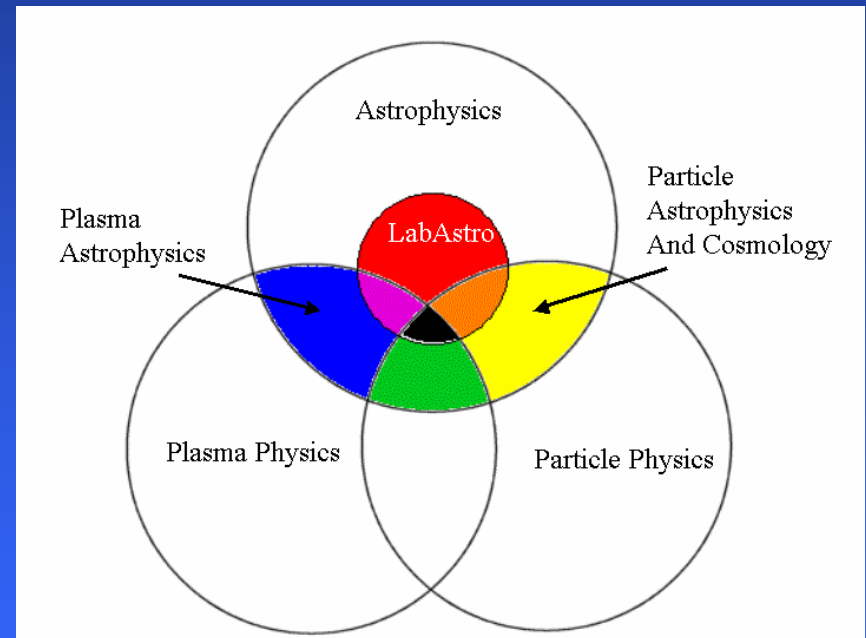


An Astrophysical Plasma Wakefield Accelerator

Alfven Wave Induced
Plasma Wakefield Acceleration

Laboratory Astrophysics at SLAC

- Study in a Laboratory setting:
 - Fundamental physics
 - Astrophysical Dynamics
 - Phenomenon relevant to astrophysics experiments: calibration of techniques, etc



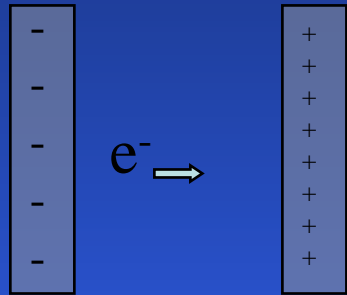
Plasma Wakefield Acceleration

- ARD-B research focus area here at SLAC.
- Laser or beam induced plasma wakefield acceleration.
- At a certain distance behind wake causing disturbance there is a “constant” region of acceleration.

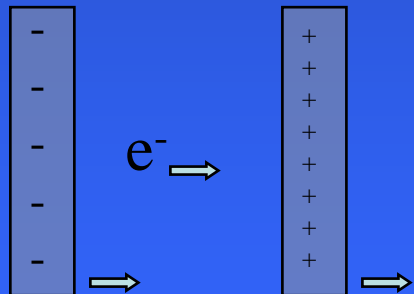


Lake Como

Wake-Field Accelerator

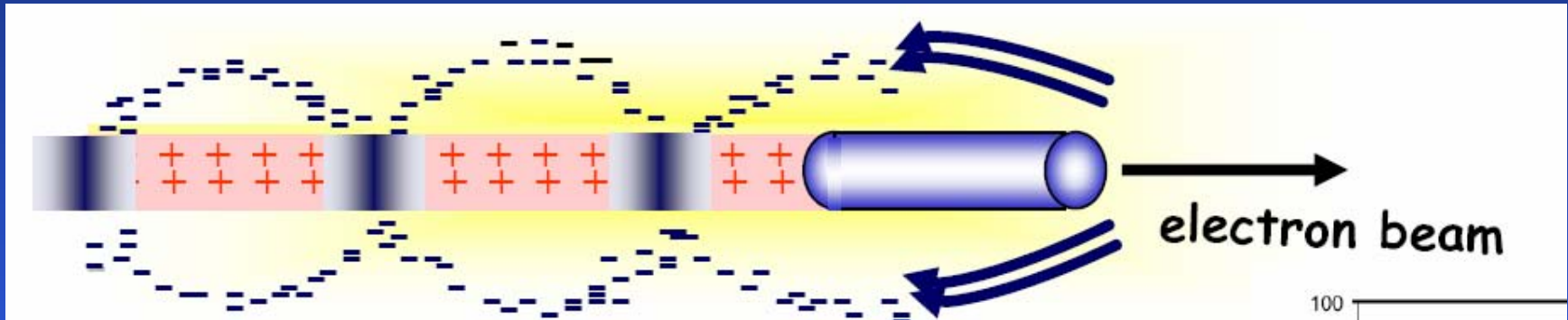


When $v_{\text{grad}} \approx c$ the acceleration gradient
Will co-move with the particle ($v_{\text{part}} \approx c$)
Providing a great deal of acceleration!



In our case the gradient is provided by
an Alfven wave wake-field. The velocity
of the gradient is v_{alfven} .

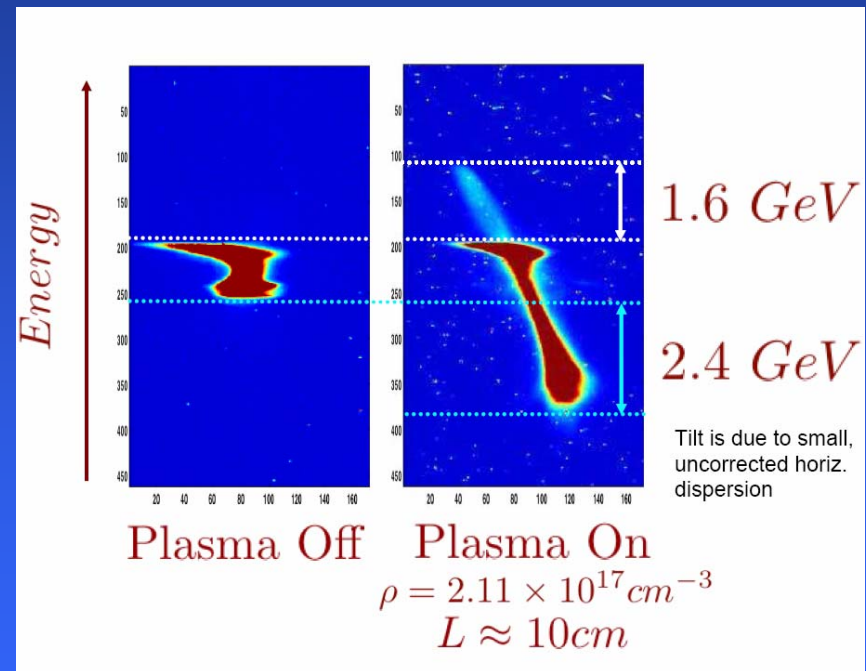
Beam Induced Plasma Wakefield Accelerator



- E164 (X,XX,Y,...)
- Leading edge of beam pulse ionizes ambient plasma and creates wake.
- Portion of trailing edge is inside the acceleration region of the wake.
- (Note: Also a large deceleration region.)

E-164 Results

- Plasma Wakefields generate extremely large scale gradients.
- >10 GeV per meter observed (~ 2 GeV over 10 cm in E164X).
- Compare to ~ 10 GeV per km at SLAC.



Working Group 1

WHAT MAKES AN IDEAL ACCELERATOR?

LESSONS FROM TERRISTRIAL ACCELERATORS

- **Continuous interaction** between the particle and the accelerating longitudinal EM field (Lorentz inv.)
 - Gain energy in macroscopic distance
- Particle-field interaction **process non-collisional**
 - Avoid energy loss through inelastic scatterings
- To reach ultra high energy, **linear acceleration** (minimum bending) is the way to go
 - Avoid severe energy loss through synchrotron radiation
- **Are these criteria applicable to celestial accelerators?**

LINEAR VS. CIRCULAR

SLAC CERN



A Brief History of Plasma Wakefields

Motivated by the challenge of high energy physics

- Laser driven plasma acceleration
T. Tajima and J. M. Dawson (1979)
- Particle-beam driven plasma wakefield acceleration
PC, Dawson et al. (1984)
- **Extremely efficient:**

$$eE \geq \sqrt{n [\text{cm}^{-3}]} \text{ eV/cm}$$

For $n=10^{18} \text{ cm}^{-3}$, $eE=100 \text{ GeV/m} \rightarrow \text{TeV collider in 10 m!}$

- * Plasma wakefield acceleration principle **experimentally verified**. Actively studied worldwide

Concepts For Plasma-Based Accelerators

□ Laser Wake Field Accelerator(LWFA)

A single short-pulse of photons

□ Plasma Beat Wave Accelerator(PBWA)

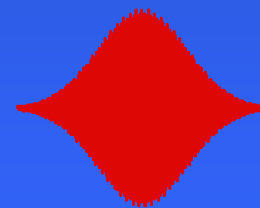
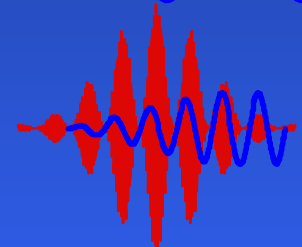
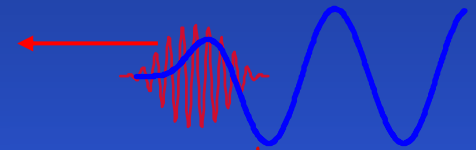
Two-frequencies, i.e., a train of pulses

□ Self Modulated Laser Wake Field Accelerator(SMLWFA)

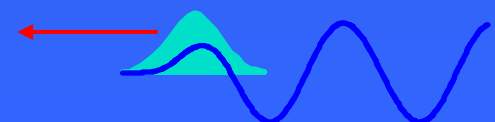
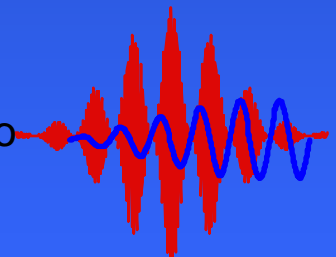
Raman forward scattering instability

□ Plasma Wake Field Accelerator(PWFA)

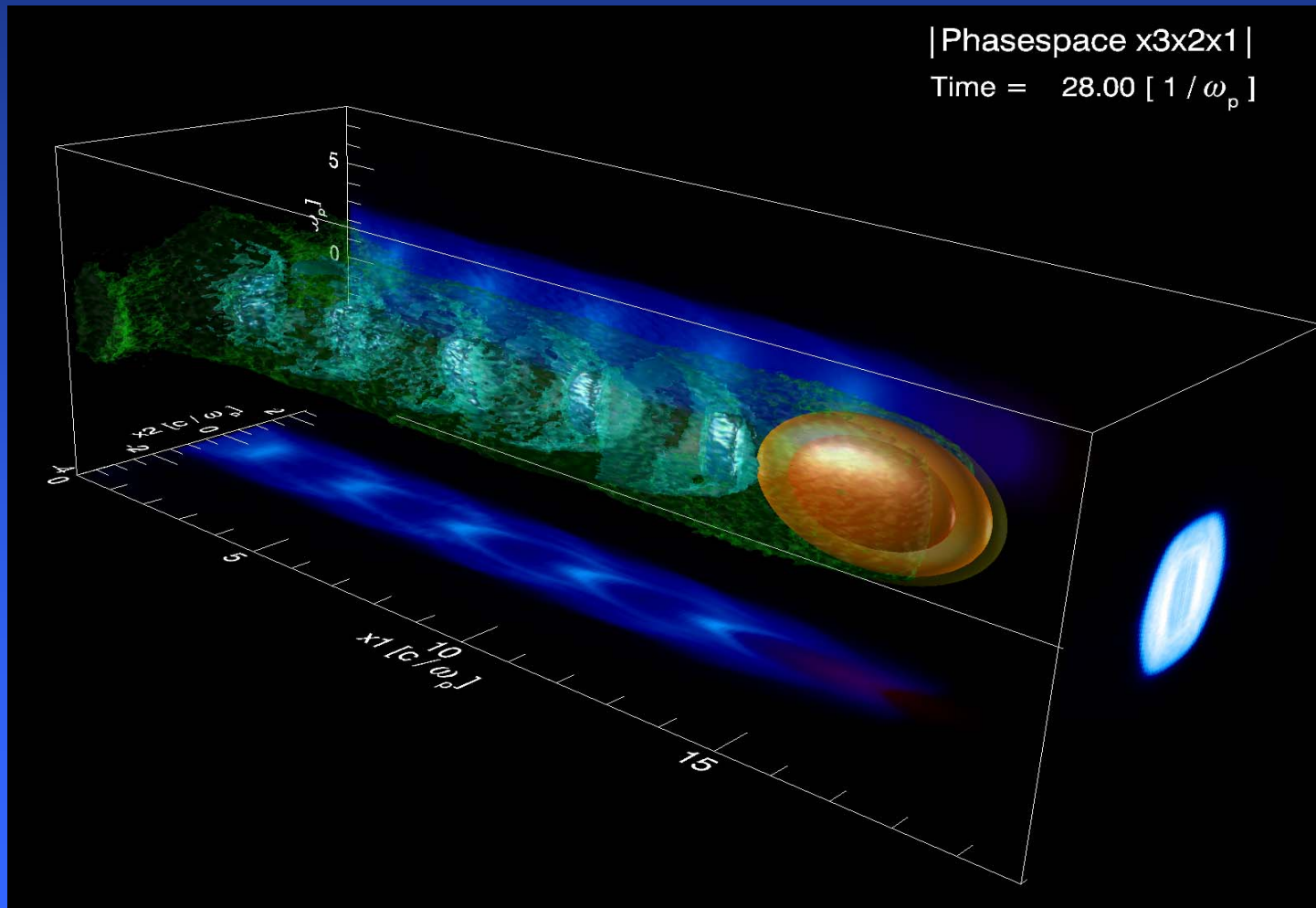
A high energy electron (or positron) bunch



evolves to



Plasma Wakefield Simulation (SLAC E-157 Collaboration)



Cosmic Accelerators

- 1 of the 11 questions for the new century from Turner Committee:
 - **“How do cosmic accelerators work and what are they accelerating?”**
- 2 of their 7 recommendations:
 - **“Determine the origin of the highest-energy gamma-rays, neutrinos, and cosmic rays.”**
 - **“Discern the physical principles that govern extreme astrophysical environments through the laboratory study of high energy-density physics.** The Committee recommends that the agencies cooperate in bringing together the different scientific communities that can foster this rapidly developing field.”

