
The Top Pair Threshold: Status & New Developments

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thanks to T. Teubner, A. Manohar, I. Stewart, P. Ruiz-Femenia
C. Reisser, C. Farrell, M. Stahlhofen



Outline

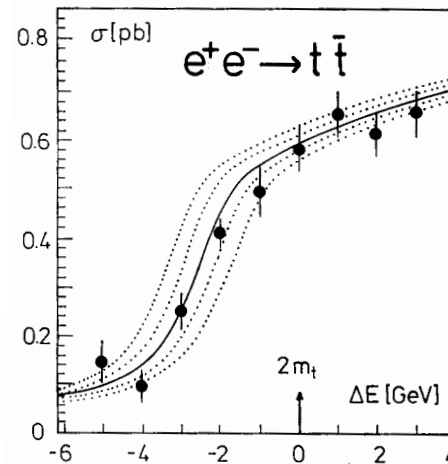
- Brief introductions
 - measurements & requirements
 - theory issues
- Status of QCD corrections
- Electroweak effects
 - hard electroweak corrections
 - finite lifetime effects & phase space matching
- Other Application:
 - $e^+e^- \rightarrow t\bar{t}H$
 - $e^+e^-, \gamma\gamma \rightarrow \tilde{t}\tilde{t}$
 - m_b from $\sigma(e^+e^- \rightarrow \text{hadrons})$



Threshold Measurements

Threshold Scan: $\sqrt{s} \simeq 350 \text{ GeV}$ (Phase I)

- ▷ count number of $t\bar{t}$ events
- ▷ color singlet state
- ▷ background is non-resonant
- ▷ physics well understood
(renormalons, summations)



$$\rightarrow \delta m_t^{\text{exp}} \simeq 50 \text{ MeV}$$

$$\rightarrow \delta m_t^{\text{th}} \simeq 100 \text{ MeV}$$

What mass?

$$\sqrt{s}_{\text{rise}} \sim 2m_t^{\text{thr}} + \text{pert.series}$$

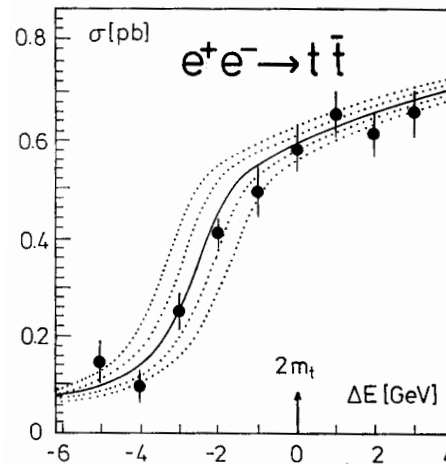
(short distance mass: $1S \leftrightarrow \overline{MS}$)



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(short distance mass: $1S \leftrightarrow \overline{MS}$)

Simulations

$$\mathcal{L} = 300 \text{ fb}^{-1}, 9 + 1 \text{ scan points}$$

Peralta, Martinez, Miquel

$$(\delta m_t)^{\text{stat}} \sim 20 \text{ MeV}$$

$$(\delta \lambda_t / \lambda_t)^{\text{stat}} = 15 - 50\%$$

$$(\delta \alpha_s(M_Z))^{\text{stat}} = 0.001$$

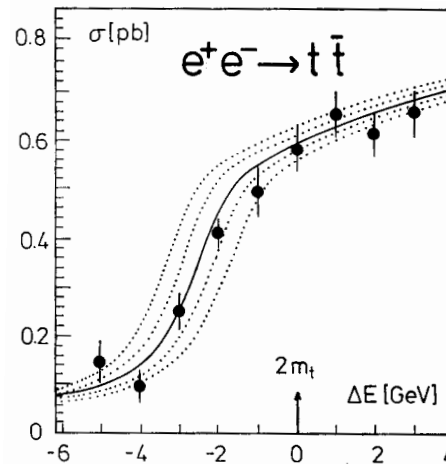
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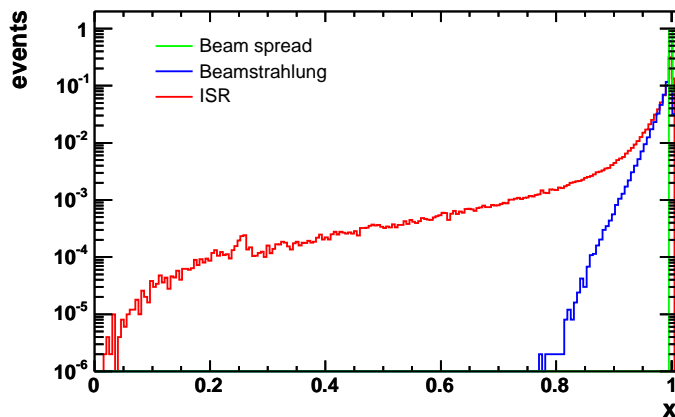
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(short distance mass: $1S \leftrightarrow \overline{MS}$)

Simulations

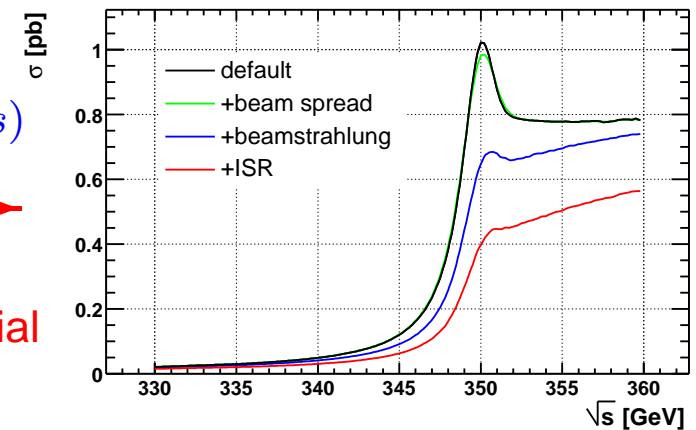
Influence of Luminosity spectrum (LO QED effect!)

Boogert



$$\sigma(s) = \int_0^1 dx L(x) \sigma^0(x^2 s)$$

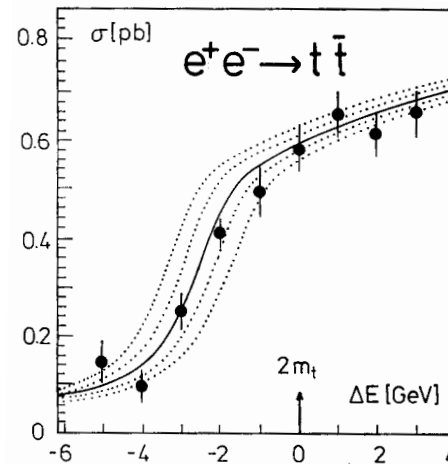
- shape of $\mathcal{L}(x)$ essential
- bhabha acollinearity



Threshold Measurements

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What mass?

$$\sqrt{s}_{\text{rise}} \sim 2m_t^{\text{thr}} + \text{pert.series}$$

(short distance mass: $1S \leftrightarrow \overline{MS}$)

Simulations & theory goal:

$$(\delta\sigma/\sigma)^{\text{theo}} \leq 3\%$$

$$(\delta m_t)^{\text{stat}} \sim 20 \text{ MeV}$$

$$(\delta m_t)^{\text{syst}} \simeq 50 \text{ MeV}$$

$$(\delta m_t)^{\text{theo}} \simeq 100 \text{ MeV}$$

$$(\delta\lambda_t/\lambda_t)^{\text{stat}} = 15 - 50\%$$

$$(\delta\lambda_t/\lambda_t)^{\text{syst}} = ?$$

$$(\delta\lambda_t/\lambda_t)^{\text{theo}} \sim ?$$

$$(\delta\alpha_s(M_Z))^{\text{stat}} = 0.001$$

$$(\delta\alpha_s(M_Z))^{\text{syst}} = 0.002$$

$$(\delta\alpha_s(M_Z))^{\text{theo}} \sim ?$$

$$(\delta\Gamma_t)^{\text{stat}} = 50 \text{ MeV}$$

$$(\delta\Gamma_t)^{\text{syst}} = 15 \text{ MeV}$$

$$(\delta\Gamma_t)^{\text{theo}} \sim ?$$

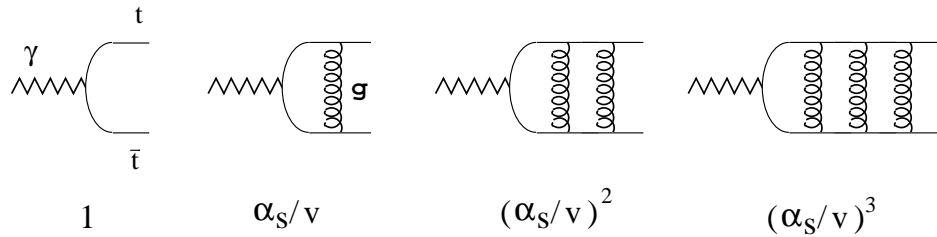


Theory Issues

$$m_t \text{ (hard)} \gg p \sim m_t v \text{ (soft)} \gg E \sim m_t v^2 \text{ (ultrasoft)}$$

- perturbation theory in α_s breaks down

$$(\alpha_s/v)^n$$



“Coulomb singularities”

→ Schrödinger Equation

- perturbation theory in α_s breaks down → large logs $(\alpha_s \ln v)^n$

$$m_t = 175 \text{ GeV}, \quad p \sim 25 \text{ GeV}, \quad E \sim 4 \text{ GeV} \quad \Rightarrow \quad \ln \left(\frac{m_t^2}{E^2} \right) = 8 \quad \rightarrow \text{RGE's}$$

⇒ “multi-scale problem”

$\left\{ \begin{array}{l} \text{vNRQCD: 1-step matching} \\ \text{pNRQCD} \rightarrow \text{A. Siegner's talk} \end{array} \right.$



Theory Issues

- $\Gamma_t \approx 1.5 \text{ GeV} \gg \Lambda_{\text{QCD}} \Rightarrow v = \sqrt{\frac{E}{m}} \rightarrow v_{\text{eff}} = \sqrt{\frac{E+i\Gamma_t}{m}}$
(Fadin,Khoze)

$\Rightarrow m_t \gg p = mv_{\text{eff}} \gg E = mv_{\text{eff}}^2 \gtrsim \Lambda_{\text{QCD}}$ always true !

