no reference to how these states appear in other $K\pi$ interactions (an effective data description)

 $\begin{array}{c} \mathbf{m} = 856 \pm 17 \\ k \ \Gamma = 464 \pm 28 \end{array} \quad K_0^* (1430) \quad \begin{array}{c} \mathbf{m} = 1461 \pm 4 \\ \Gamma = 177 \pm 8 \end{array}$

Isobar and K-matrix fits show
same "hot spots" in the adaptive binning scheme
good agreement in vector-tensor fit parameters

What else can we infer from F-vectors?



Phase comparison





Results (I)

- The hypothesis of the **two -body dominance** is **consistent** with the **high statistics** $D^+ \rightarrow K^-\pi^+\pi^+$
- The first determination in D decays of the I=1/2 and I=3/2 for the S-wave $K\pi$ system has been performed
- Our results show close consistency with Kπ scattering data, and consequently, with Watson's theorem predictions for two-body Kπ interactions in the low Kπ mass region where elastic processes dominate.

Results (II)

• Our K-matrix representation fits along the real energy axis inputs on scattering data and ChPT in close agreement with those used by Descotes-Genon and Moussallam (Eur. Phys. J C48 (2006) 553) that locate k with

mass
$$(653 \pm 15) \text{ MeV/c}^2$$

and

width
$$(557 \pm 24)$$
 MeV/c²

different from isobar • Whatever k is revealed by our data, it is the same as that found in scattering data

Results (III)

- Our K-matrix description gives a fit quality globally good.
- However it deteriorates at higher $K\pi$ mass
 - Two channels: $K\pi$ and $K\eta$ ':
 - Reliable info on the former, poor constraints on the latter

• Improvements: using a number of D-decay chains with $K\pi$ final state interactions and inputting all these in one combined analysis in which several inelastic channels are included in the K-matrix formalism.



Conclusions

- Dalitz plot analysis is teaching us much about hadronic decays. It will definitely keep us company over the next few years
- Some complications have already emerged
 - expecially in the charm field

others (unexpected) will only become clearer when we delve deeper into the beauty sector

- **B**_s will be a new chapter (PLB645 (2007) 201: $B_s \rightarrow K\pi\pi, B_s \rightarrow KK\pi$)
- There will be work for both theorists and experimentalists
 Synergy invaluable!

The are no shortcuts toward ambitious and high-precision studies and NP search

Back-up slides

Isobar analysis of D⁺ $\rightarrow \pi^{+}\pi^{+}\pi^{-}\pi^{-}$ would instead require

a new scalar meson: σ $m = 442.6 \pm 27.0 \text{ MeV/c}$ $\Gamma = 340.4 \pm 65.5 \text{ MeV/c}$





$$M = (F_{1/2})_{1}(s) + F_{3/2}(s) + \sum_{n} a_{n}e^{i\delta_{n}}A_{n}$$

for J >0
$$A_{n} = F_{D}F_{r} \times |p_{1}|^{J}|p_{3}|^{J}P_{J}\cos(\theta_{13}^{r}) \times \frac{1}{m_{r}^{2} - m_{12}^{2} - im_{r}\Gamma_{r}}$$

$$\begin{bmatrix} F = 1 \\ F = (1 + R^2 p^2)^{-\frac{1}{2}} \\ F = (9 + 3R^2 p^2 + 3R^4 p^4)^{-\frac{1}{2}} \end{bmatrix} \begin{cases} Spin \ 0 \\ Spin \ 1 \\ Spin \ 2 \end{cases} \begin{cases} P_J = 1 \\ P_J = (-2\overline{p_3} \cdot \overline{p_1}) \\ P_J = 2(p_3 p_1)^2(3\cos^2 \theta_{13} - 1) \end{cases}$$

Isobar fit parameters

Table 2

Fit fractions, phases, and coefficients from the isobar fit to the FOCUS $D^+ \rightarrow K^- \pi^+ \pi^+$ data. The first error is statistic, the second error is systematic from the experiment, and the third error is systematic induced by model input parameters for higher resonances

Channel	Fit fraction (%)	Phase \mathcal{E}_{ℓ} (deg)	Coefficient
non-resonant	29.7 ± 4.5	325 ± 4	1.47 ± 0.11
	$\pm 1.5 \pm 2.1$ (see text)	$\pm 2 \pm 1.2$	$\pm 0.06 \pm 0.06$
$K^{\pi}(892)\pi^{+}$	13.7 ± 0.9	O (fixed)	1 (fixed)
	$\pm 0.6 \pm 0.3$		
$K^*(1410)\pi^+$	0.2 ± 0.1	350± 34	0.12 ± 0.03
	$\pm 0.1 \pm 0.04$	±17±15	$\pm 0.003 \pm 0.01$
$K^*(1680)\pi^+$	1.8 ± 0.4	3 ± 7	0.36 ± 0.04
	$\pm 0.2 \pm 0.3$	±4±8	$\pm 0.02 \pm 0.03$
$K_{2}^{+}(1430)\pi^{+}$	0.4 ± 0.05	319±8	0.17 ± 0.01
weeks weeks	$\pm 0.04 \pm 0.03$	±2±2	$\pm 0.01 \pm 0.01$
$K_n^*(1430)\pi^+$	17.5 ± 1.5	36±5	1.13 ± 0.05
	$\pm 0.8 \pm 0.4$	±2±1.2	$\pm 0.01 \pm 0.02$
кπ+	22.4±3.7	199±6	1.28 ± 0.10
	$\pm 1.2 \pm 1.5$ (see text)	±1±5	$\pm 0.015\pm 0.04$
	Mass (MeV/c^2)	Width (MeV/ c^2)	
$K_0^{+}(1430)$	$1461 \pm 4 \pm 2 \pm 0.5$	177±8±3±1.5	
ĸ	$856 \pm 17 \pm 5 \pm 12$	$464 \pm 28 \pm 6 \pm 21$	