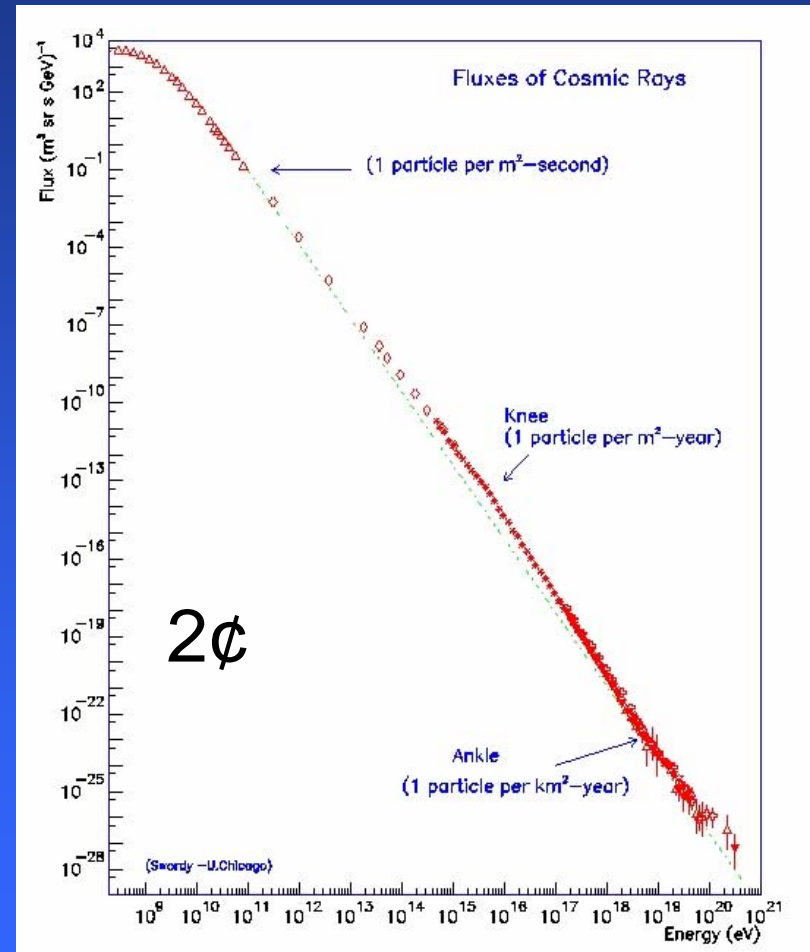
Cosmic Acceleration Mechanisms:

Addressing the Bottom-Up Scenario for Acceleration of Ordinary Particles

- Conventional cosmic acceleration mechanisms encounter limitations:
 - Fermi acceleration (1949) (= stochastic accel. bouncing off B-fields)
 - Diffusive shock acceleration (1970's) (a variant of Fermi mechanism)
 - Limitations for UHE: field strength, diffusive scattering inelastic
 - Eddington acceleration (= acceleration by photon pressure)
 - Limitation: acceleration diminishes as $1/\gamma$
- Examples of new ideas:
 - Zevatron (= unipolar induction acceleration) (R. Blandford, astro-ph/9906026, June 1999)
 - Alfvén-wave induced wakefield acceleration in relativistic plasma (Chen, Tajima, Takahashi, Phys. Rev. Lett. 89 , 161101 (2002).
 - Additional ideas by M. Baring, R. Rosner, etc.

Cosmic Acceleration

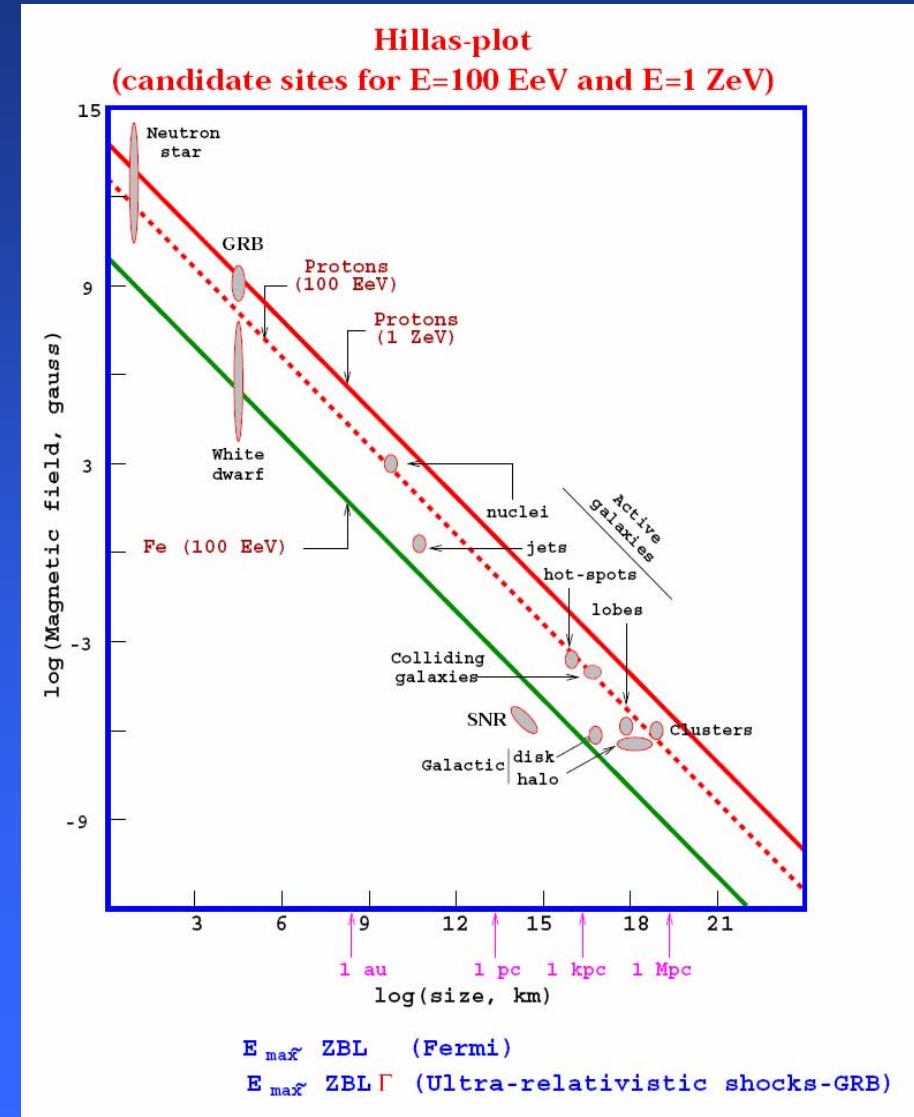
- Need to produce particles with energy in excess of 10^{20} eV.



- 94 ϕ

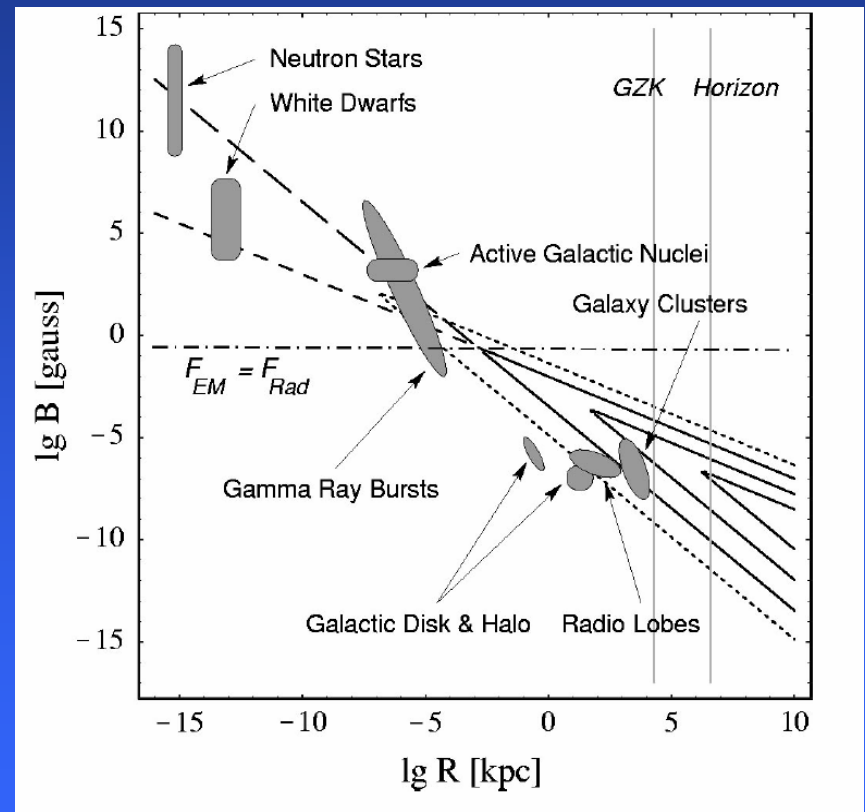
Hillas Diagram

- UHECR created as:
 - Protons receive multiple kicks of energy
 - Protons are confined inside of the acceleration region by magnetic fields.
- In order to achieve UHE:
 - many kicks required.
 - Large field needed to contain the particle for enough kicks to occur.
- Summary:
 - Need large B fields (small acceleration region) OR large acceleration region (small B field)

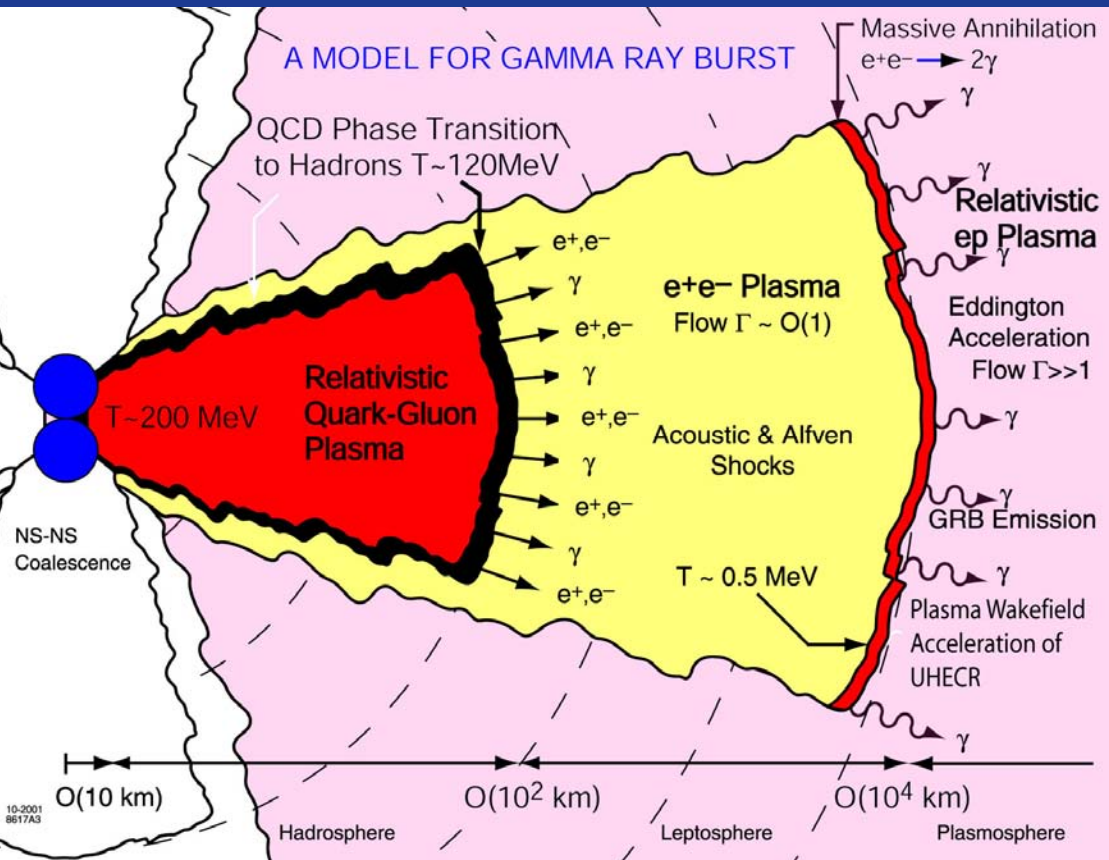


Synchrotron losses!