\Rightarrow top threshold entirely perturbative ! \rightarrow “Schrödinger theory”

- $E \sim \Gamma_t$: top quarks are always produced off-shell !

\rightarrow methods for on-shell production do not apply for

- theoretical computations
- experimental analysis

“theory for unstable particles”

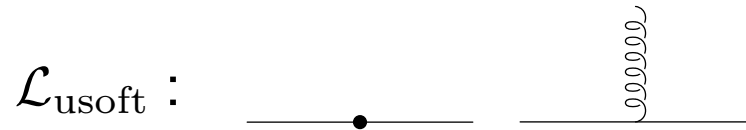


vNRQCD (stable quarks)

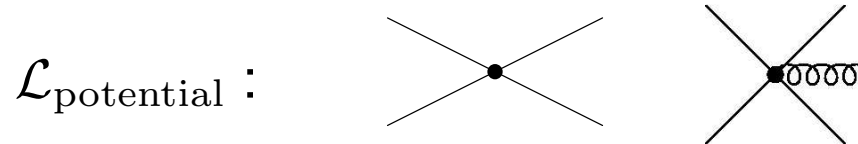
$$\underline{\mathcal{L} = \mathcal{L}_{\text{usoft}} + \mathcal{L}_{\text{potential}} + \mathcal{L}_{\text{soft}}}$$

Luke, Manohar, Rothstein, Stewart, A.H.

$$D^\mu = \partial^\mu + i\mu_U^\epsilon g_s(m\nu^2) A^\mu$$



$$\psi_{\mathbf{p}}^\dagger(x) \left\{ iD^0 - \frac{(\mathbf{p} - i\mathbf{D})^2}{2m_t} - \delta m_t \right\} \psi_{\mathbf{p}}(x)$$



$$\mu_S^{2\epsilon} V(\nu) \psi_{\mathbf{p}'}^\dagger \psi_{\mathbf{p}} \chi_{-\mathbf{p}'}^\dagger \chi_{-\mathbf{p}}$$



$$\mu_S^{2\epsilon} U_{\mu\nu}(\nu) \psi_{\mathbf{p}'}^\dagger A_q^\mu A_{q'}^\nu \psi_{\mathbf{p}}$$

$$\mu_U \propto \mu_S^2/m$$

$$V = \left[\frac{\mathcal{V}_c(\nu)}{\mathbf{k}^2} + \frac{\mathcal{V}_k(\nu)\pi^2}{m|\mathbf{k}|} + \frac{\mathcal{V}_r(\nu)(\mathbf{p}^2 + \mathbf{p}'^2)}{2m^2\mathbf{k}^2} + \frac{\mathcal{V}_2(\nu)}{m^2} + \frac{\mathcal{V}_s(\nu)}{m^2} \mathbf{S}^2 + \frac{\mathcal{V}_\Lambda(\nu)}{m^2} \Lambda + \frac{\mathcal{V}_t(\nu)}{m^2} T \right]$$

$$\mathbf{k} \equiv \mathbf{p} - \mathbf{p}'$$

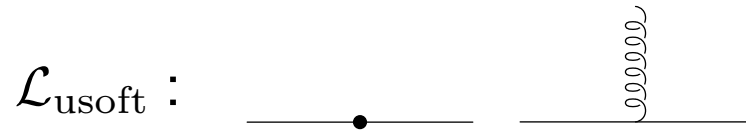


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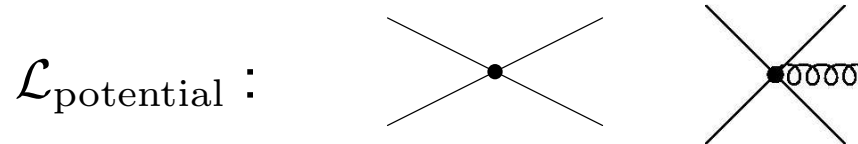
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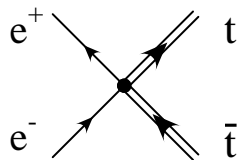
$$\mu_S^{2\epsilon} V(\nu) \psi_{\mathbf{p}'}^\dagger \psi_{\mathbf{p}} \chi_{-\mathbf{p}'}^\dagger \chi_{-\mathbf{p}}$$



$$\mu_S^{2\epsilon} U_{\mu\nu}(\nu) \psi_{\mathbf{p}'}^\dagger A_q^\mu A_{q'}^\nu \psi_{\mathbf{p}}$$

$$\mu_U \propto \mu_S^2/m$$

external currents: (production & annihilation)



$$O_{\mathbf{p}} = C_{V,A}(\nu) \cdot [\bar{e} \gamma^i (\gamma_5) e] [\psi_{\mathbf{p}}^\dagger \sigma^i \tilde{\chi}_{-\mathbf{p}}^*] + \dots \quad t\bar{t} (^3S_1)$$



Cross Section at NNLL Order

Schematic:

$$\begin{aligned}\sigma_{\text{tot}} &\propto \text{Im} \left[\int d^4x e^{-i\hat{q}x} \langle 0 | T \mathbf{O}_{\mathbf{p}}^\dagger(0) \mathbf{O}_{\mathbf{p}'}(x) | 0 \rangle \right] \\ &\propto \text{Im} \left[(C_A(\nu)^2 + C_V(\nu)^2) G(0, 0, \sqrt{s}) \right]\end{aligned}$$

$$\left(-\frac{\nabla^2}{m} - \frac{\nabla^4}{4m^3} + V(\mathbf{r}) - (\sqrt{s} - 2m - 2\delta m) - i\Gamma_t \right) G(\mathbf{r}, \mathbf{r}') = \delta^{(3)}(\mathbf{r} - \mathbf{r}')$$



Cross Section at NNLL Order

Schematic:

$$\sigma_{\text{tot}} \propto \text{Im} \left[\int d^4x e^{-i\hat{q}x} \langle 0 | T \mathbf{O}_{\mathbf{p}}^\dagger(0) \mathbf{O}_{\mathbf{p}'}(x) | 0 \rangle \right]$$

$$\propto \text{Im} \left[(C_A(\nu)^2 + C_V(\nu)^2) G(0, 0, \sqrt{s}) \right]$$

$$\left(-\frac{\nabla^2}{m} - \frac{\nabla^4}{4m^3} + V(\mathbf{r}) - (\sqrt{s} - 2m - 2\delta m) - i\Gamma_t \right) G(\mathbf{r}, \mathbf{r}') = \delta^{(3)}(\mathbf{r} - \mathbf{r}')$$

