NEW EVIDENCE OF 4Q STRUCTURE IN THE X SYSTEM AD POLOSA INFN ROMA I - LA SAPIENZA

ISOSPIN VIOLATION AND TWO X'S

 $\frac{\mathcal{B}(X \to \pi^+ \pi^- \pi^0 J/\psi)}{\mathcal{B}(X \to \pi^+ \pi^- J/\psi)} = 1.0 \pm 0.4 \pm 0.3$

FROM EARLY OBSERVATIONS BY BELLE AND BABAR (`03-`04)

MOLECULES

4-QUARKS



NO PROBLEM WITH ISOSPIN VIOLATION :: <u>1 STATE</u> :: <u>SMALL</u> DECAY RATE TO DDπ



NEED <u>TWO STATES</u>, AND MAKE ISOSPIN VIOLATION POSSIBLE

FIND THESE TWO X'S IN DATA

A MASS DIFFERENCE X_{U} - X_{D} OF ABOUT ~ 5 MEV WAS PREDICTED :: THEY COULD APPEAR IN B⁺ AND B⁰ SEPARATELY

 $B^+ \to K^+ X_u$ with rate Γ_1 $B^+ \to K^+ X_d$ with rate Γ_2 suppose $\Gamma_1 \gg \Gamma_2 \triangleright \Gamma_4 \gg \Gamma_3$ $B^0 \to K^0 X_u$ with rate Γ_3 $B^0 \to K^0 X_d$ with rate Γ_4



DIFFERENCE IN MASS FROM DATA NOT SIGNIFICATIVE

X(3872): STILL SOME SURPRISES



- 2.5 σ away from the X(3872) world average! $M(J/\psi\pi^+\pi^-) = 3871.2 \pm 0.5$ MeV (World Average)
- If X(3872), J^P = 2⁺ disfavored

V. Poireau DIS 2007 April 2007

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Belle: Phys. Rev. Lett. 97 (2006) 162002 BaBar: preliminary

ARE THERE TWO DIFFERENT X PARTICLES?

:: OUR HYPOTHESIS: TWO X, GENERICALLY PRODUCED IN B⁺⁰ ::

 $X_u \equiv X$ state decaying into $D^0 \bar{D}^0 \pi^0 = X(3876)$ $X_d \equiv X$ state decaying into $J/\psi \pi^+ \pi^- = X(3872)$

:: THE TWO NEUTRAL STATES IN THE 4Q-COMPLEX ::

 $X^{+} = [cu][\bar{c}\bar{d}] \quad X^{-} = [cd][\bar{c}\bar{u}]$ $X_{u} = [cu][\bar{c}\bar{u}] \quad X_{d} = [cd][\bar{c}\bar{d}]$

IT IS TRICKY THAT XD TURNS OUT TO BE LIGHTER THAN XU (MAYBE ELECTROSTATICS IS RESPONSIBLE FOR THIS)

HOW FAR IS THIS PICTURE CONSISTENT WITH A FOUR QUARK MODEL?

HOWEVER, THE ASSUMPTION, THAT XU AND XD WOULD DECAY IN J WITH SIMILAR BRANCHING RATIOS WAS NOT JUSTIFIED AND THE EARLIER SCHEME IS SUPERSEDED BY THE ONE PRESENTED HERE.

A REMARKABLE FACT

 $b+(u)
ightarrowar{c}+car{s}+(u)+qar{q}~~(\Delta I=0)~$ that the inverted

 $\mathcal{A}(B^+ \to K^+ X_u) = \mathbf{V} + \mathbf{S} = \mathcal{A}(B^0 \to K^0 X_d)$ $\mathcal{A}(B^+ \to K^+ X_d) = \mathbf{V} = \mathcal{A}(B^0 \to K^0 X_u)$ $\mathcal{A}(B^+ \to K^0 X^+) = \mathbf{S} = \mathcal{A}(B^0 \to K^+ X^-)$