- Mikhail V. Medvedev,**
"Constraint on Electromagnetic Acceleration of Highest Energy Cosmic Rays",
***Physical Review E* 67, 045401 (2003).**
Abstract
- The energetics of electromagnetic **acceleration** of ultra-high-energy cosmic rays (UHECRs) is constrained both by confinement of a particle within an **acceleration** site and by radiative energy losses of the particle in the confining magnetic fields. We demonstrate that the detection of $\sim 3 \times 10^{20}$ eV events is inconsistent with the hypothesis that compact cosmic accelerators with high magnetic fields can be the sources of UHECRs. **This rules out the most popular candidates, namely spinning neutron stars, active galactic nuclei (AGNs), and γ -ray burst blast waves.** Galaxy clusters and, perhaps, AGN radio lobes remain the only possible (although not very strong) candidates for **UHECR acceleration** sites. Our analysis places no limit on linear accelerators. With the data from the future {it Auger} experiment one should be able to answer whether a conventional theory works or some new physics is required to explain the origin of UHECRs.
- PACS numbers:* 41.60.-m, 96.40.-z



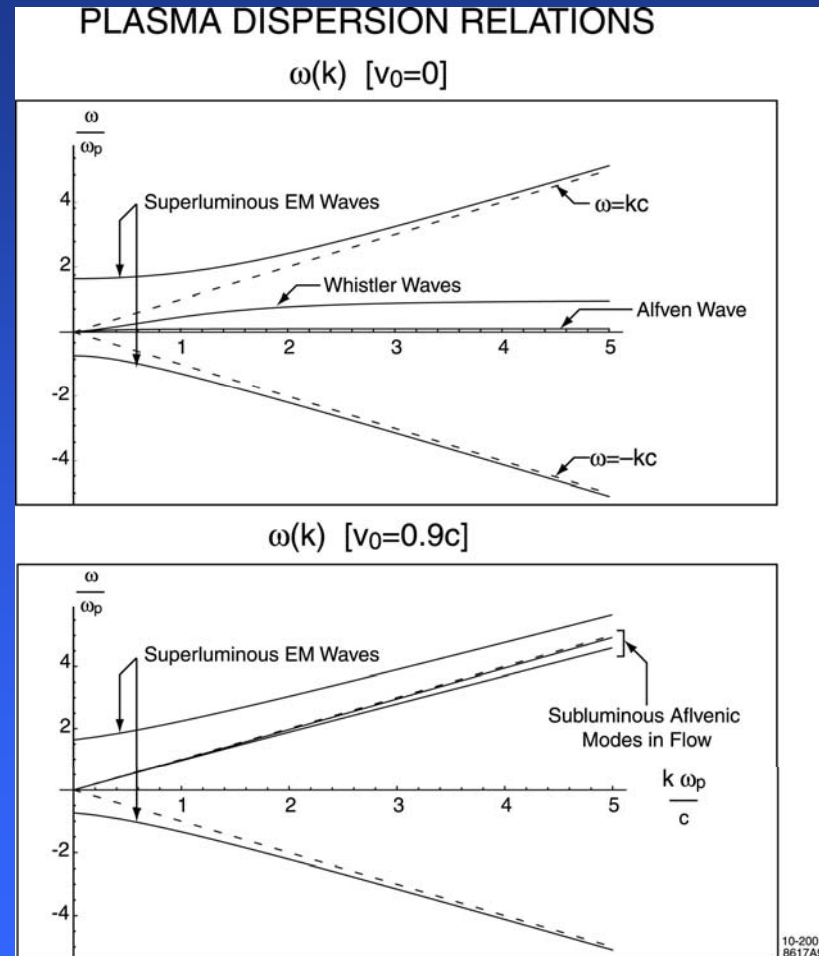
Plasma Wakefields in Astrophysics



- In accelerator research, great care must be taken to get “entire” bunch accelerated.
- In astrophysics, only a few need to be trapped in the accelerating portion of wakefield.

Alfven Shock Induced PWFA!

- In a non-relativistic plasma,
– $E_A/B_A = v_A/c \ll 1$.
- In a relativistic plasma,
– $E_A/B_A = v_A/c \approx 1$.



Alfven Wave Induced Wake Field Simulations

K. Reil (SLAC), PC and R. Sydora (U of Alberta)

Dispersion relation for EM waves in magnetized plasma:

$$\frac{k^2 c^2}{\omega^2} = 1 - \frac{2\omega_{pe}^2}{\omega^2 - \Omega_e^2} \quad \omega_{pe}^2 = 4\pi e^2 n/m$$

$$\Omega_c = eB/mc$$

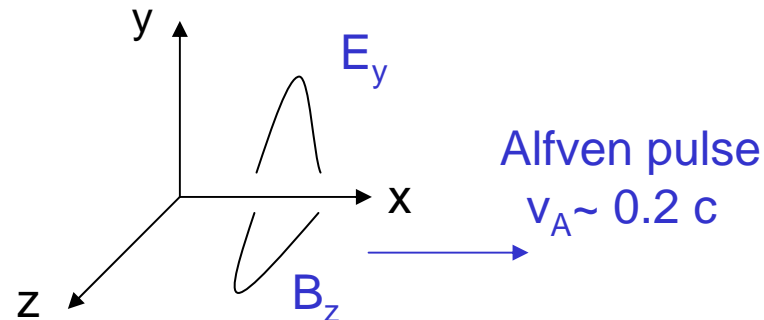
E.M. waves propagate when

(i) $\omega > \sqrt{2\omega_{pe}^2 + \Omega_e^2}$, the high frequency branch

(ii) $\omega < \Omega_e$, the low frequency (or Alfven) branch

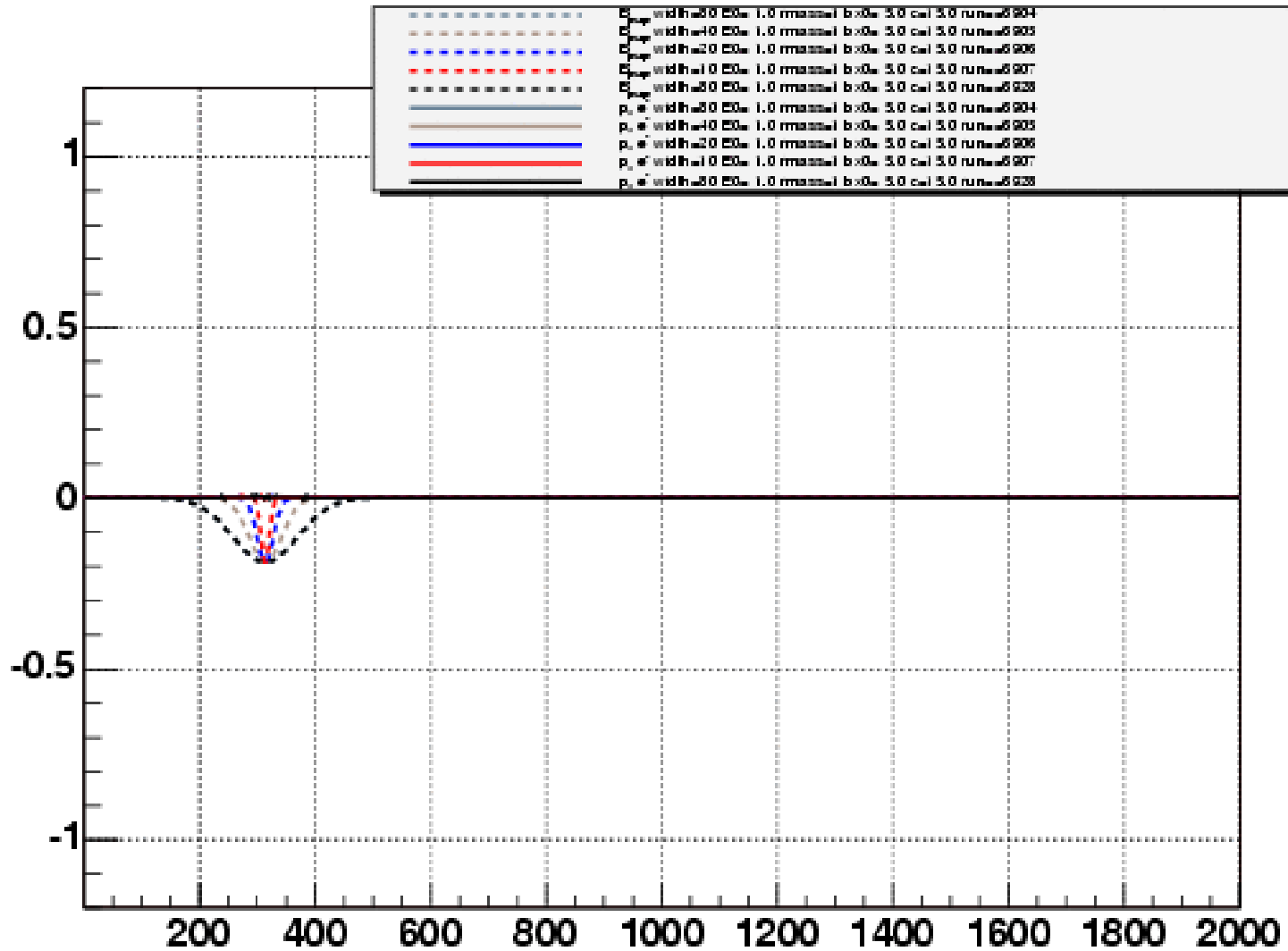
$$\omega \simeq kV_A \left(1 - \frac{k^2 \omega_{pe}^2 V_A^2}{\Omega_e^4 c^2}\right), \text{ where } V_A = \frac{c}{\sqrt{(1 + 2\frac{\omega_{pe}^2}{\Omega_e^2})}}$$

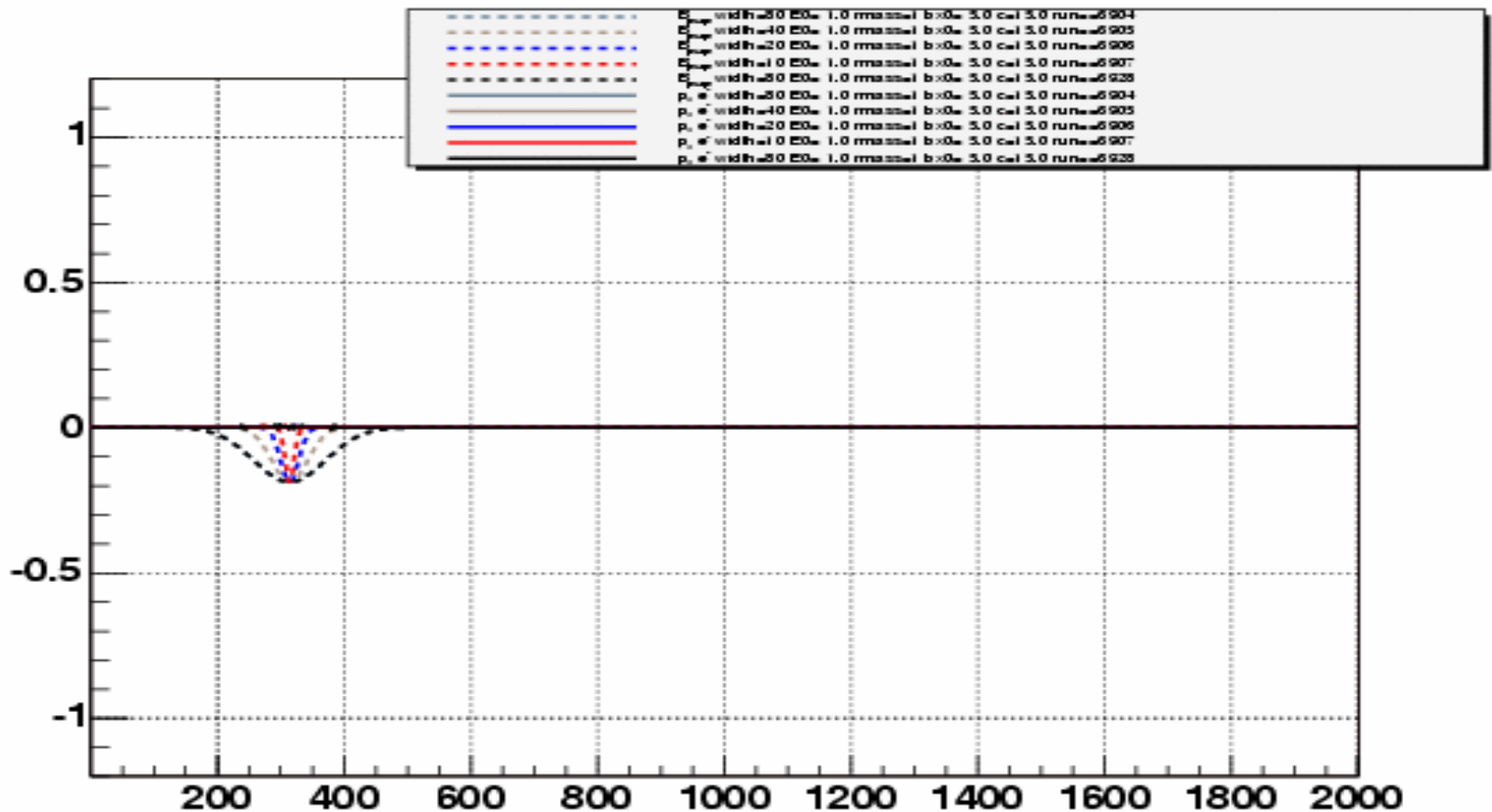
Simulation geometry:



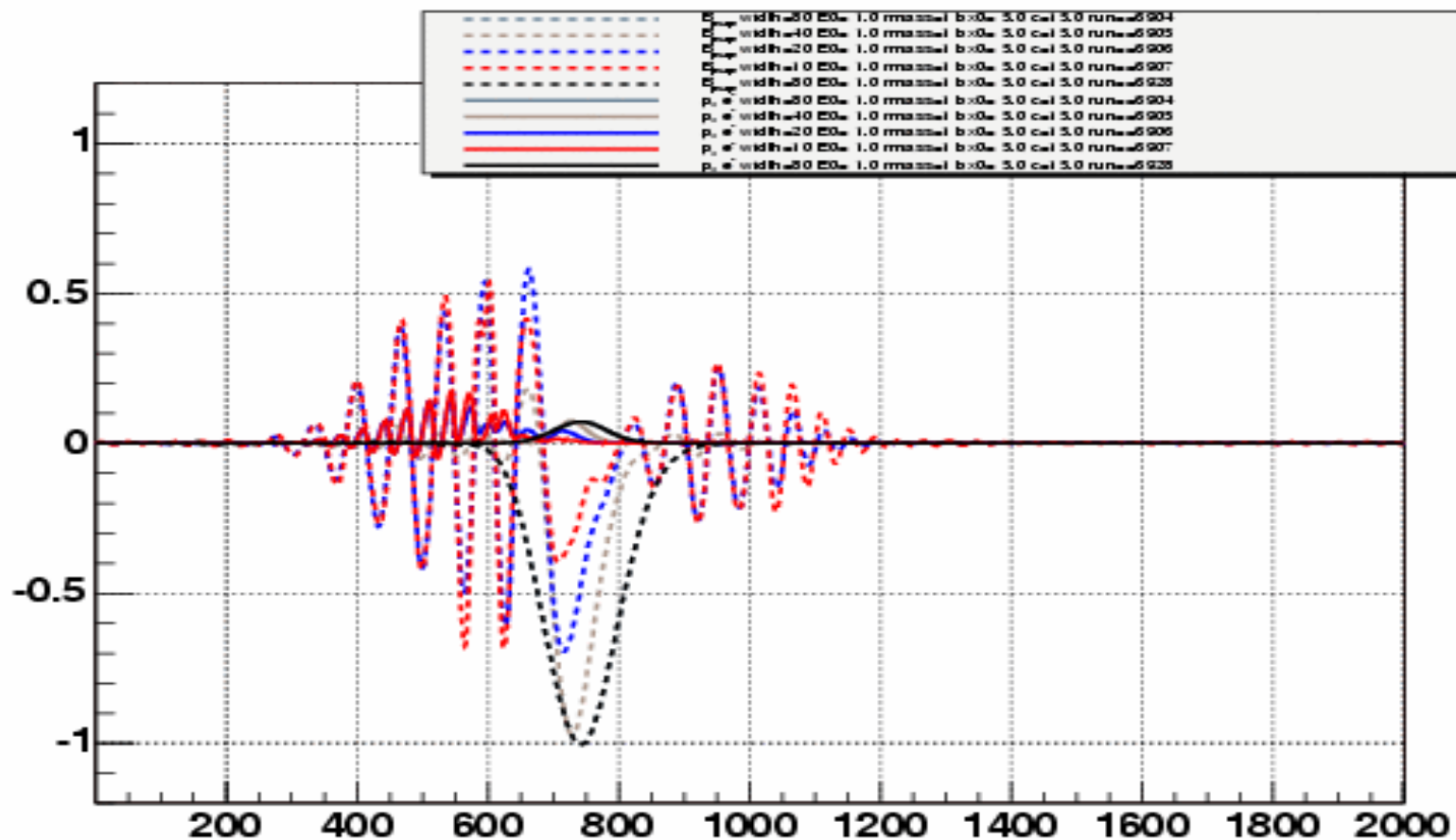
Simulation parameters for plots:

- e+ e- plasma ($m_i = m_e$)
- Zero temperature ($T_i = T_e = 0$)
- $\Omega_{ce}/\omega_{pe} = 1$ (normalized magnetic field in the x-direction)
- Normalized electron skin depth c/ω_{pe} is 15 cells long
- Total system length is $273 c/\omega_{pe}$
- $dt = 0.1 \omega_{pe}^{-1}$ and total simulation time is $300 \omega_{pe}^{-1}$
- Alfven pulse width is about $11 c/\omega_{pe}$
- 10 macroparticles per cell

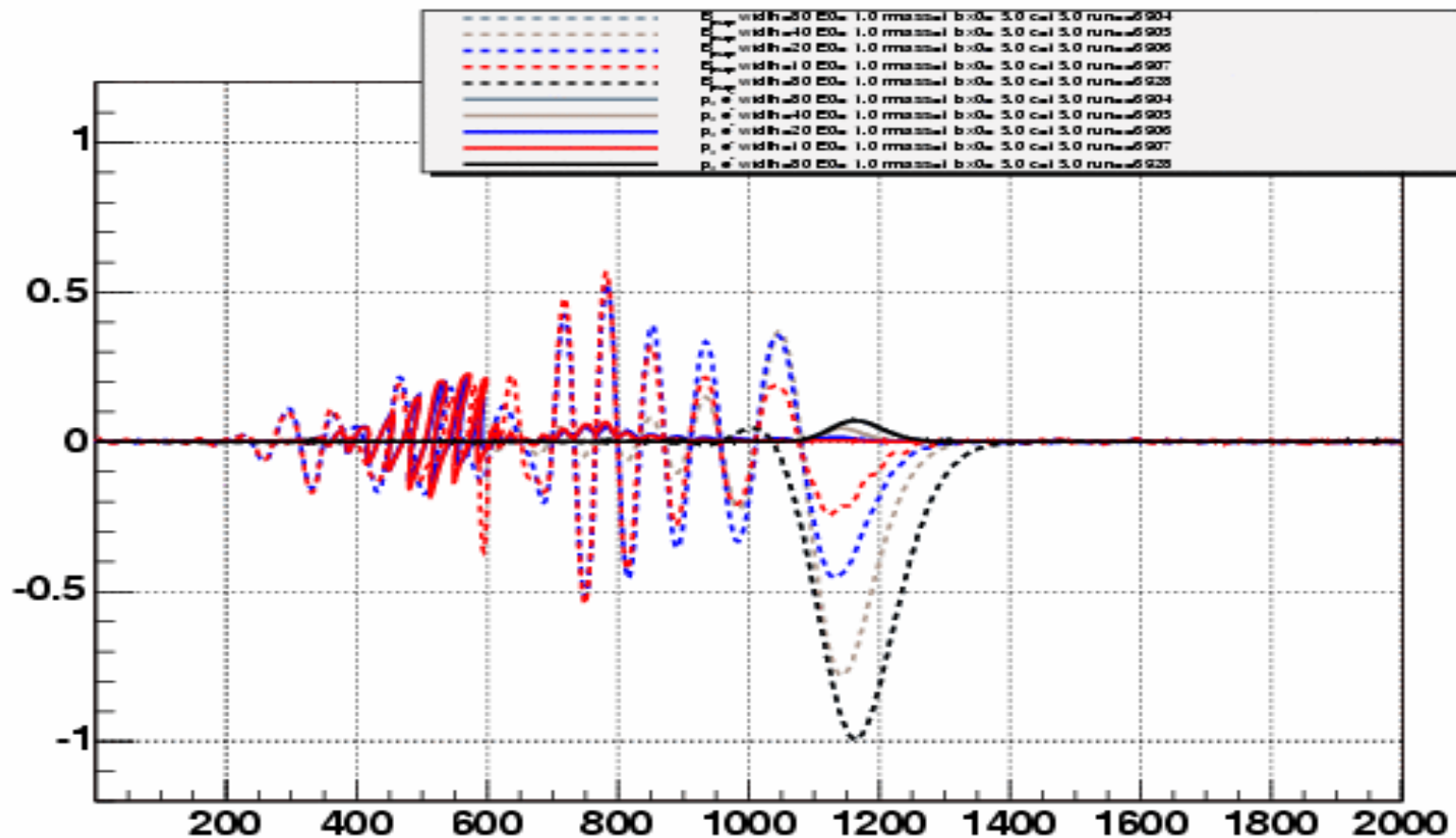




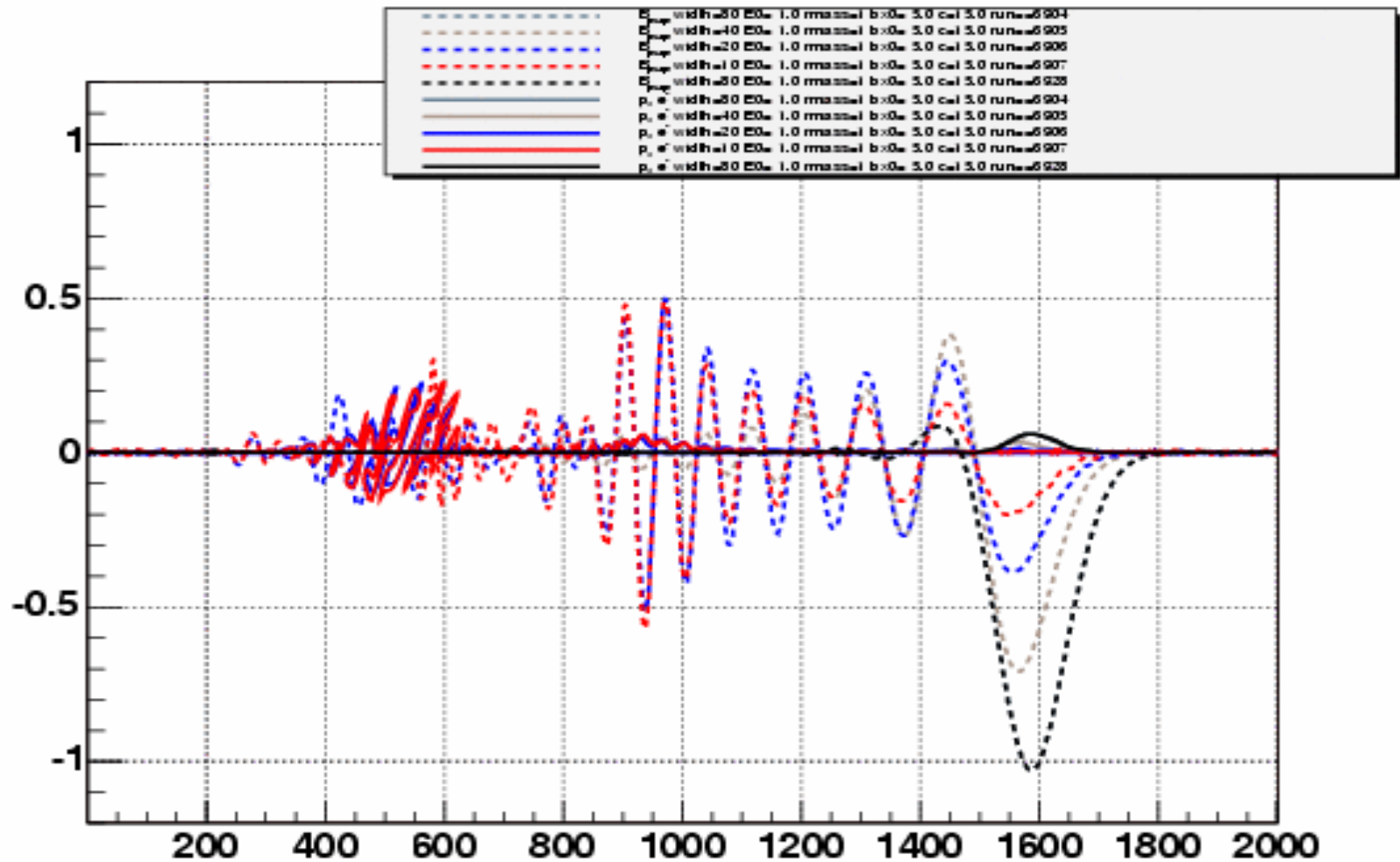
$$R_{\text{mass}}=1 \quad T=150 \quad \omega_p^{-1}$$



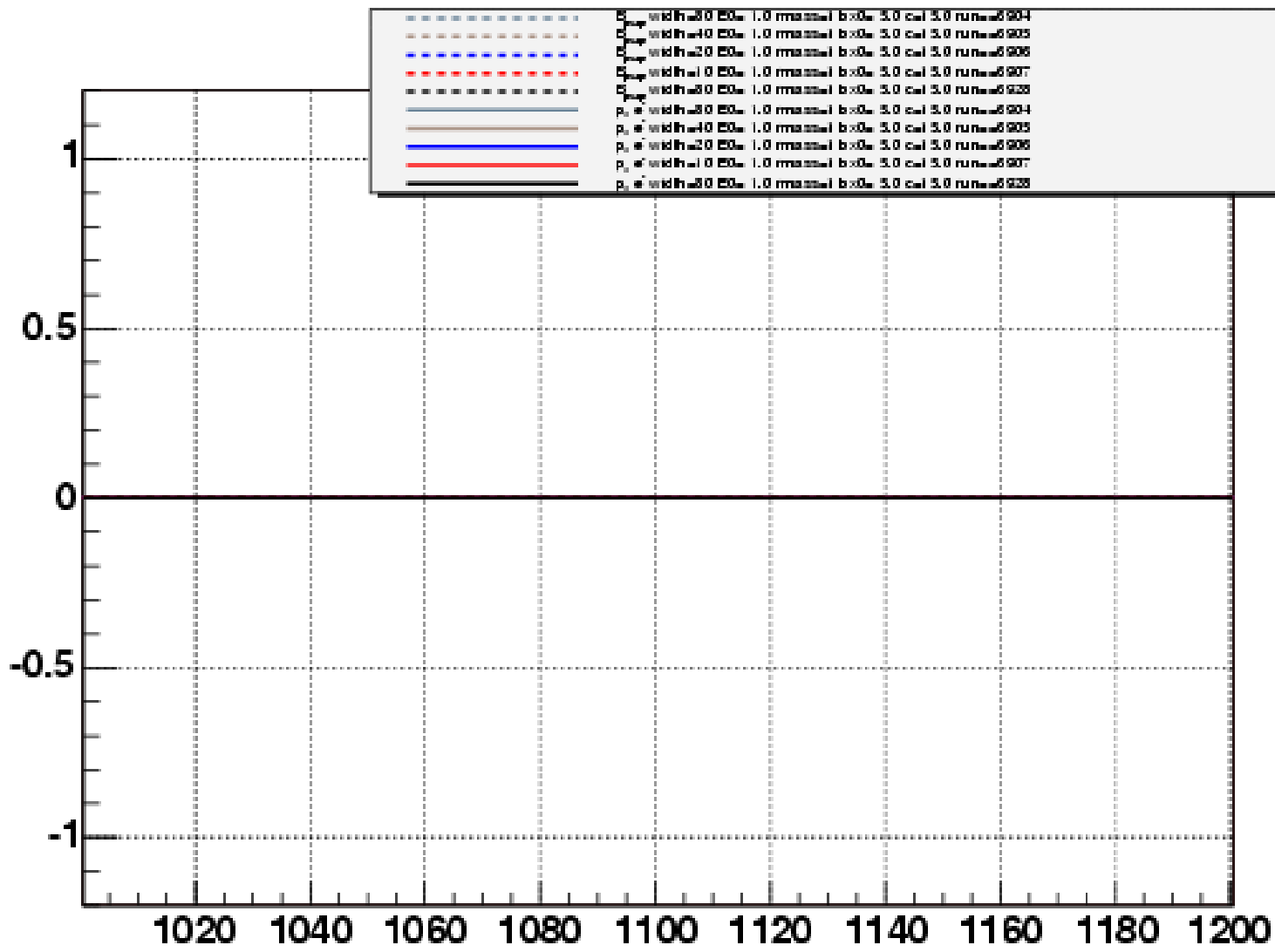
$$R_{\text{mass}}=1 \quad T=275 \, \omega_p^{-1}$$