fully known
at NNLL order ✓

Manohar, Stewart; AH '99-'03
Pineda, Soto '00-'01
Peter '94, Schröder '98

NLL ✓ Luke et al. '99

NNLL (matching) ✓ Benke et al; Czarnecki et al '99

NNLL (non-mixing) ✓ AH '03

NNLL (mixing) **ess. unknown**

spin-dependent (soft) Penin et al. '04

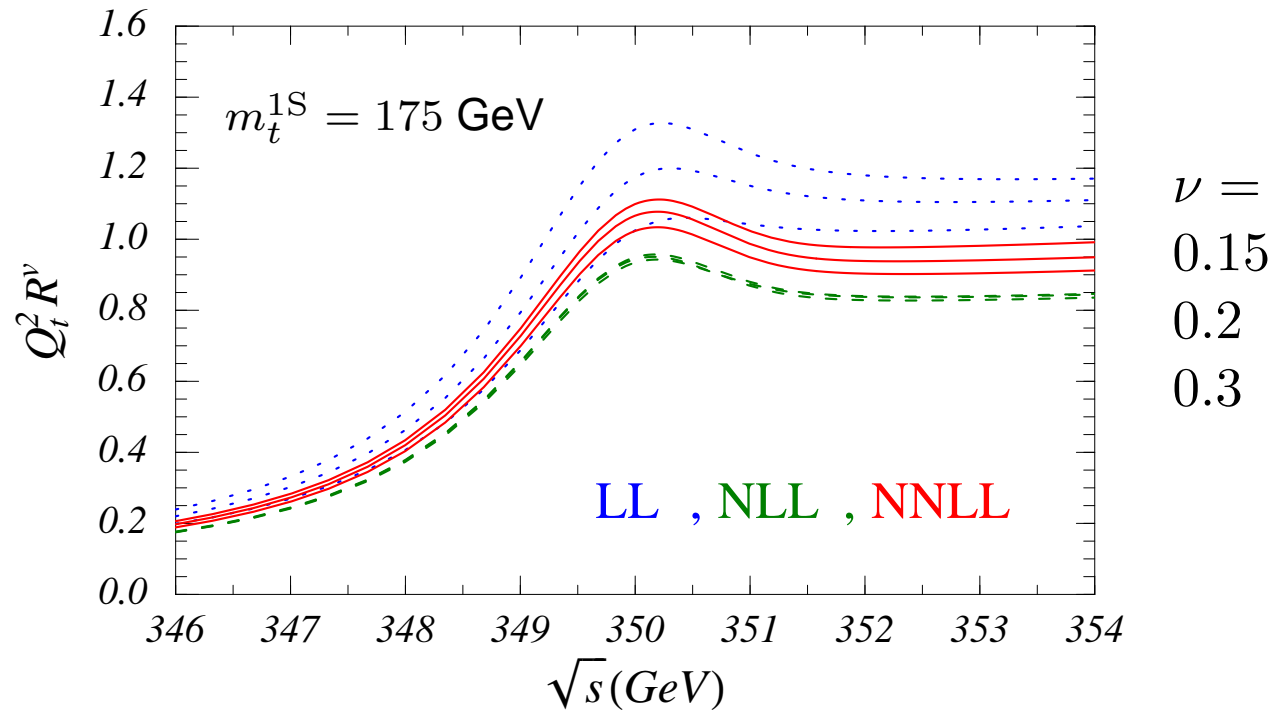
usoft n_f Stahlhofen, AH '05



Cross Section at NNLL Order

1S mass - RG-improved, with NNLL non-mixing terms

Manohar, Stewart, Teubner, AH



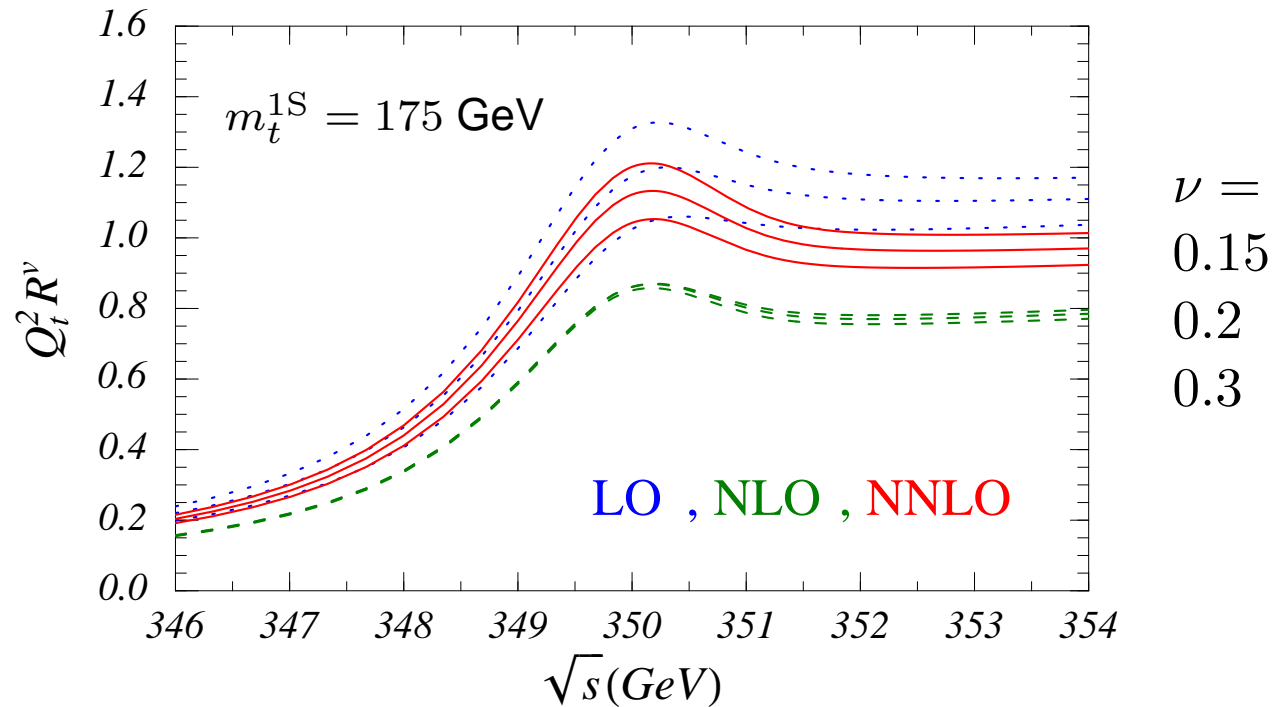
- RGI expansion shows better convergence
 - theory error: $\delta\sigma_{t\bar{t}}/\sigma_{t\bar{t}} \sim \pm 6\%$ goal: 3%
- full NNLL (mixing) running of $C(\nu)$ required → Stahlhofen, AH (usoft w.i.p.)



Cross Section at NNLL Order

1S mass - fixed order approach

Teubner,AH; Melnikov, Yelkovski;Yakovlev;
Beneke,Signer,Smirnov; Sumino, Kiyo



- peak position stable (threshold masses: 1S, PS, ...)
- large sensitivity to factorization/renormalization scale setting
- NNLO partial results: Penin etal. '02 '05, Beneke etal. '05, Eiras etal. '05



NRQCD (unstable quarks)

“inclusive treatment”

- ⇒ Optical Theory: effective complex indices of refraction for absorptive processes
- ⇒ NRQCD: contributions from **real Wb final states** included in EFT matching conditions to QCD+ew. theory (=SM)
- complex matching conditions & anomalous dimensions
 - effective Lagrangian non-hermitian
 - total rates through the **optical theorem**
 - **phase space matching**
 - power counting maintained
 - expansion around mass-shell (Beneke etal. '04)
→ automatic in NRQCD

Christoph Reisser, AH; Phys. Rev. D 71, 074022 (2005)

Christoph Reisser, AH; hep-ph/0604104



Electroweak Effects

3 classes:

- “Hard” electroweak
- Electromagnetic
- Finite lifetime

AHH, hep-ph/0604185

→ No general theory for all cases and observables !

→ EFT for certain observables and given powercounting.

$$\Gamma_t \sim m_t \alpha \sim m_t \alpha_s^2 \Rightarrow v \sim \alpha_s \sim \alpha^{1/2}$$

status for σ_{tot} :

	LL	NLL	NNLL
“Hard” e.w.	✓	✓	✓
El. mag.	✓	(✓)	?
Fin. life.	✓	w.i.p.	w.i.p.

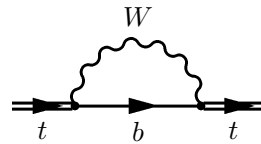


NRQCD (unstable quarks)

quark bilinears: $(\rightarrow \mathcal{L}_{\text{usoft}})$

time dilatation
correction

$$\text{Im}\Sigma_t = \frac{1}{2}\Gamma_t$$



$$= i\Sigma_t \quad \Longrightarrow \quad \delta\mathcal{L} = \psi_{\mathbf{p}}^\dagger \left[i\frac{\Gamma_t}{2} \left(1 - \frac{\mathbf{p}^2}{2m_t^2} \right) \right] \psi_{\mathbf{p}}$$

- power counting: $\Gamma_t \propto m_t g^2 \sim m_t v^2 \Rightarrow g \sim g' \sim v \sim \alpha_s$
- finite lifetime is LL order, $E \rightarrow E + i\Gamma_t$ Fadin,Khoze
- NNLL time dilatation effect
- $E \rightarrow E + i\Gamma_t$ prescription does not work beyond LL order

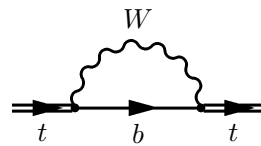


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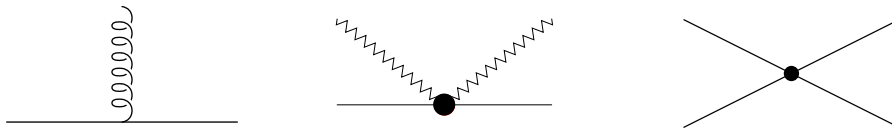


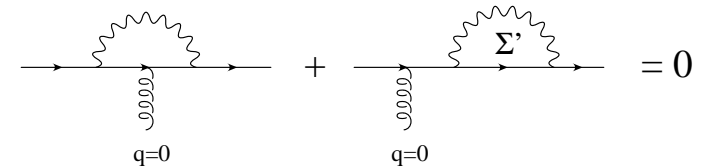
A Feynman diagram showing a quark bilinear with a top quark line and a bottom quark line connected by a W boson loop. The top quark line is labeled 't' and the bottom quark line is labeled 'b'. The W boson loop is labeled 'W'.

$$= i\Sigma_t \quad \Longrightarrow \quad \delta\mathcal{L} = \psi_{\mathbf{p}}^\dagger \left[i\frac{\Gamma_t}{2} \left(1 - \frac{\mathbf{p}^2}{2m_t^2} \right) \right] \psi_{\mathbf{p}}$$


A Feynman diagram showing a quark bilinear with a self-energy correction on the top quark line, represented by a dot on the line.

gluon interactions & potentials:





Two Feynman diagrams showing gluon exchange at $q=0$. The first diagram shows a gluon exchange between two quark lines. The second diagram shows a gluon exchange between two quark lines with a self-energy correction on the top quark line, labeled Σ' . The sum of the two diagrams is equal to zero.

$$= 0$$

- electroweak corrections either beyond NNLL order or vanish due to gauge cancellations

\rightarrow ultrasoft gluon interference effects vanish at NLL [Khoze et al., Melnikov et al.]
and NNLL order (new !)

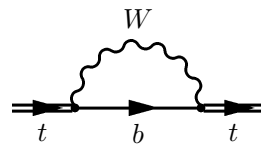



NRQCD (unstable quarks)

quark bilinears: $(\rightarrow \mathcal{L}_{\text{usoft}})$

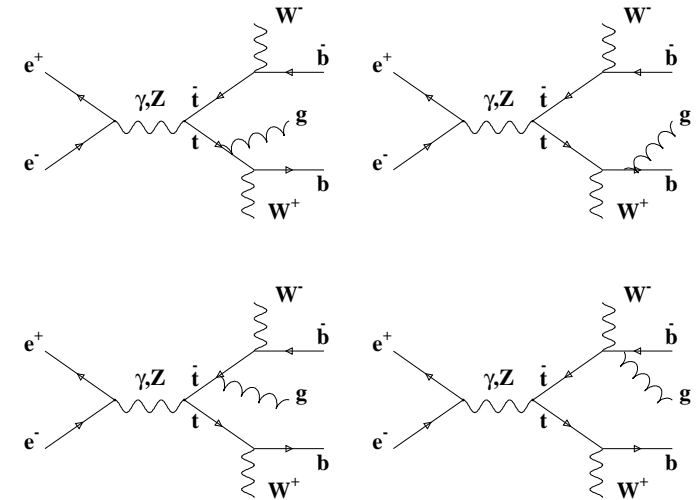
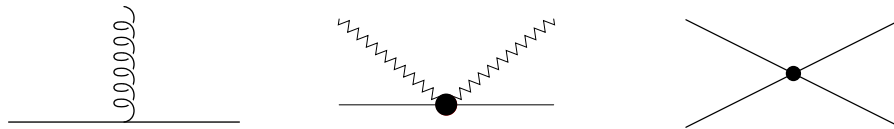
time dilatation correction

$$\text{Im}\Sigma_t = \frac{1}{2}\Gamma_t$$



$$= i\Sigma_t \quad \Longrightarrow \quad \delta\mathcal{L} = \psi_{\mathbf{p}}^\dagger \left[i\frac{\Gamma_t}{2} \left(1 - \frac{\mathbf{p}^2}{2m_t^2} \right) \right] \psi_{\mathbf{p}}$$


gluon interactions & potentials:



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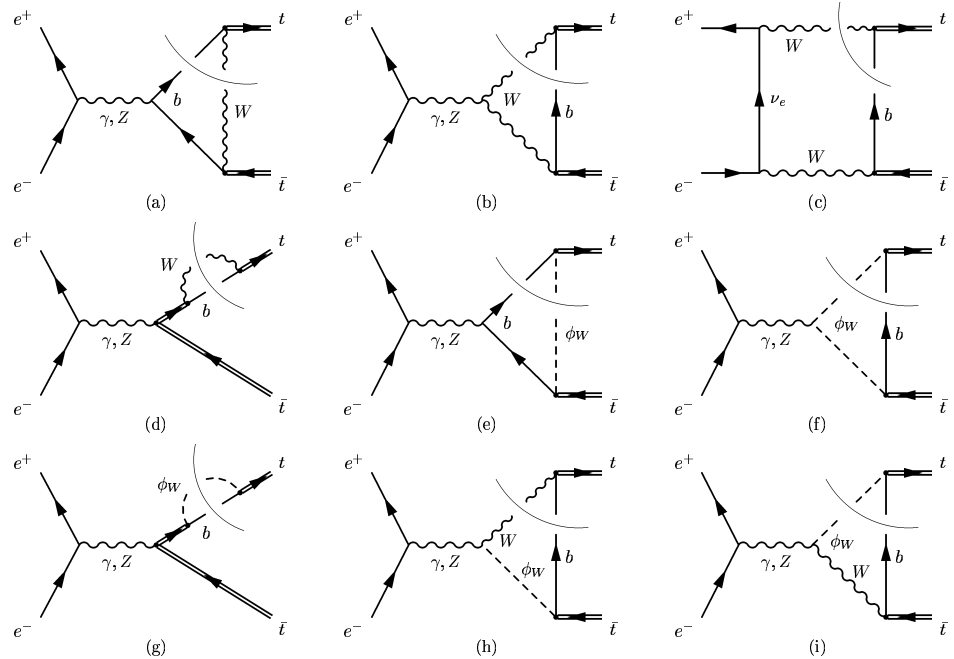
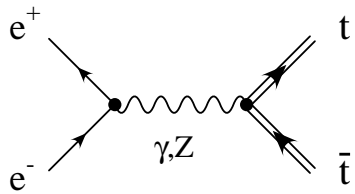
→ ultrasoft gluon interference effects vanish at NLL [Khoze et al., Melnikov et al.] and NNLL order (new !)



NRQCD (unstable quarks)

Currents:

- only bW-cuts included
- bW-cuts gauge invariant
- W treated as stable particle



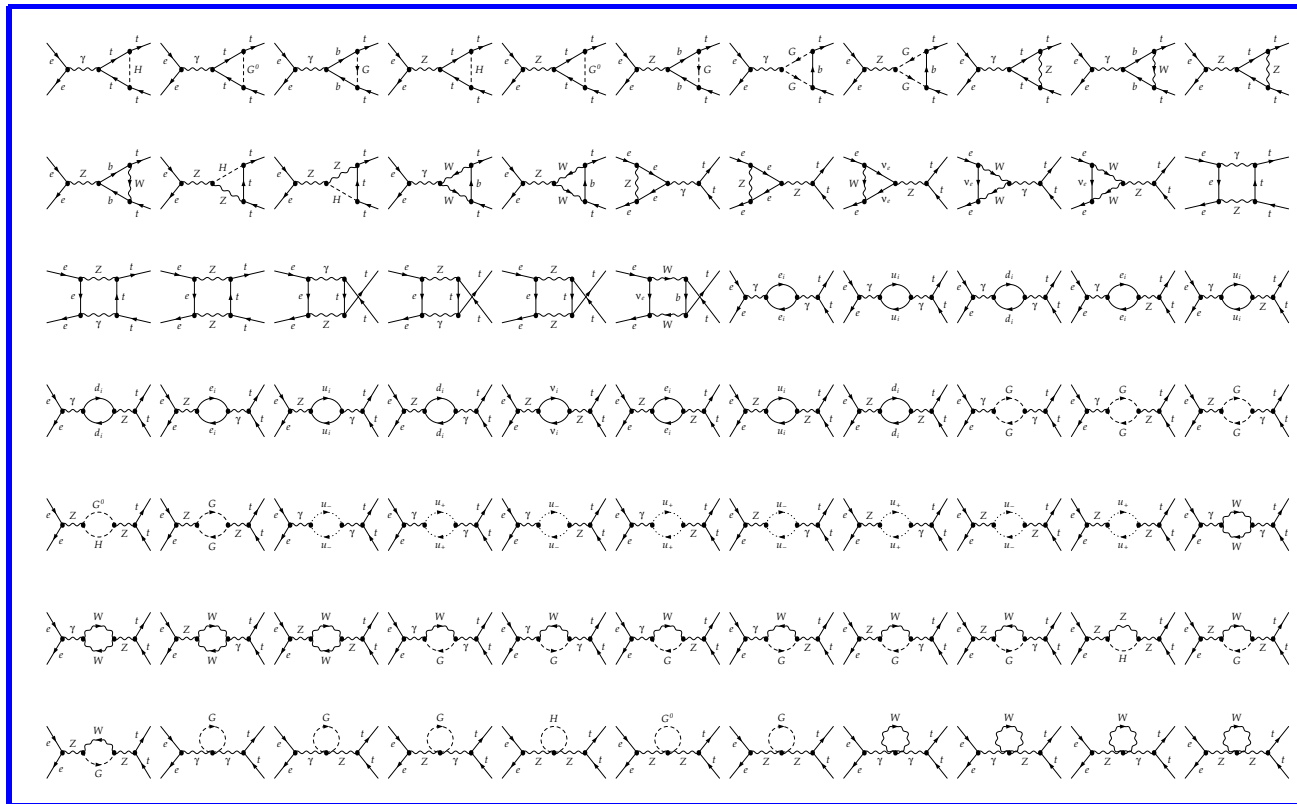
$$O_p = \left[C^{LL} + C^{NLL} + C^{NNLL} + iC_{abs}^{NNLL} \dots \right] \cdot \left(\begin{array}{cc} e^+ & t \\ e^- & \bar{t} \end{array} \right) + \dots$$

$\text{Re}[C_{ew}^{NNLL}] \rightarrow$ weak contributions: C. Reisser, AH, hep-ph/0604104



Hard Electroweak Effects

- “real” short-distance electroweak corrections: $\mathcal{O}(\alpha_{em}) \sim \text{NNLL}$
- only matching conditions to $(e^+e^-)(t\bar{t})$ operators exist !



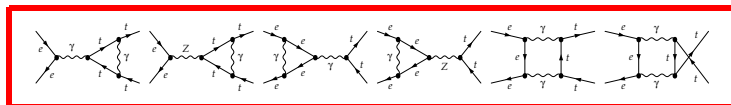
Grzadkowski, Kühn,
etal. (1987)

Guth, Kühn (1992)

updated:

Reisser, AH

hep-ph/0604104

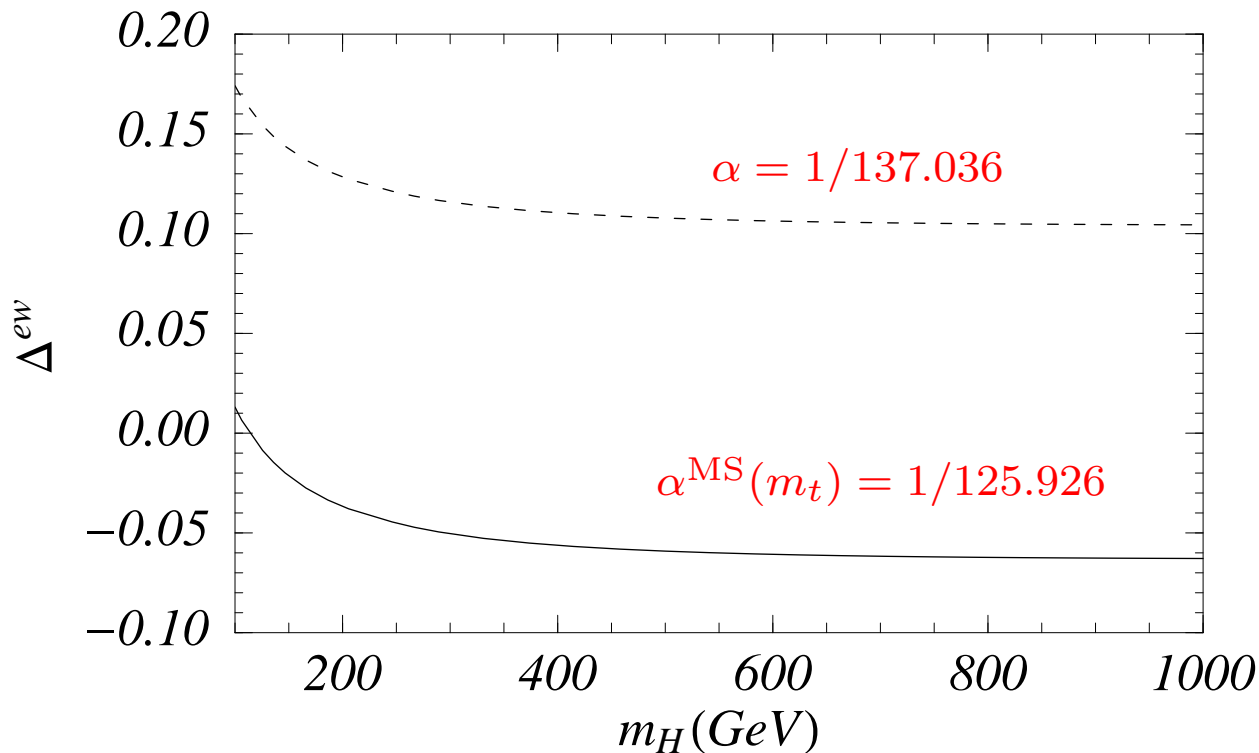


QED corrections (+ real radiation, ISR)
not yet determined → w.i.p.