AS A CONSEQUENCE WE HAVE

$$\left(\frac{B^0}{B^+}\right)_{J/\psi} = \frac{\mathcal{B}(B^0 \to K^0 X_d) \mathcal{B}(X_d \to J/\psi \pi^+ \pi^-)}{\mathcal{B}(B^+ \to K^+ X_d) \mathcal{B}(X_d \to J/\psi \pi^+ \pi^-)} = \frac{\mathcal{B}(B^0 \to K^0 X_d)}{\mathcal{B}(B^+ \to K^+ X_d)} = \frac{\mathcal{B}(B^+ \to K^+ X_d) \mathcal{B}(X_d \to J/\psi \pi^+ \pi^-)}{\mathcal{B}(B^0 \to K^0 X_u)} = \frac{\mathcal{B}(B^+ \to K^+ X_u) \mathcal{B}(X_u \to D\bar{D}\pi)}{\mathcal{B}(B^0 \to K^0 X_u)} = \left[\left(\frac{B^0}{B^+}\right)_{D\bar{D}\pi}\right]^{-1}$$

WHAT DATA TELL (X(3872) AND X(3876) APPEAR TO BE RELATED BY U \Leftrightarrow D SYMMETRY!)

	$f = J/\psi \pi^+ \pi^-$	$f = D^0 \bar{D}^0 \pi^0$
$\overline{\mathcal{B}(B^{\pm} \to K^{\pm}X)\mathcal{B}(X \to f) \times 10^5}$	1.05 ± 0.18	$10.7 \pm 3.1^{1.9}_{3.3}$
	$1.01 \pm 0.25 \pm 0.10$	
$\overline{\mathcal{B}(B^0 \to K^0 X) \mathcal{B}(X \to f) \times 10^5}$		$17.3 \pm 7.0^{3.1}_{5.3}$
	$0.51 \pm 0.28 \pm 0.07$	
$(B^0/B^+)_f$		1.62 ± 0.80
	$0.50 \pm 0.30 \pm 0.05$	$2.23 \pm 0.93 \pm 0.55$

(V)alence and (S)ea needed to build the final state Kaons :: observe that the inverted pattern with BO was already observed in our first paper





scale of the background of multihadronic decays





(RED TWISTS)

DECAYS



QUALITATIVELY WE EXPECT THAT :: (1) MUST BE SMALL (FLAVOR) :: (2) IS LARGER THAN (3)

ALTERNATIVE: TWIST C AND MAKE J/ψ

BY QUARK FLAVOR CONSERVATION X_D SHOULD DECAY IN D^+D^{*-} :: PHASE SPACE FORBIDDEN. D^0D^{*0} is suppressed twice because $UU \leftrightarrow DD$

& BECAUSE OF A SMALL `REDUCED RATE`

WE COULD TWIST HERE C AS WELL; BUT THE *CHEAPEST* ALTERNATIVE IS STILL DD*



A QUALITATIVE PICTURE OF THE BARRIERS

THE BARRIER, QUALITATIVELY FROM DATA:

$$\left(\frac{B^0}{B^+}\right)_{D\bar{D}\pi} = \left|\frac{V}{V+S}\right|^2 \simeq 2 \to \frac{S}{V} \simeq \begin{cases} -0.3\\ -1.7 \end{cases}$$

AND

$$R = \frac{\mathcal{B}(B^+ \to K^+ X_u) \mathcal{B}(X_u \to D^0 \bar{D}^0 \pi^0)}{\mathcal{B}(B^+ \to K^+ X_d) \mathcal{B}(X_d \to J/\psi \pi^+ \pi^-)} \simeq 10$$

and

$$R = \left|\frac{V+S}{V}\right|^2 \frac{\mathcal{B}(X_u \to D^0 \bar{D}^0 \pi^0)}{\mathcal{B}(X_d \to J/\psi \pi^+ \pi^-)}$$

THEREFORE IT FOLLOWS

$$\mathcal{B}(X_d \to J/\psi \pi^+ \pi^-) \simeq \frac{1}{20} \mathcal{B}(X_u \to D^0 \bar{D}^0 \pi^0)$$

VISIBILITY LIMIT (NO SUSPECT OF TWO INTERFERING STRUCTURES)

rosso inventato...