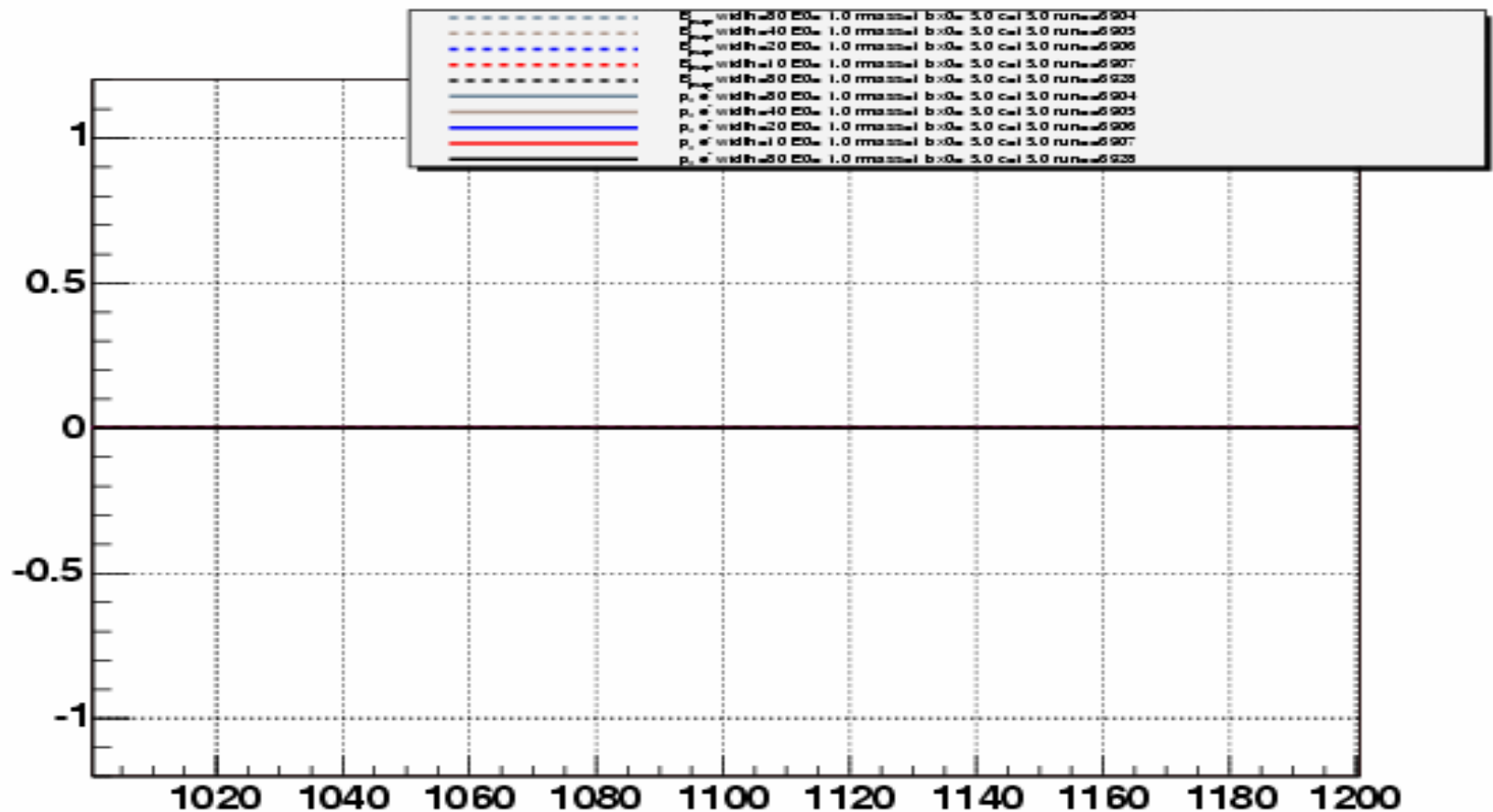
$$R_{\text{mass}}=1 \quad T=400 \quad \omega_p^{-1}$$



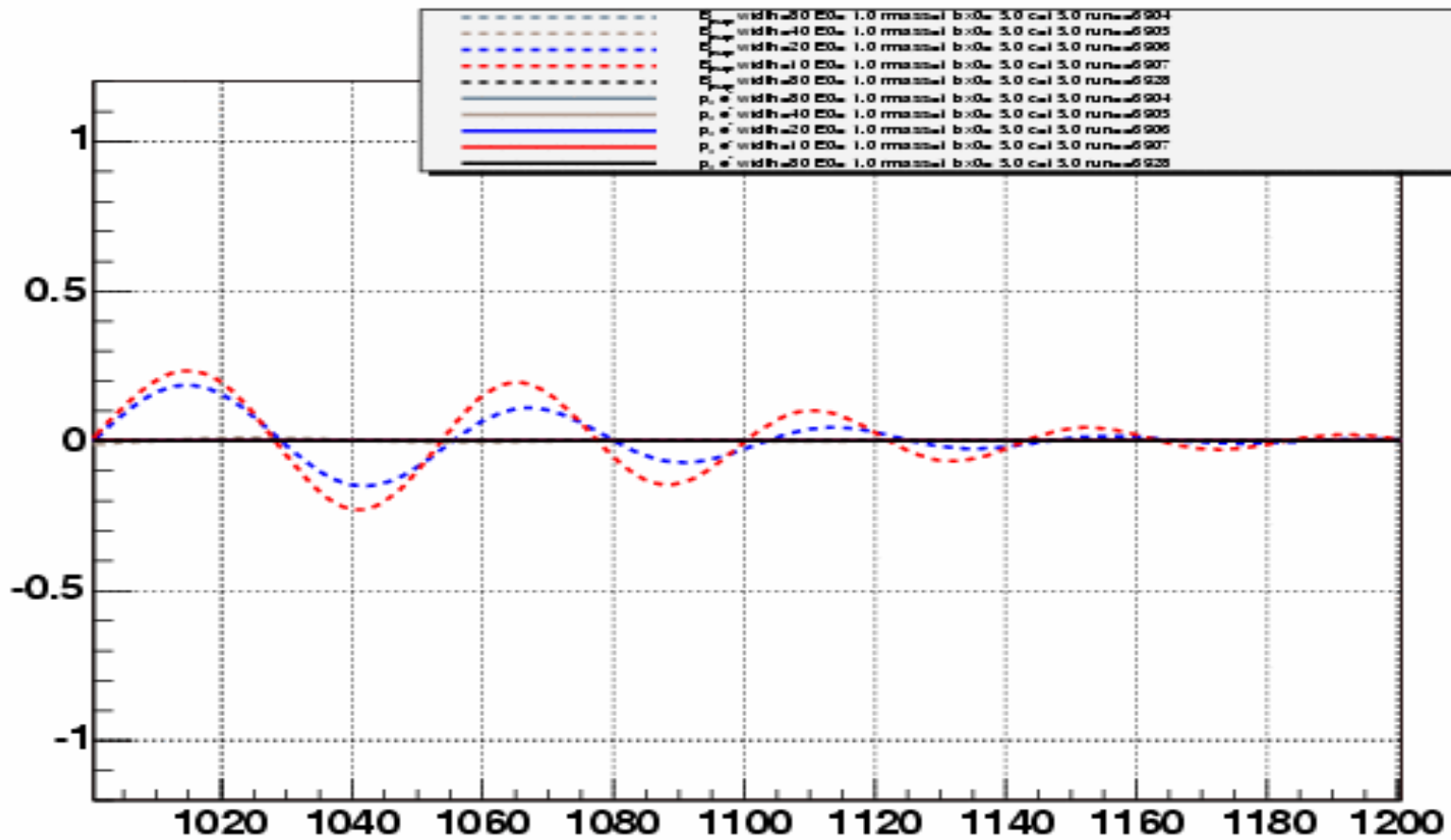
$$R_{\text{mass}} = 1$$



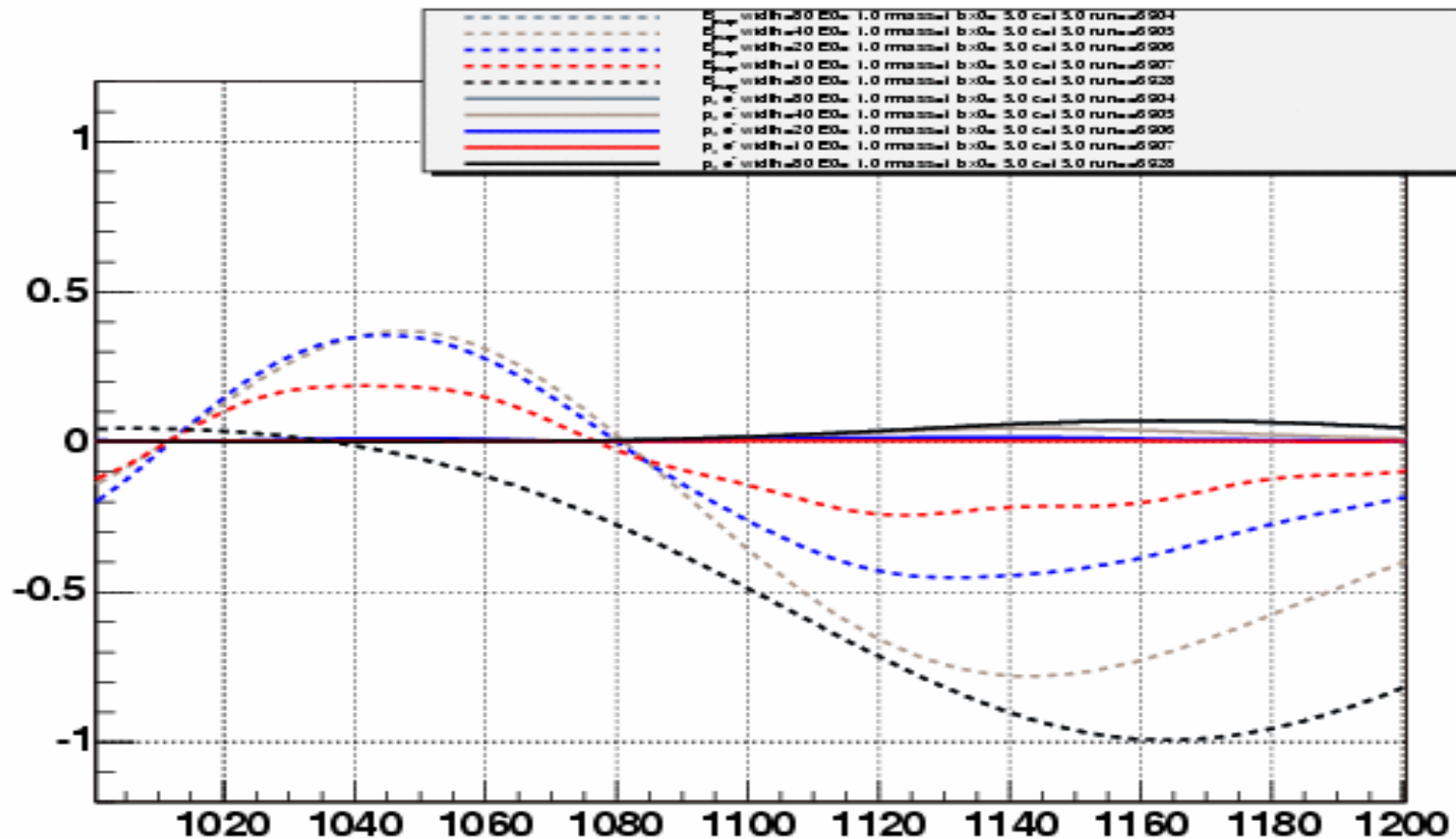
$R_{\text{mass}}=1 \quad T=25 \quad \omega_p^{-1}$ (Zoomed)