Hard Electroweak Effects

- “real” short-distance electroweak corrections: $\mathcal{O}(\alpha_{\text{em}}) \sim \text{NNLL}$
- global normalization correction: $\Delta^{\text{ew}} \sim \text{Re}[2c(1)_{\text{ew}}^{\text{NNLL}}]$



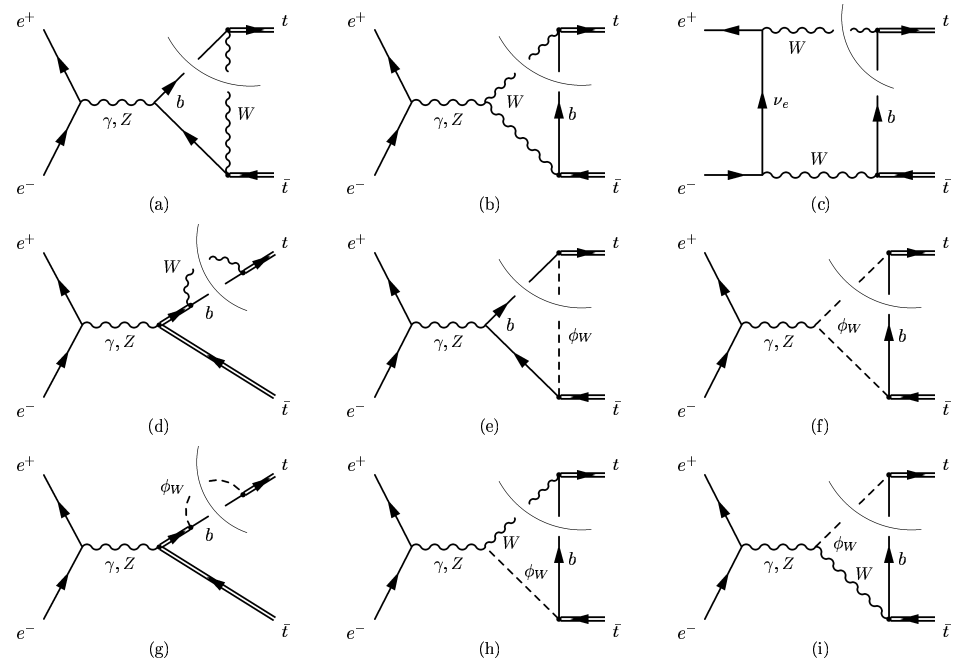
alt. scheme for α_{QED} :
 $\alpha^{\overline{\text{MS}}, n_f=8}(\mu = m_t)$

$m_t = 175 \text{ GeV}$



Finite Lifetime Effects

Currents:



Hard electroweak &
QCD matching corrections

bW^+ and $\bar{b}W^-$ cuts

$$\mathbf{O}_p = \left[C^{LL} + C^{NLL} + C^{NNLL} + iC_{abs}^{NNLL} \dots \right] \cdot \left(\begin{array}{c} e^+ \\ e^- \end{array} \rightarrow \begin{array}{c} t \\ \bar{t} \end{array} \right) + \dots$$



Finite Lifetime Effects

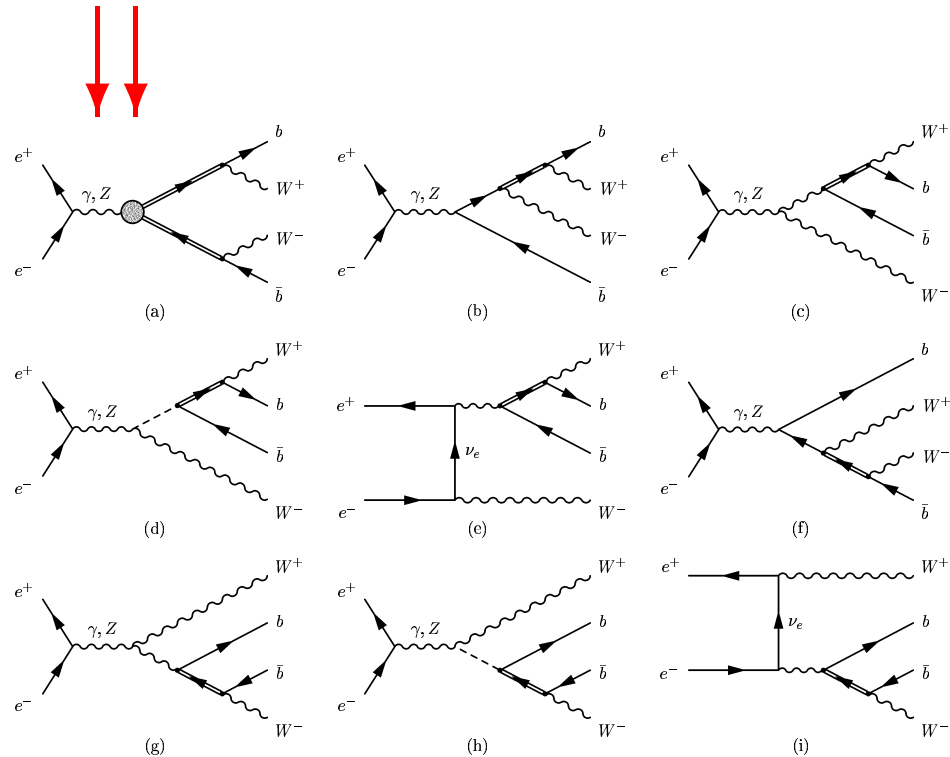
Currents:

$$\mathbf{O}_p = \left[C^{LL} + C^{NLL} + C^{NNLL} + iC_{abs}^{NNLL} \dots \right] \cdot \left(\begin{array}{cc} e^+ & t \\ e^- & \bar{t} \end{array} \right) + \dots$$

$$\sigma_{tot} \propto \text{Im} [C(\nu)^2 G(0, 0, \sqrt{s} + i\Gamma_t)]$$

- accounts for **irreducible interference** contributions:

resonant \leftrightarrow non-resonant
 $W^+ W^- b \bar{b}$ final states



Finite Lifetime Effects

Currents:

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$$\sigma_{tot} \propto \text{Im} [C(\nu)^2 G(0, 0, \sqrt{s} + i\Gamma_t)]$$



- $(\Delta\sigma_{tot}^\Gamma) \sim \alpha_s \Gamma_t \frac{1}{\epsilon} \Rightarrow$ **NLL** logarithmic phase space UV divergences

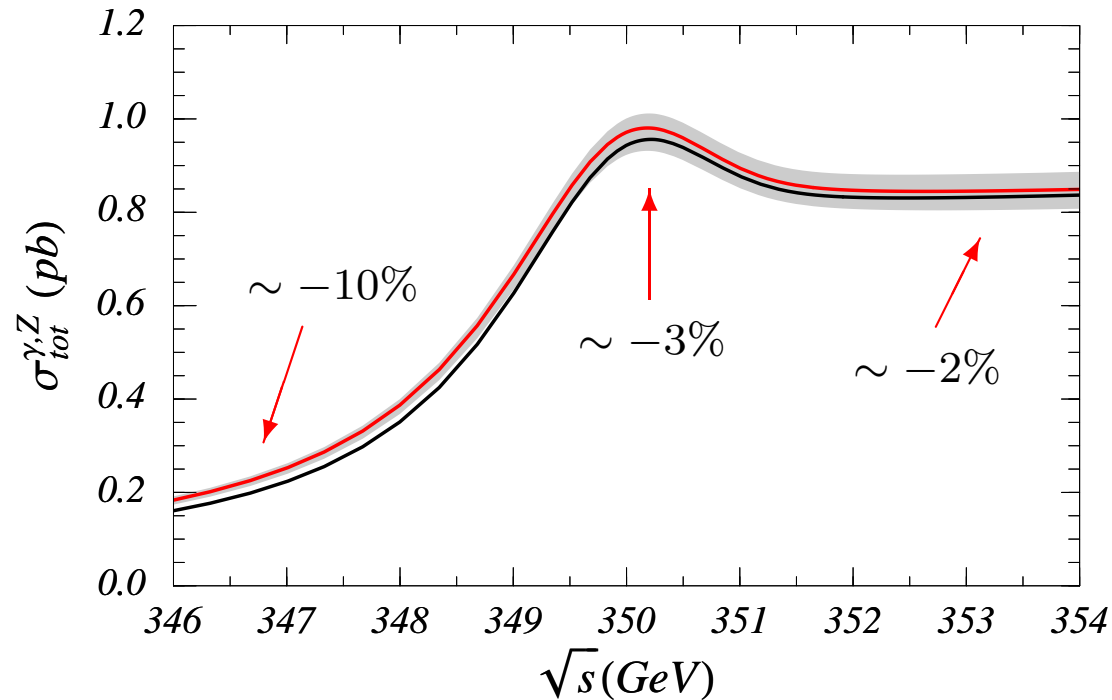
→ **NLL** anom. dim. for $(e^+e^-)(e^+e^-)$ operator → $iC(\mu) \cdot \left(\begin{array}{cc} e^+ & e^- \\ e^- & e^+ \end{array} \right)$ ✓

→ matching for $iC(\mu)$:

- physical $W^+W^-b\bar{b}$ phase space → “Phase Space Matching” w.i.p.
- $W^+W^-b\bar{b}$ final state without tops input: Rieman, Kolodziej '05



Total Cross Section



Reisser, AHH

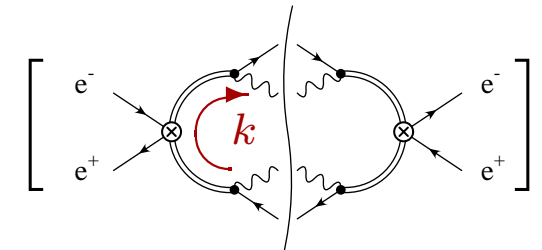
- also included summation of phase space logs $\sim (\alpha_s \ln v)^n$
- finite lifetime corrections comparable to NNLL QCD corrections
- shift in the peak position: 30 – 50 MeV $(\delta m_t^{\text{ex}} \approx 50 \text{ MeV})$