we put the X_u here because we trust that it decays in psi

$$\mathcal{B}(B^+ \to K^+ X_u) \mathcal{B}(X_u \to \psi \pi \pi) < \frac{1}{3} \times \mathcal{B}(B^+ \to K^+ X_d) \mathcal{B}(X_d \to \psi \pi \pi)$$

$$\Rightarrow \qquad \mathcal{B}(X_u \to \psi \pi \pi) \lesssim \frac{1}{30} \times \mathcal{B}(X_u \to D\bar{D}\pi)$$

EFFECTIVE DECAY LAGRANGIAN

 $\mathcal{L}_{eff} = \lambda_{\psi V}^{u} \ \frac{1}{M_{\rho}} \epsilon^{\mu\nu\rho\sigma} (p_{V})_{\mu} V_{\nu} \psi_{\rho} X_{\sigma}^{(u)} + \lambda_{D^{*}D}^{u} \ X^{(u)\mu} (\bar{D}_{\mu}^{*0} D^{0} - \bar{D}^{0} D_{\mu}^{*0}) \simeq$

 $\simeq \boldsymbol{\lambda}_{\psi V}^{\boldsymbol{u}} \mathbf{X}^{(\boldsymbol{u})} \cdot (\mathbf{V} \times \boldsymbol{\psi}) + \boldsymbol{\lambda}_{D^* D}^{\boldsymbol{u}} \mathbf{X}^{(\boldsymbol{u})} \cdot (\mathbf{D}^{*0} D^0 - \mathbf{D}^{*0} \bar{D}^0)$

FROM WHAT DISCUSSED ABOVE WE EXPECT

 $\lambda^u_{\psi V} \ll \lambda^u_{D^*D}$

how the limit in the latter slide transforms in a value for the couplings?

LET US COMPUTE THE REDUCED RATES Y OF THE DECAYS ACCORDING TO:

 $\Gamma(X_{u,d} \to f) = |\lambda|^2 \gamma(f)_{u,d}$

THE REDUCED RATES FOUND FOR THE 3-BODY DECAYS AT HAND ARE

	$\int f = J/\psi \pi^+ \pi^-$	$f = D^0 \bar{D}^0 \pi^0$	$\int f = D^+ D^- \pi^0$	$\int f = J/\psi \pi^+ \pi^0$	f :	
$\overline{\mathbf{X}(3876) = X_u \to f}$	0.59	0.26	$4.5 \cdot 10^{-7}$			N.B. THE O IS
$\mathbf{X}(3872) = X_d \to f$	0.56	0.0102	0			MUCH BROADE
$\overline{X^+(3877)} \to f$				1.2		THAN D*
$X^+(3876) \rightarrow f$				1.2		(FACTOR

 $\Gamma(X_u \to D^0 \bar{D}^0 \pi^0) >> \Gamma(X_u \to J/\psi \pi^+ \pi^-) \simeq$ $\simeq \Gamma(X_d \to J/\psi \pi^+ \pi^-) >> \Gamma(X_d \to D^0 \bar{D}^0 \pi^0)$

COUPLINGS

USING THE RESULT OBTAINED BEFORE

$$\mathcal{B}(X_u \to \psi \pi \pi) \lesssim \frac{1}{30} \times \mathcal{B}(X_u \to D\bar{D}\pi)$$

AND THE RATIO OF REDUCED RATES 0.26/0.59 FROM TABLE, WE GET INDEED

$$\frac{\lambda^u_{\psi V}}{\lambda^u_{D^*D}} \sim 0.13$$

and
$$\Gamma(X_u \to J/\psi \pi^+ \pi^-) \simeq \Gamma(X_d \to J/\psi \pi^+ \pi^-) \le 0.1 \text{ MeV}$$

but we expect X+ just in between of Xu and Xd. u-->d but a repulsion transforms into attraction