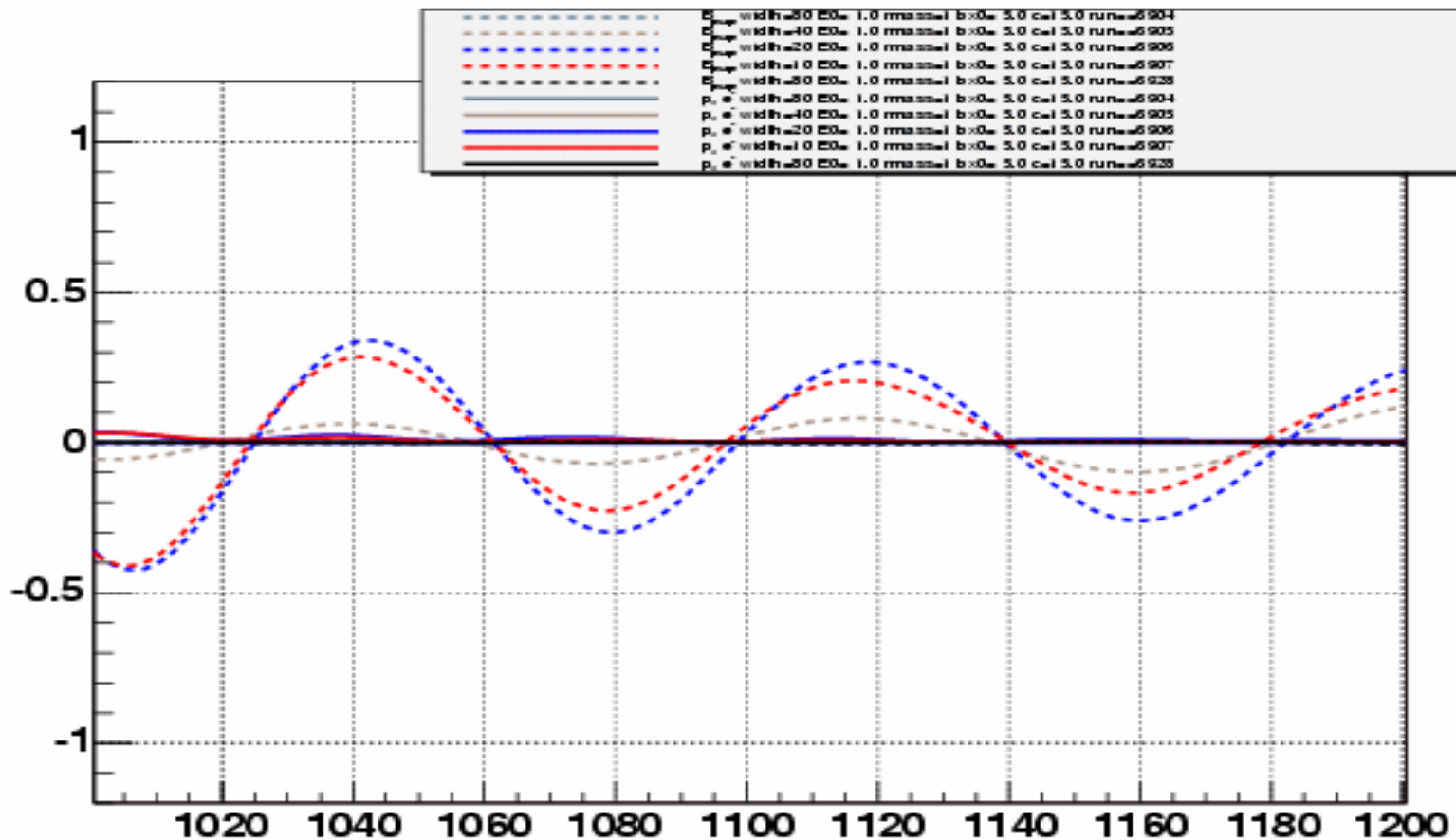
$R_{\text{mass}}=1$ $T=150 \omega_p^{-1}$ (Zoomed)

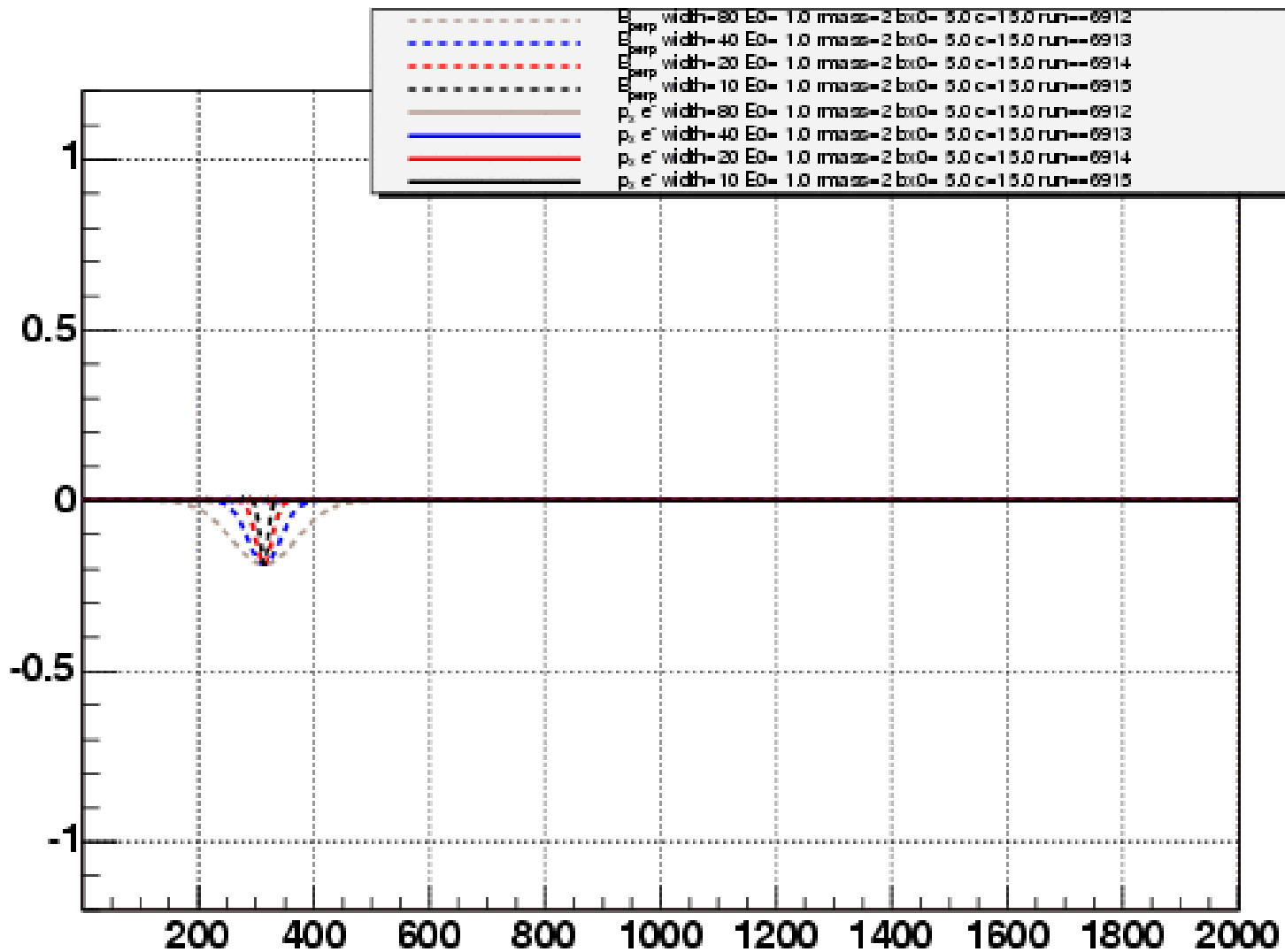


$R_{\text{mass}}=1$ $T=275 \omega_p^{-1}$ (Zoomed)

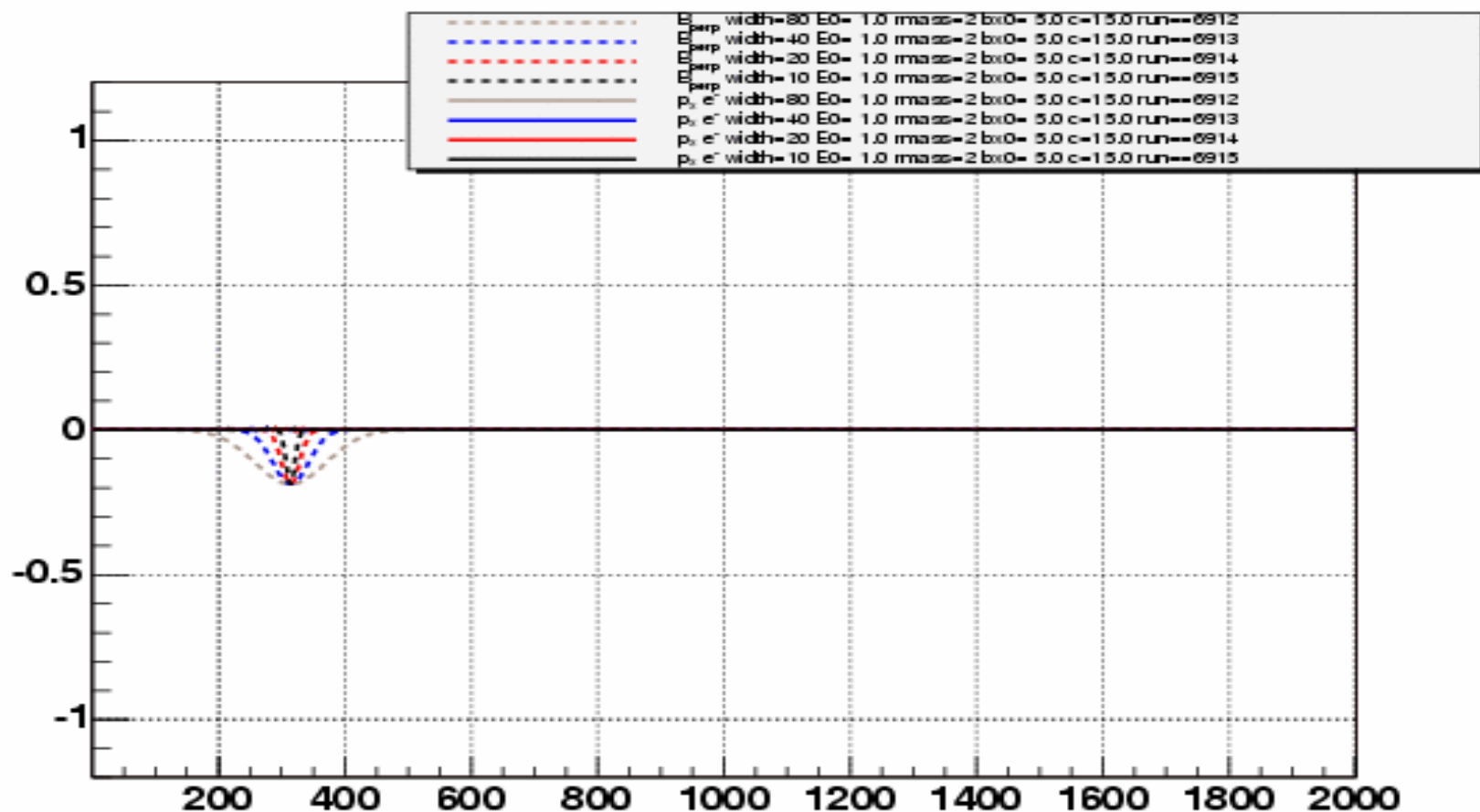


$R_{\text{mass}}=1$ $T=400$ ω_p^{-1} (Zoomed)