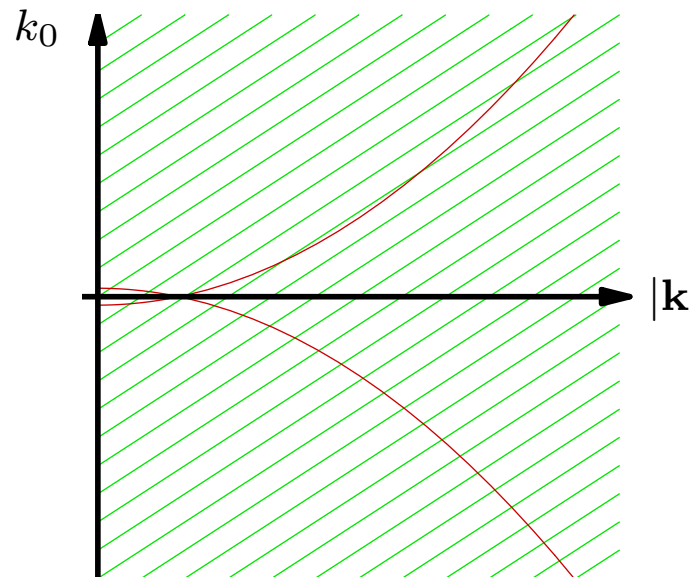


Phase Space Matching

Reisser, Ruiz-Femenia, AH

- required for σ_{tot} from optical theorem in EFT
 - finite imaginary matching conditions for all EFT operators
- removal of unphysical EFT phase space contributions

$$\sigma_{\text{tot}} \sim \left[\begin{array}{c} e^- \\ e^+ \end{array} \right] \left[\begin{array}{c} e^- \\ e^+ \end{array} \right] \int_{-\infty}^{+\infty} dk_0 \int_0^{+\infty} d|\mathbf{k}| \frac{|\mathbf{k}|^2 \Gamma_t^2}{|k_0 - \frac{\mathbf{k}^2}{2m_t^2} + i\Gamma_t|^2 | -k_0 - \frac{\mathbf{k}^2}{2m_t^2} + i\Gamma_t|^2}$$




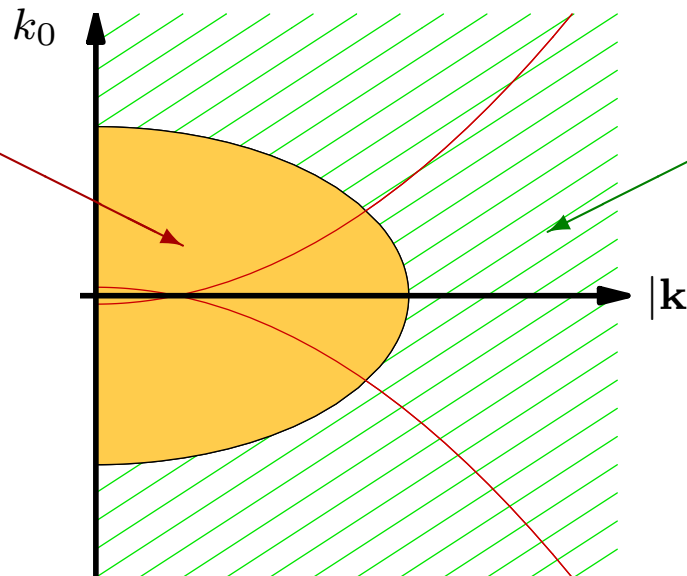
Phase Space Matching

Reisser, Ruiz-Femenia, AH

- required for σ_{tot} from optical theorem in EFT
 - finite imaginary matching conditions for all EFT operators
- removal of unphysical EFT phase space contributions

$$\sigma_{\text{tot}} \sim \left[\begin{array}{c} e^- \\ e^+ \end{array} \right] \left[\begin{array}{c} \text{loop} \\ \text{loop} \end{array} \right] \int_{-\infty}^{+\infty} dk_0 \int_0^{+\infty} d|\mathbf{k}| \frac{|\mathbf{k}|^2 \Gamma_t^2}{|k_0 - \frac{\mathbf{k}^2}{2m_t^2} + i\Gamma_t|^2 | -k_0 - \frac{\mathbf{k}^2}{2m_t^2} + i\Gamma_t|^2}$$

$(q^2 - m_t^2) \ll m_t^2$
nonrel. exp. valid
double resonant



unphysical region of EFT
single/non-resonant
subtracted in local exp.

Phase Space Matching

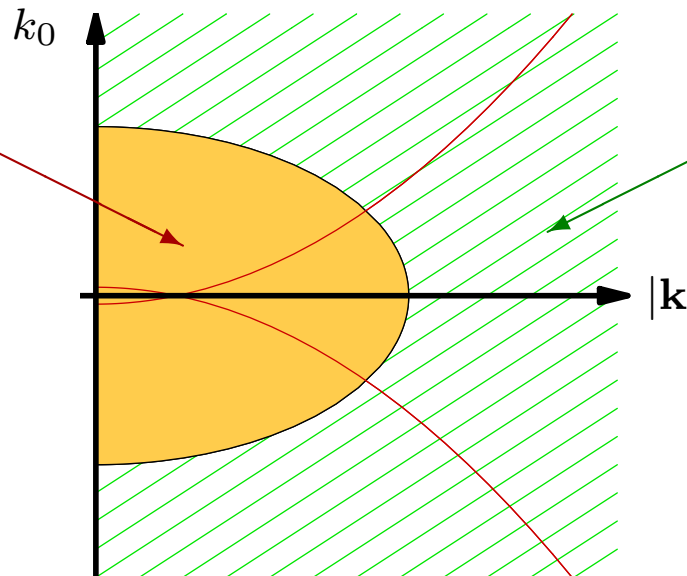
Reisser, Ruiz-Femenia, AH

- required for σ_{tot} from optical theorem in EFT
- finite imaginary matching conditions for all EFT operators
 → removal of unphysical EFT phase space contributions

$$\sigma_{\text{tot}} \sim \left[\text{Diagram 1} \right] + \text{Im} \left[\text{Diagram 2} + \text{Diagram 3} + \dots \right]$$

The equation shows the total cross-section σ_{tot} as a sum of a tree-level diagram and the imaginary part of a series of loop diagrams. Diagram 1 is a tree-level exchange of a photon between two electron-positron pairs. Diagram 2 is a loop diagram with a vertex $iC(\nu)$. Diagram 3 is a loop diagram with a vertex $i\tilde{C}(\nu) \frac{D^2}{m^2}$.

$(q^2 - m_t^2) \ll m_t^2$
 nonrel. exp. valid
 double resonant



unphysical region of EFT
 single/non-resonant
 subtracted in local exp.

$C(1), \tilde{C}(1), \dots$
 (match. cond. for $4e$ ops.)



Phase Space Matching

Reisser, Ruiz-Femenia, AH

- required for σ_{tot} from optical theorem in EFT
- finite imaginary matching conditions (counterterms) for all EFT operators
→ removal of unphysical EFT phase space contributions
- imaginary (finite) renormalization
- separation of hard and soft phase space contributions
 - ⊕ makes summation of phase space logs meaningful
 - ⊖ no fully differential predictions
- practical issue: “mild” power counting breaking effects
→ dependence on experimental cuts
- S-wave production: contributes at NLL order
- P-, D-, . . . wave states: mandatory for consistent prediction at LL order (e.g. $e^+e^- \rightarrow t\bar{t}$)

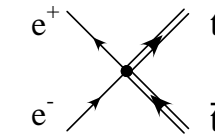
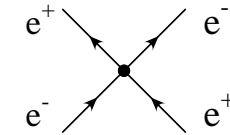


What's left to do ... a lot!

NNLL Total Cross Section:

(A) electroweak corrections

- NLL + NNLL phase space matching $\rightarrow \tilde{C}^{eeee}(\mu = m_t)$
- NNLL running of $\tilde{C}^{eeee}(\mu)$
- $\mathcal{O}(\alpha_s)$ corr.'s to $iC_{\text{abs}}^{eett}(\mu = m_t)$ for NNLL mixing effects
- QED effects: ISR, beamstrahlung, Coulomb singularities



(B) QCD corrections (RGE-improved)

- Current uncertainty: $(\delta\sigma_{t\bar{t}}/\sigma_{t\bar{t}})^{\text{QCD}} \sim \pm 6\%$
 \rightarrow full NLL running of potentials

Fully Differential Cross Sections:

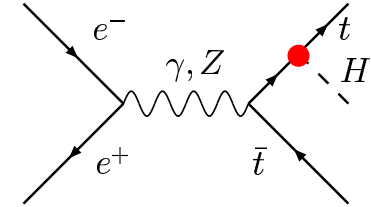
- \rightarrow efficiencies ? really needed ? interesting theoretically ?
- \rightarrow more involved treatment of unstable particles



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ top-Yukawa coupling



- Theory Status: $\sigma(e^+e^- \rightarrow t\bar{t}H)$
 - Born ✓
 - 1-loop ew. ✓
 - $\mathcal{O}(\alpha_s)$ fixed-order ✓

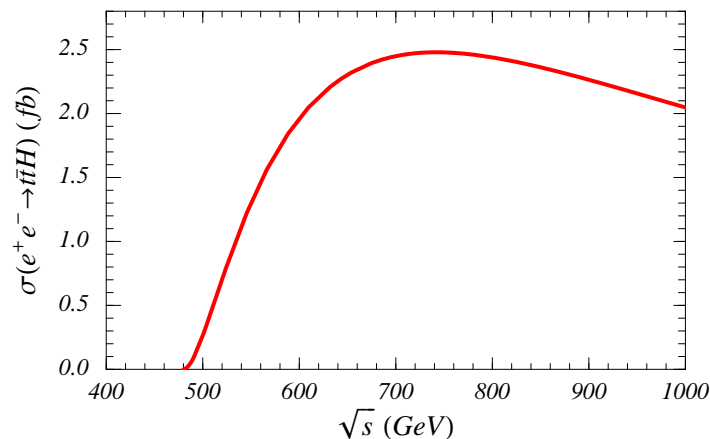
[Gaemers et al., Djouadi et al.]