PICTURE

- X_u : $\Gamma(\text{multi-g}) \approx 1 \text{ MeV}$, $\Gamma(D^0 D^0 \pi^0) \approx 1-3 \text{ MeV} (B=0.5 \text{ to } 1)$; $B(\psi \pi \pi) = \text{negl.}$
- X_d: Γ (multi-g) \approx 1 MeV, $\Gamma(\psi\pi\pi)\approx$ 0.1 MeV (*B*=0.05); Γ (DD π)= 0.
- X⁺: Γ(multi-g)≈0.1-1 MeV, Γ(ψππ)≈0.2 MeV; Γ(DDπ)= strongly mass dependent, may be dominant for M>3876

THE YET UNOBSERVED X⁺⁻

EXPERIMENTAL BOUNDS

 $\mathcal{B}(B^+ \to K^0 X^+) \mathcal{B}(X^+ \to J/\psi \pi^+ \pi^0) \le 2.2 \times 10^{-5}$ $\mathcal{B}(B^0 \to K^+ X^-) \mathcal{B}(X^- \to J/\psi \pi^- \pi^0) \le 0.54 \times 10^{-5}$

USING PREVIOUS RESULTS WE GET

$$\frac{\mathcal{B}(X^{-} \to \psi \pi^{-} \pi^{0})}{\mathcal{B}(X_{d} \to \psi \pi \pi)} \equiv \frac{\mathcal{B}(B^{0} \to K^{+} X^{-}) \mathcal{B}(X^{-} \to \psi \pi^{-} \pi^{0})}{\mathcal{B}(B^{0} \to K^{+} X^{-}) \mathcal{B}(X_{d} \to \psi \pi \pi)} \leq \\
\leq \frac{0.54 \times 10^{-5}}{\mathcal{B}(B^{0} \to K^{+} X^{-}) \mathcal{B}(X_{d} \to \psi \pi \pi)} \frac{\mathcal{B}(B^{0} \to K^{0} X_{d})}{\mathcal{B}(B^{0} \to K^{0} X_{d})} = \frac{0.54}{0.51} \frac{\mathcal{B}(B^{0} \to K^{0} X_{d})}{\mathcal{B}(B^{0} \to K^{+} X^{-})} \simeq \\
\simeq \left| \frac{V + S}{S} \right|^{2} \times \frac{0.54}{0.51}$$

I.E., THE LIMIT

 $\mathcal{B}(X^+ \to J/\psi \pi^+ \pi^0) \le \left| \frac{V+S}{V} \right|^2 \times \frac{0.54}{0.51} \times \mathcal{B}(X(3872) \to J/\psi \pi^+ \pi^-) \simeq 0.25$

SUMMARY

ARE X(3872) AND X(3876) TWO DIFFERENT PARTICLES? WE GUESS SO AND IDENTIFY THEM AS THE XD AND XU OF THE 4Q MODEL

INDEED THEY CAN EFFECTIVELY BE ACCOMODATED AS THE NEUTRAL COMPONENTS OF A COMPLEX OF FOUR STATES CONTAINING ALSO TWO CHARGED PARTICLES

MAYBE THE CHARGED PARTNERS HAVE TO BE SEARCHED IN OPEN CHARM FINAL STATES: $X^+ \rightarrow D^+ D^- \pi^0$

SEE MAIANI, POLOSA, RIQUER ARXIV:0707.3354

MASSES



BUT...

BELLE AND, VERY RECENTLY, BABAR REPORT A PEAK IN DDπ AT A MASS 3875 ~2.5 MEV AWAY FROM X(3872)! TWO STATES?

MOLECULES? DD π IS OBSERVED TO OCCUR AT A LARGER RATE THAN J ρ

$B_E \ ({\rm MeV})$	$D^0 \overline{D}^0 \pi^0$	$\pi^+\pi^- J/\psi$	$\pi^+\pi^-\pi^0 J/\psi$
0.7	67	1290	720
1.0	66	1215	820
2.0	57	975	1040

By Swanson hep-ph/0311229

WHICH AGREES WITH THE SIMPLE EXPECTATION

 $\Gamma(DD\pi) \sim \Gamma(D^{*0}) = 70 \text{ KeV}$

Y(4260), discovered by BaBar in 2005, J^{PC}=1⁻⁻



$ee \rightarrow J/\psi \pi \pi cross-section$



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