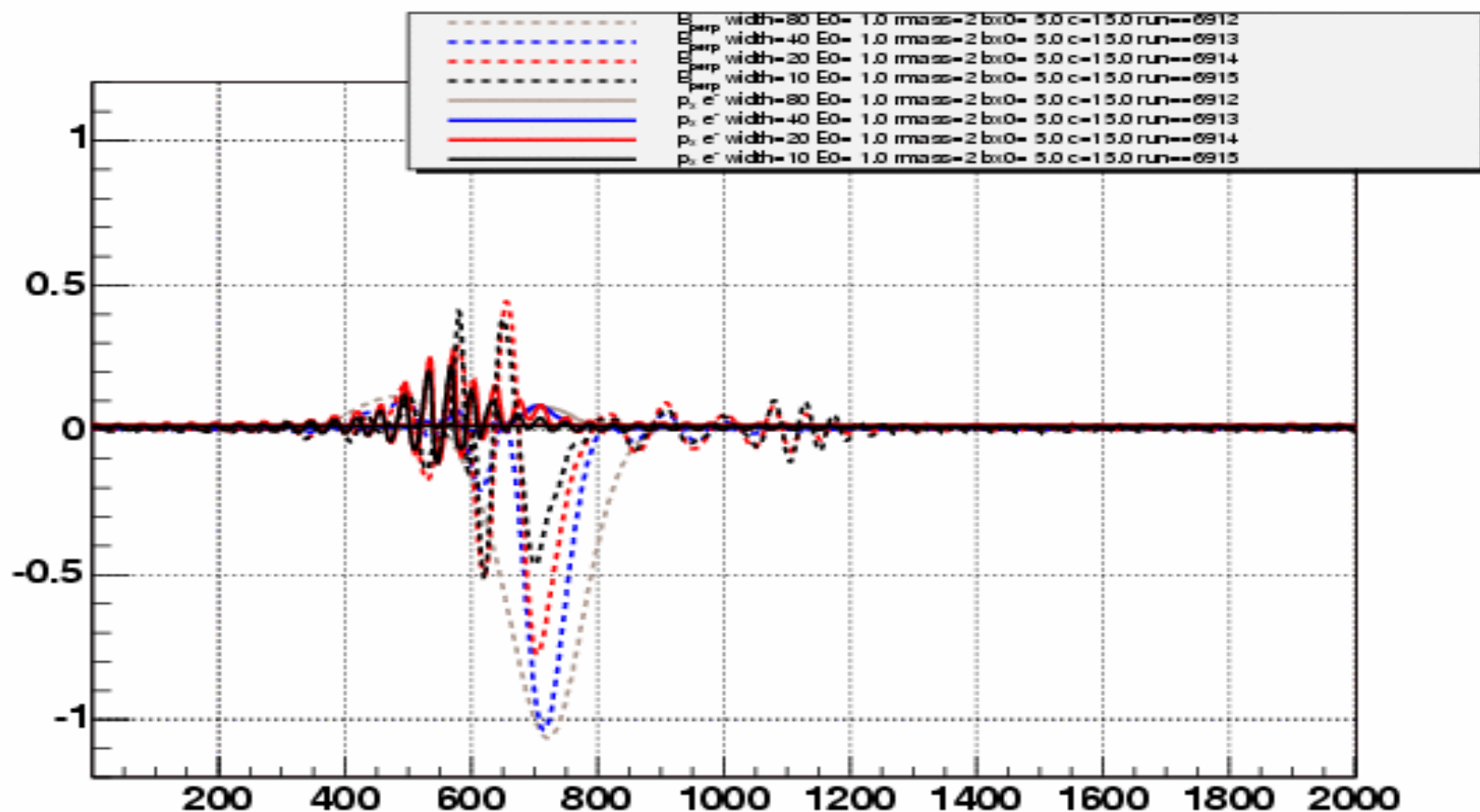




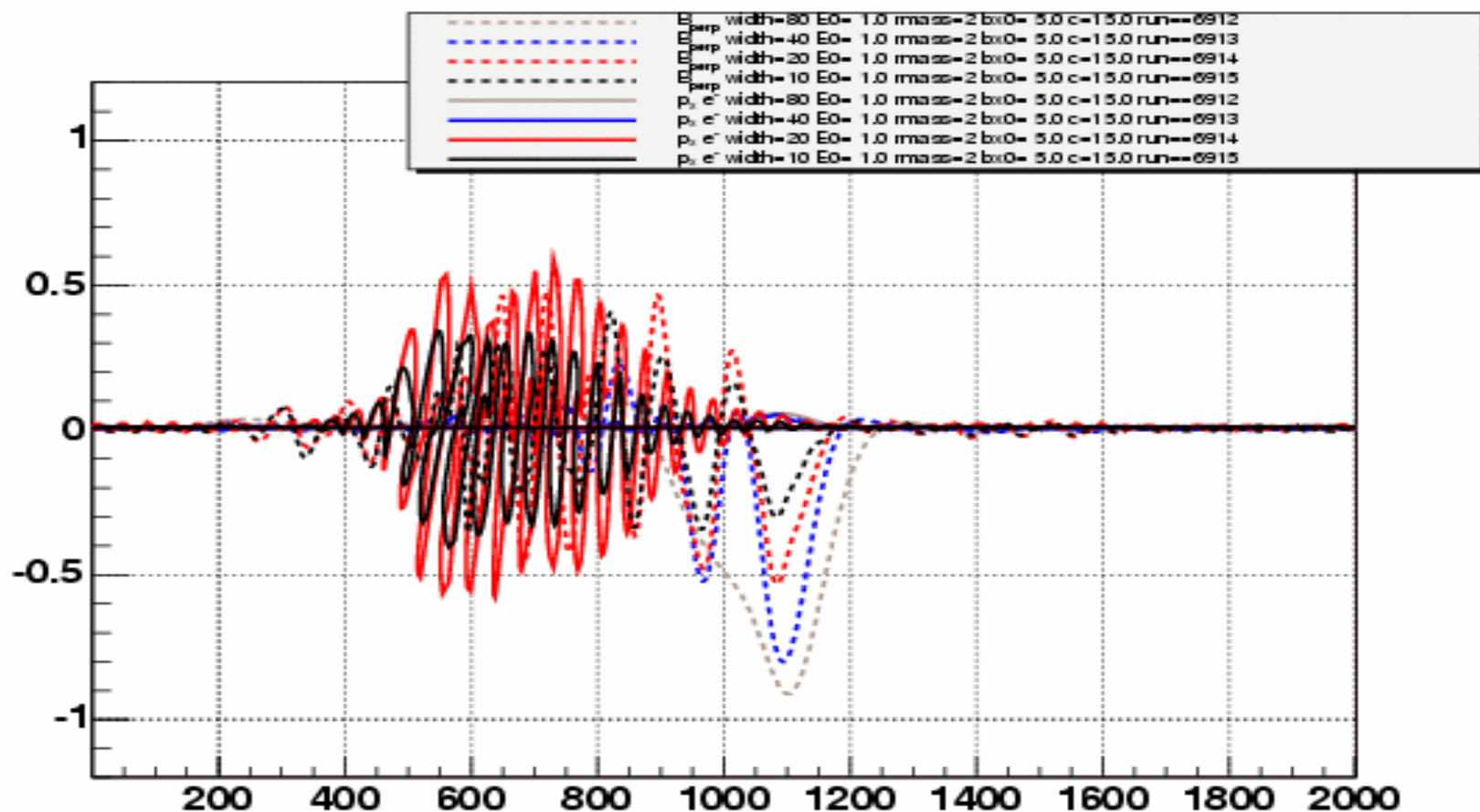
$$R_{\text{mass}} = 2 \quad T = 25 \quad \omega_p^{-1}$$



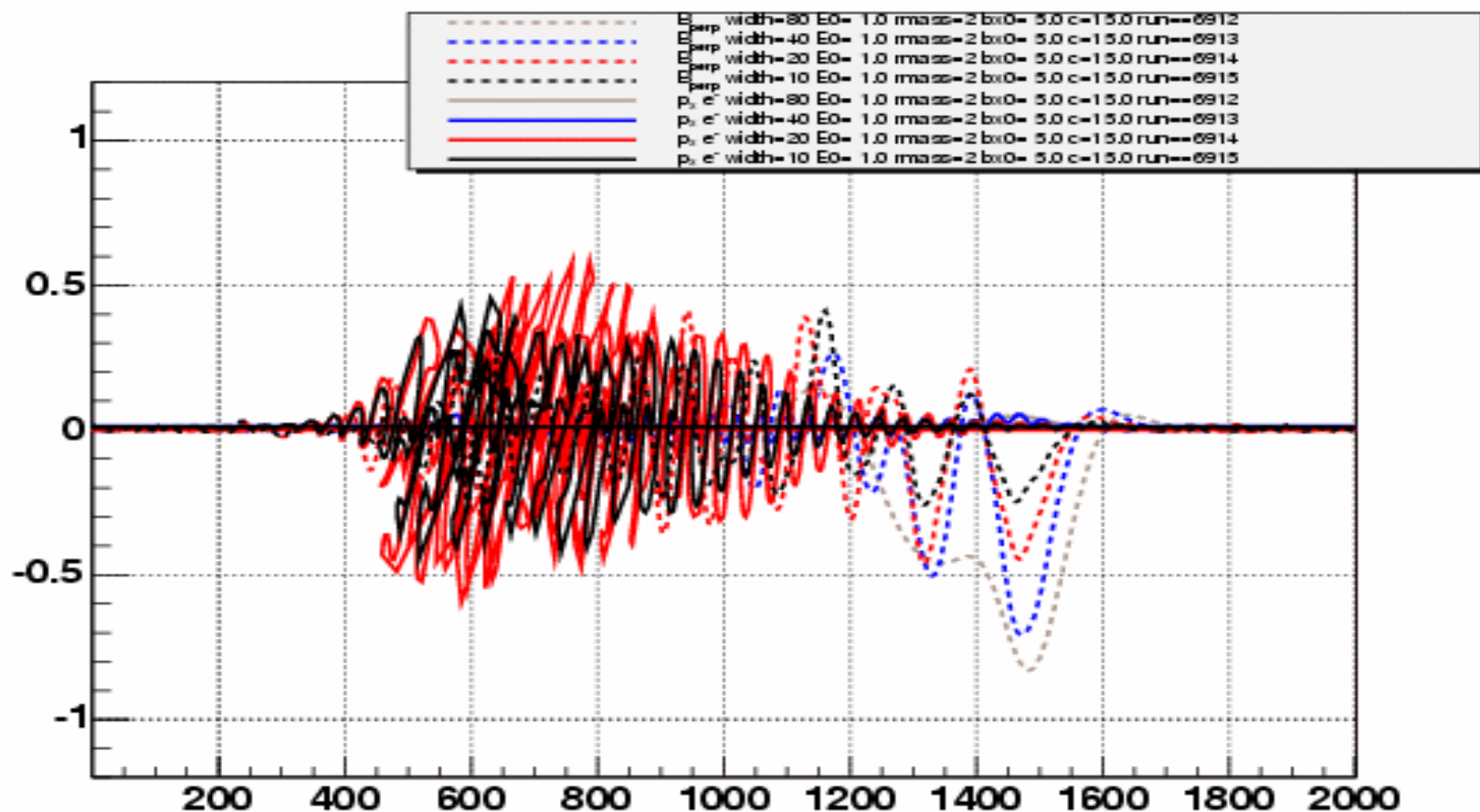
$$R_{\text{mass}} = 2 \quad T = 150 \quad \omega_p^{-1}$$