[Denner et al., Belanger et al., You et al.]

[Dittmaier et al., Dawson et al.]

NLL large- E_H QCD endpoint corrections ✓

[C. Farrell, AHH]



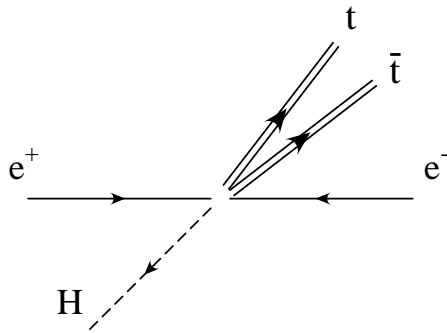
- tiny cross section for $\sqrt{s} = 500$ GeV
- measurement of Yukawa coupling difficult
- $\delta Y_t / Y_t \sim 30\%$ feasible at $\sqrt{s} = 500$ GeV [A. Juste, 2002]



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ region of large Higgs energy



→ $t\bar{t}$ collinear

→ QCD effects localized in $t\bar{t}$ system

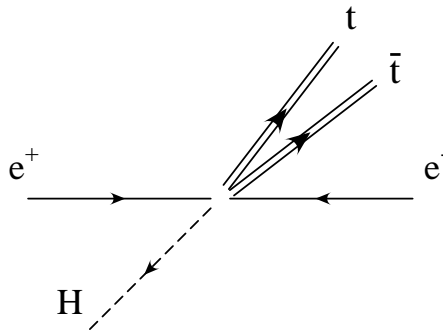
⇒ $t\bar{t}$ dynamics non-relativistic



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

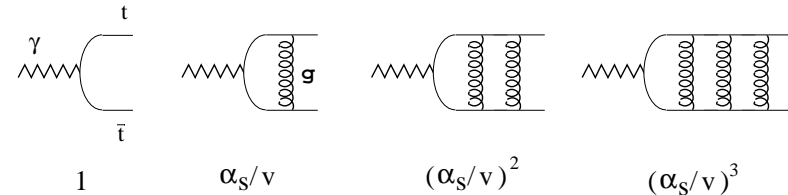
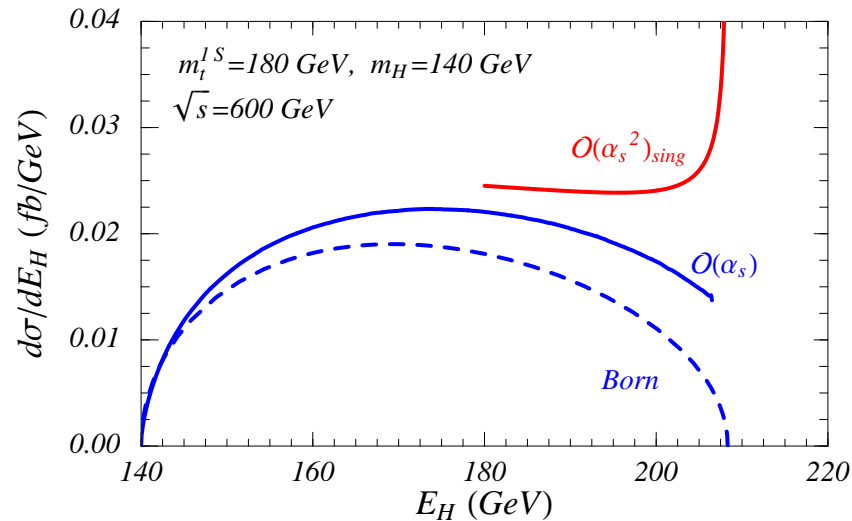
→ region of large Higgs energy



→ $t\bar{t}$ collinear

→ QCD effects localized in $t\bar{t}$ system

⇒ $t\bar{t}$ dynamics non-relativistic



→ singularities: $\sim (\alpha_s/v)^n$,

$\sim (\alpha_s \ln v)^n$

→ fixed order expansion breaks down

⇒ summation of singular terms

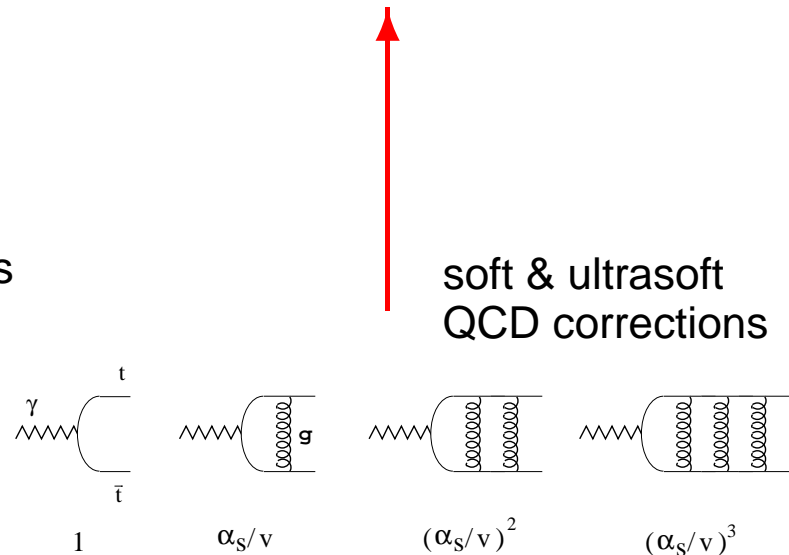
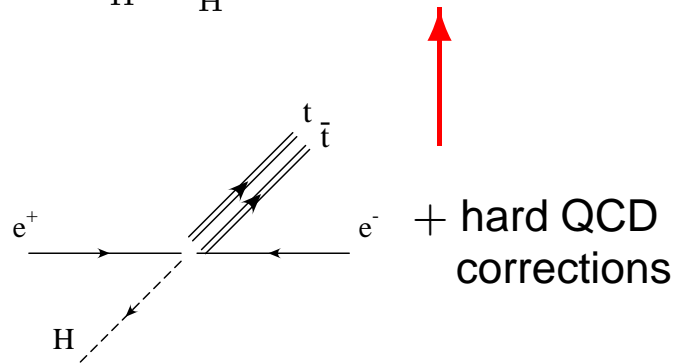


Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ factorization formula

$$\left(\frac{d\sigma}{dE_H}\right)_{E_H \approx E_H^{\max}} \sim C^2(\mu, \sqrt{s}, m_t, m_H) \times \text{Im}[G(0, 0, v, \mu)]$$



NLL formalism: Cailin Farrell, AH; Phys.Rev.D72,014007 (2005)

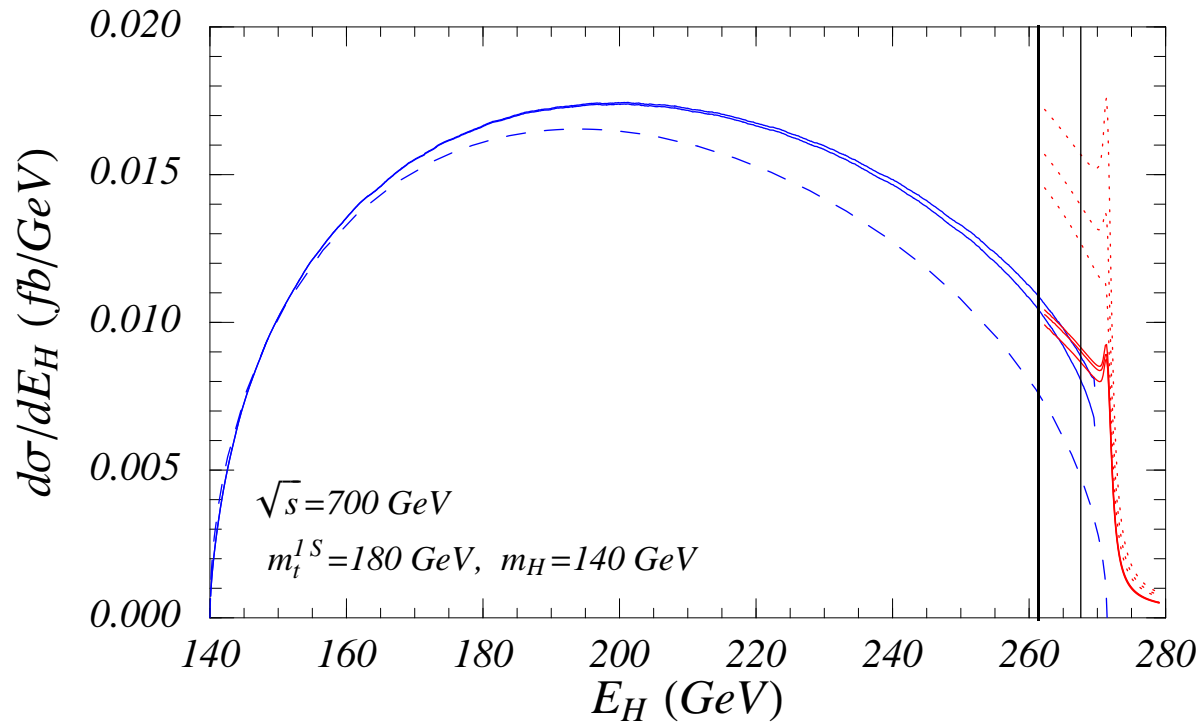
Cailin Farrell, AH; hep-ph/0604166



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ NLL Higgs energy spectrum



Farrell, AH '05

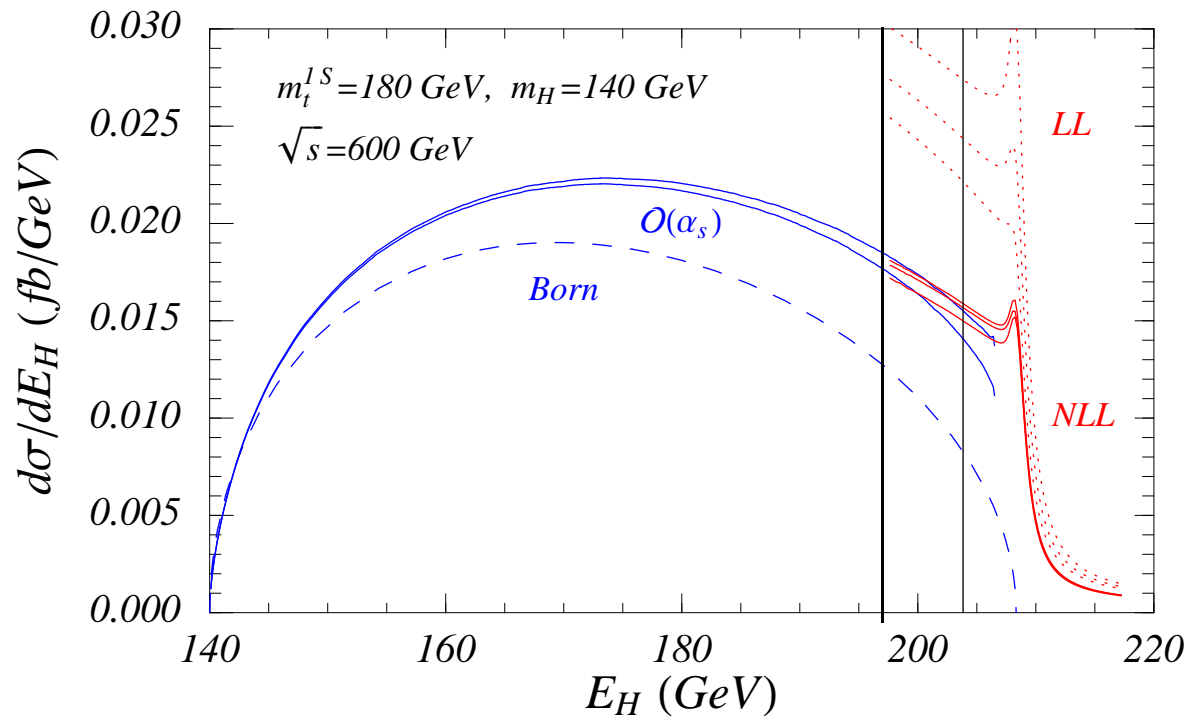
- large E_H endpoint regions increases for smaller \sqrt{s} / larger m_H
- NLL matching: full theory results from Denner, Dittmaier, Roth, Weber '04



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ NLL Higgs energy spectrum



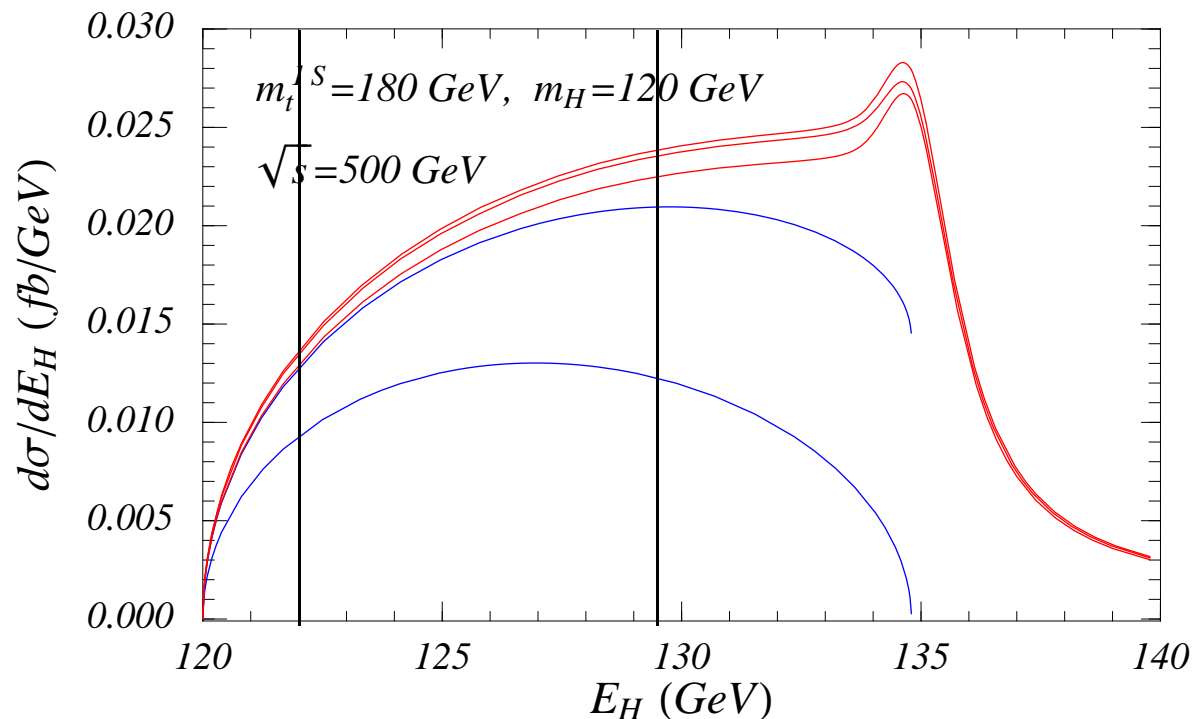
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Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ NLL Higgs energy spectrum



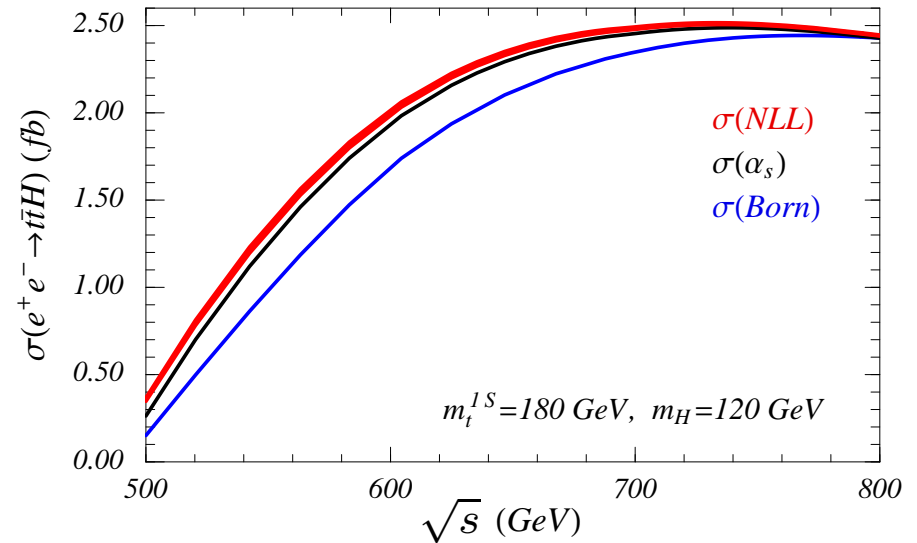
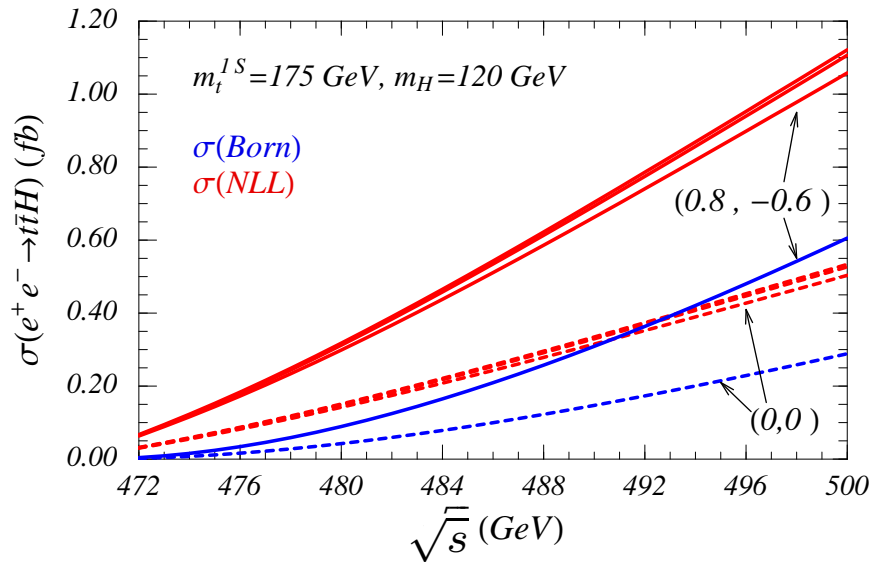
- large E_H endpoint regions increases for smaller \sqrt{s} / larger m_H
- NLL matching: full theory results from [Denner, Dittmaier, Roth, Weber '04](#)
- $\sqrt{s} = 500 \text{ GeV}$: extension of factorization formula for low- E_H endpoint



Other Application

$$e^+e^- \rightarrow t\bar{t}H$$

→ total cross section



500 GeV:

- **factor 2 enhancement** over tree level from summation of $(\alpha_s/v)^n, (\alpha_s \ln v)^n$ terms
- **another factor of 2 enhancement** for $P_- = -80\%, P_+ = +60\%$
- essential for realistic studies for ILC (phase I) **Juste '02, '06**

$$\Rightarrow (\delta\lambda_t/\lambda_t)_{500 \text{ GeV}}^{\text{ILC}} \sim 30\% \quad \rightsquigarrow \quad 10 - 15\%$$



Conclusion

- Top pair total rate predictions become more and more realistic
 - missing NNLL QCD corrections w.i.p. $\rightarrow d\sigma/\sigma < \pm 6\%$
 - full set of electroweak corrections (“hard” ✓, “QED” ?)
 - top quark finite lifetime: imaginary Wilson coefficients
 - $\rightarrow bW$ final states
 - \rightarrow phase space matching
- reaching goals $\delta m_t \sim 100$ MeV is no totally free lunch
still considerable work has to be invested
to get final number
- Many interesting applications of threshold physics exist.



Colors

This is blue

This is red

This is brown

This is magenta

This is Dark Green

This is Dark Blue

This is Green

This is Cyan

Test how this color looks

Test how this color looks

Test how this color looks

Test how this color looks

Test how this color looks