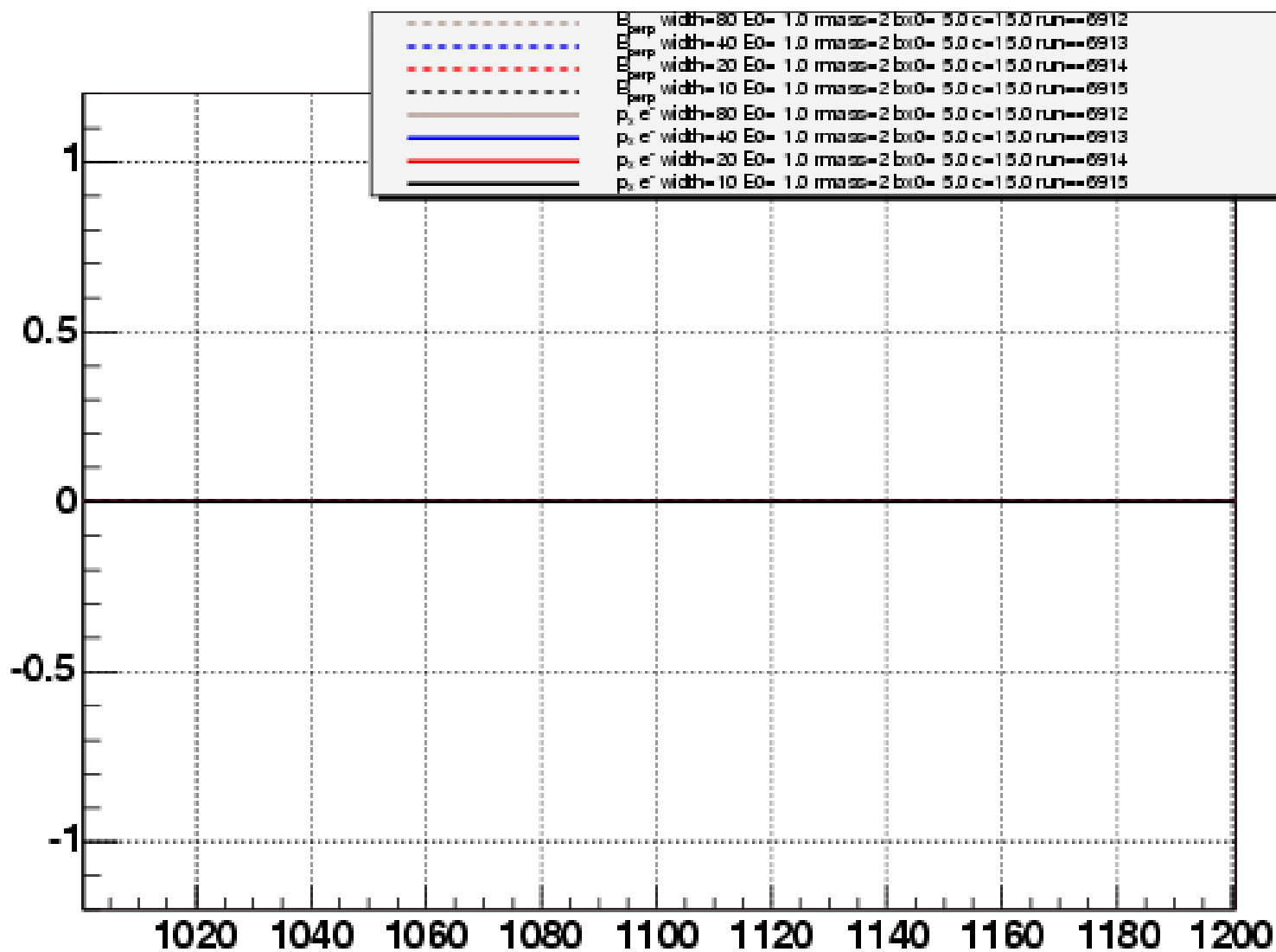


$$R_{\text{mass}} = 2 \quad T = 275 \quad \omega_p^{-1}$$

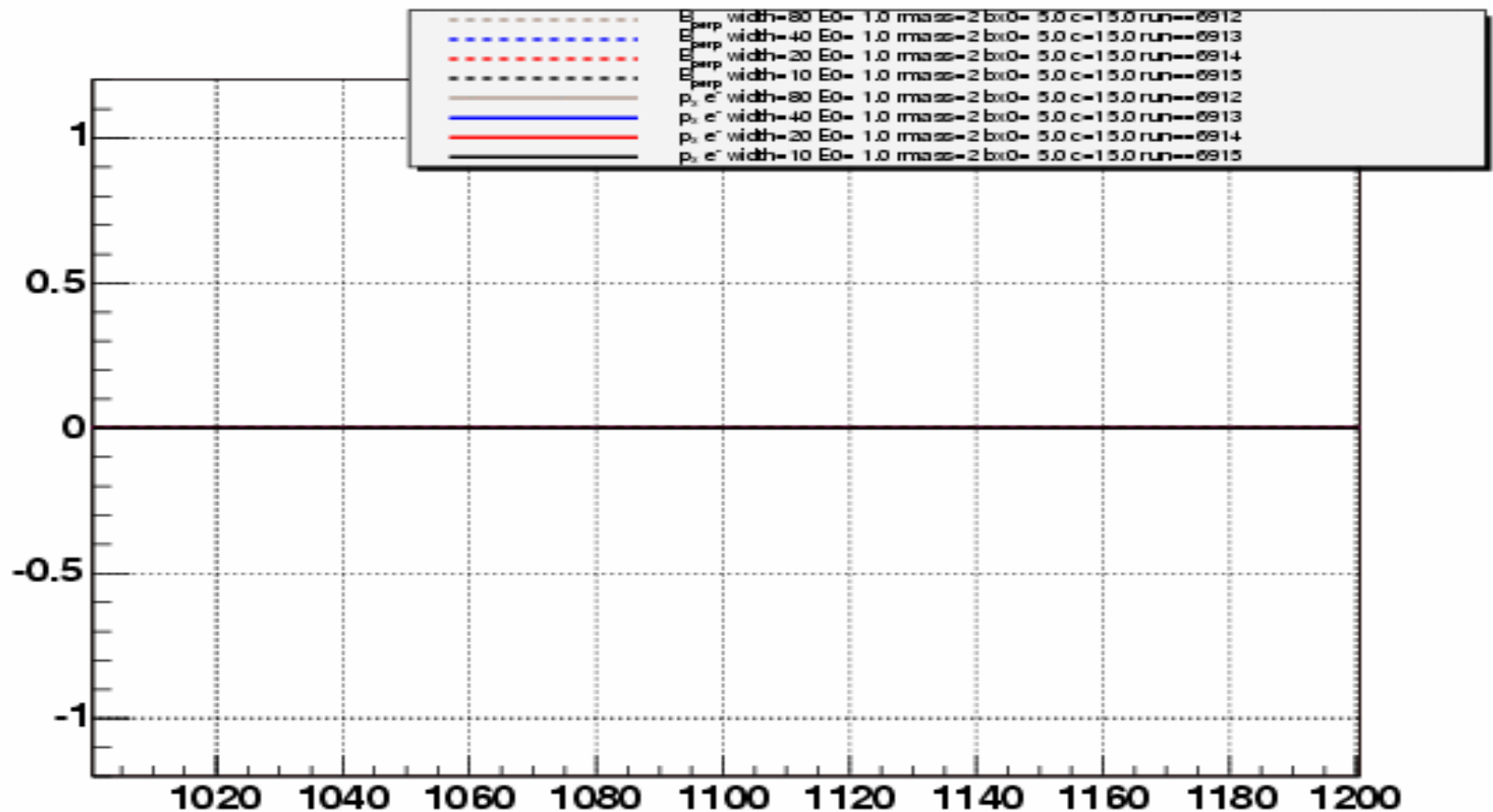


$$R_{\text{mass}} = 2 \quad T = 400 \quad \omega_p^{-1}$$

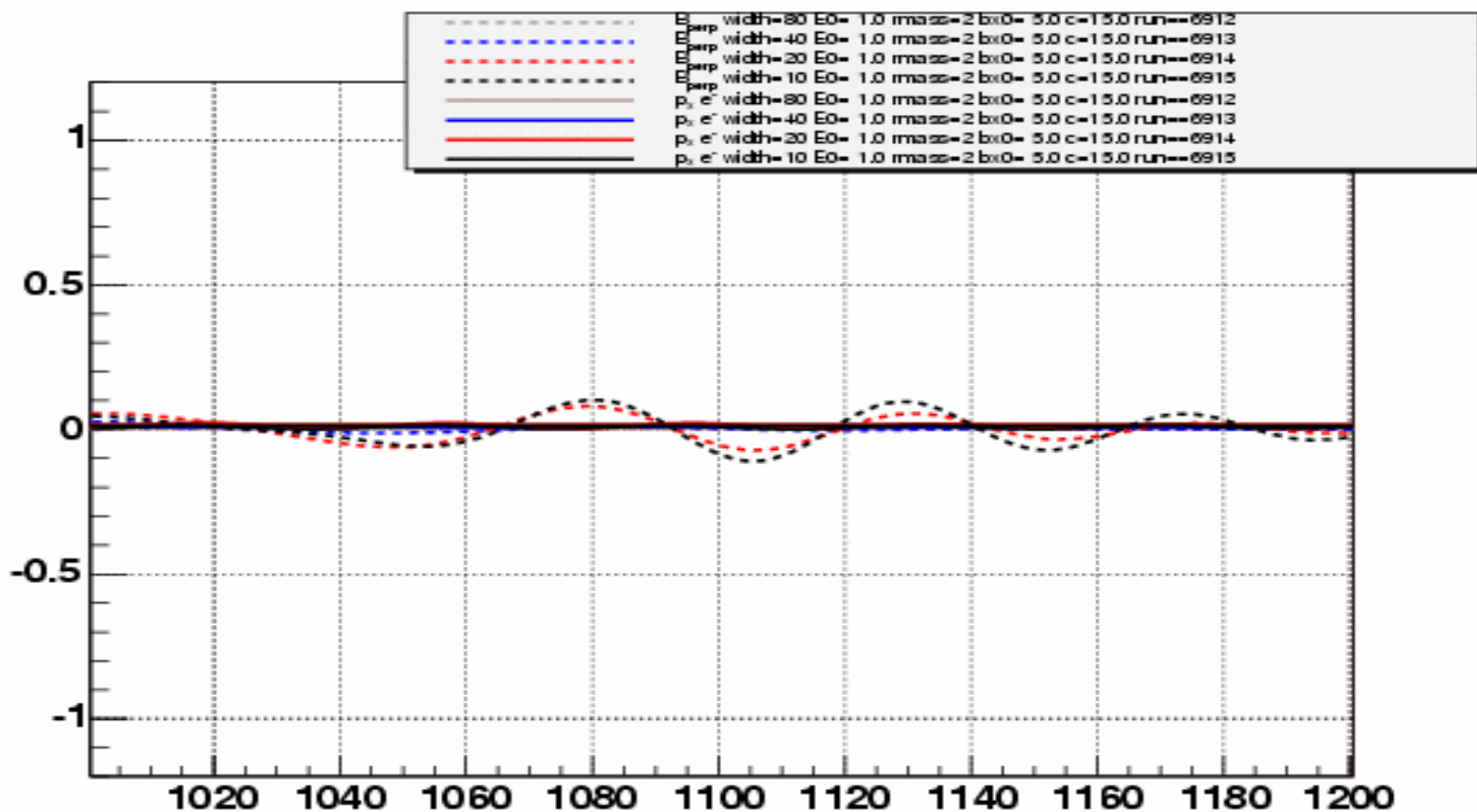




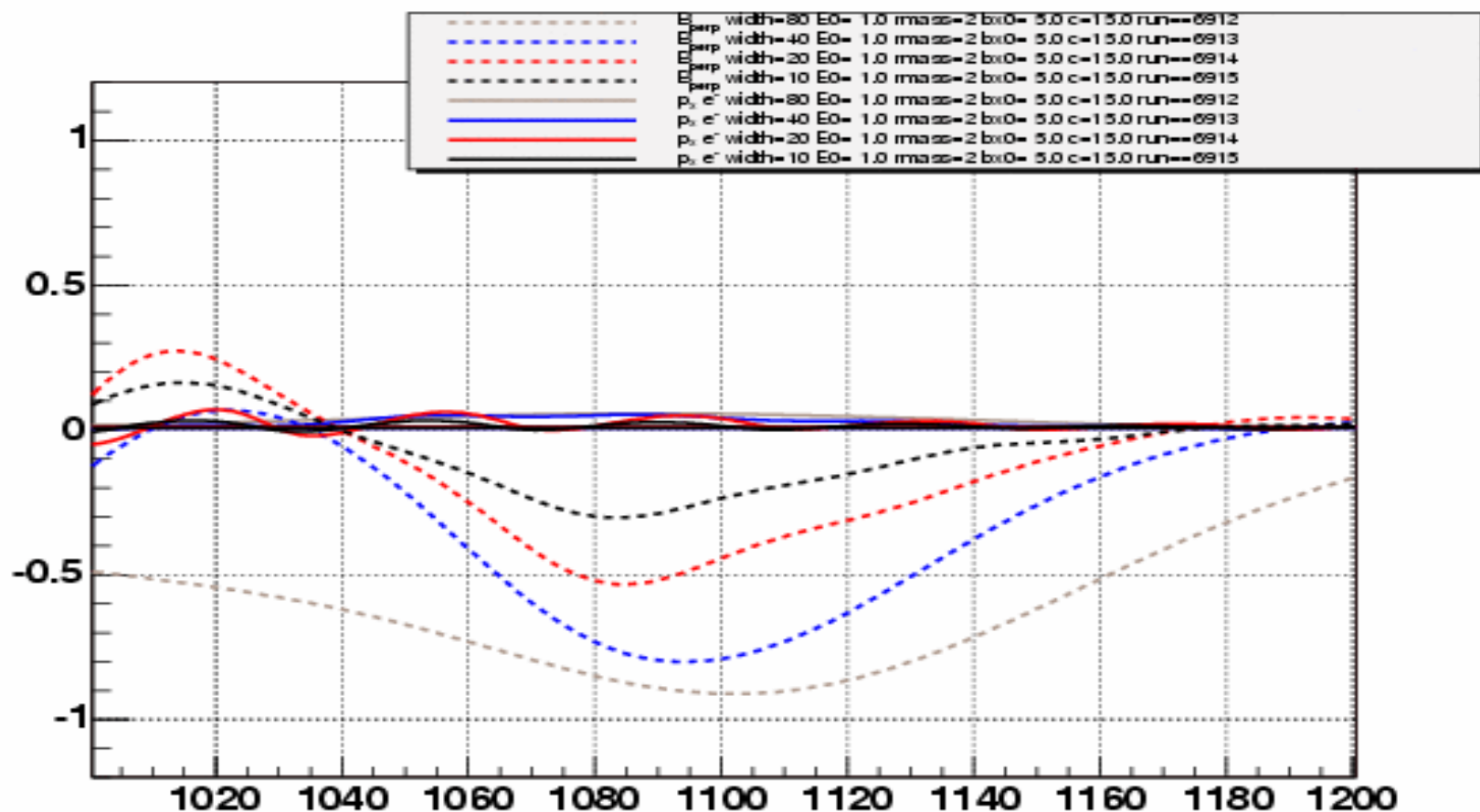
$R_{\text{mass}}=2 \quad T=25 \quad \omega_p^{-1}$ (Zoomed)



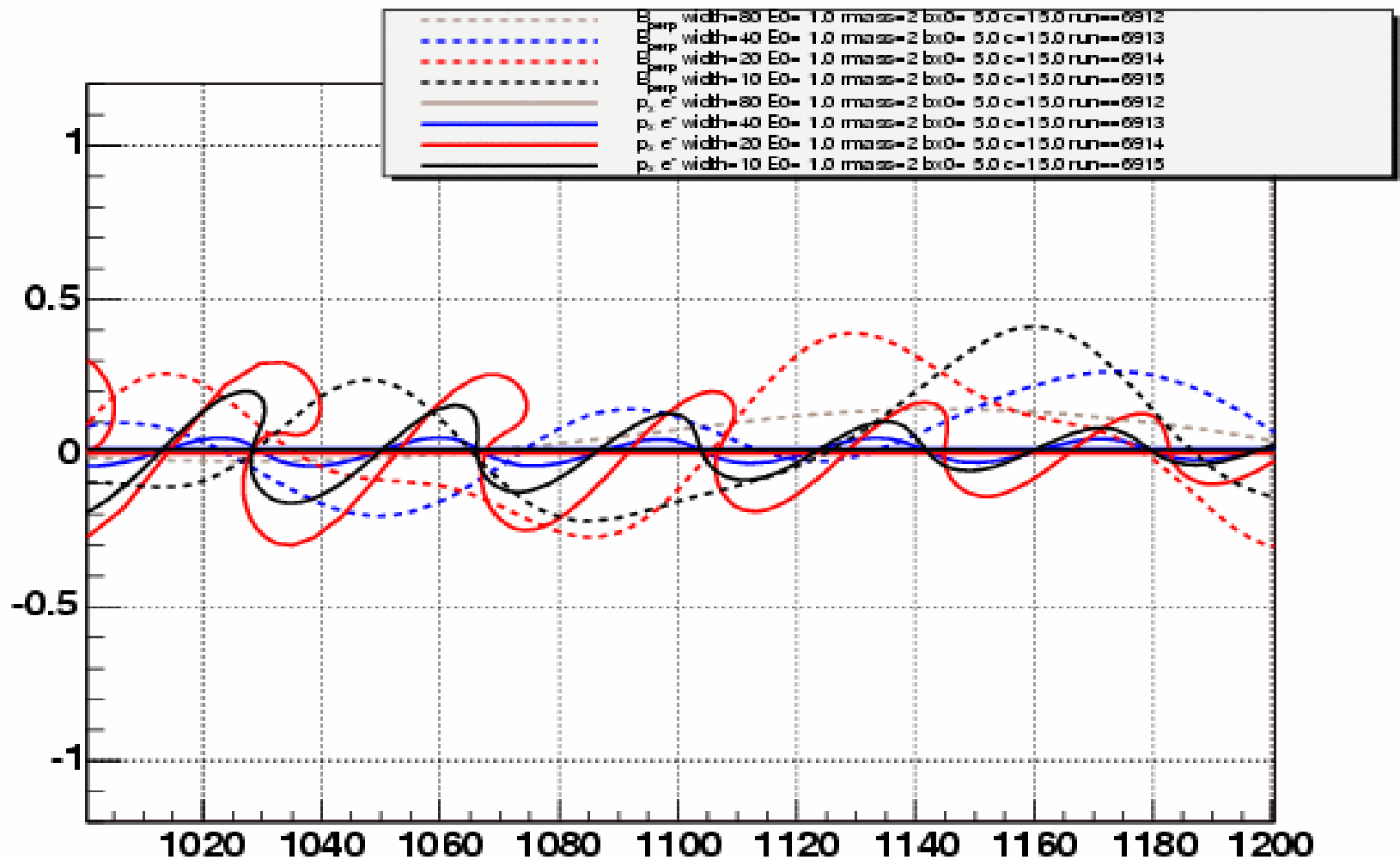
$R_{\text{mass}}=2 \quad T=150 \quad \omega_p^{-1}$ (Zoomed)

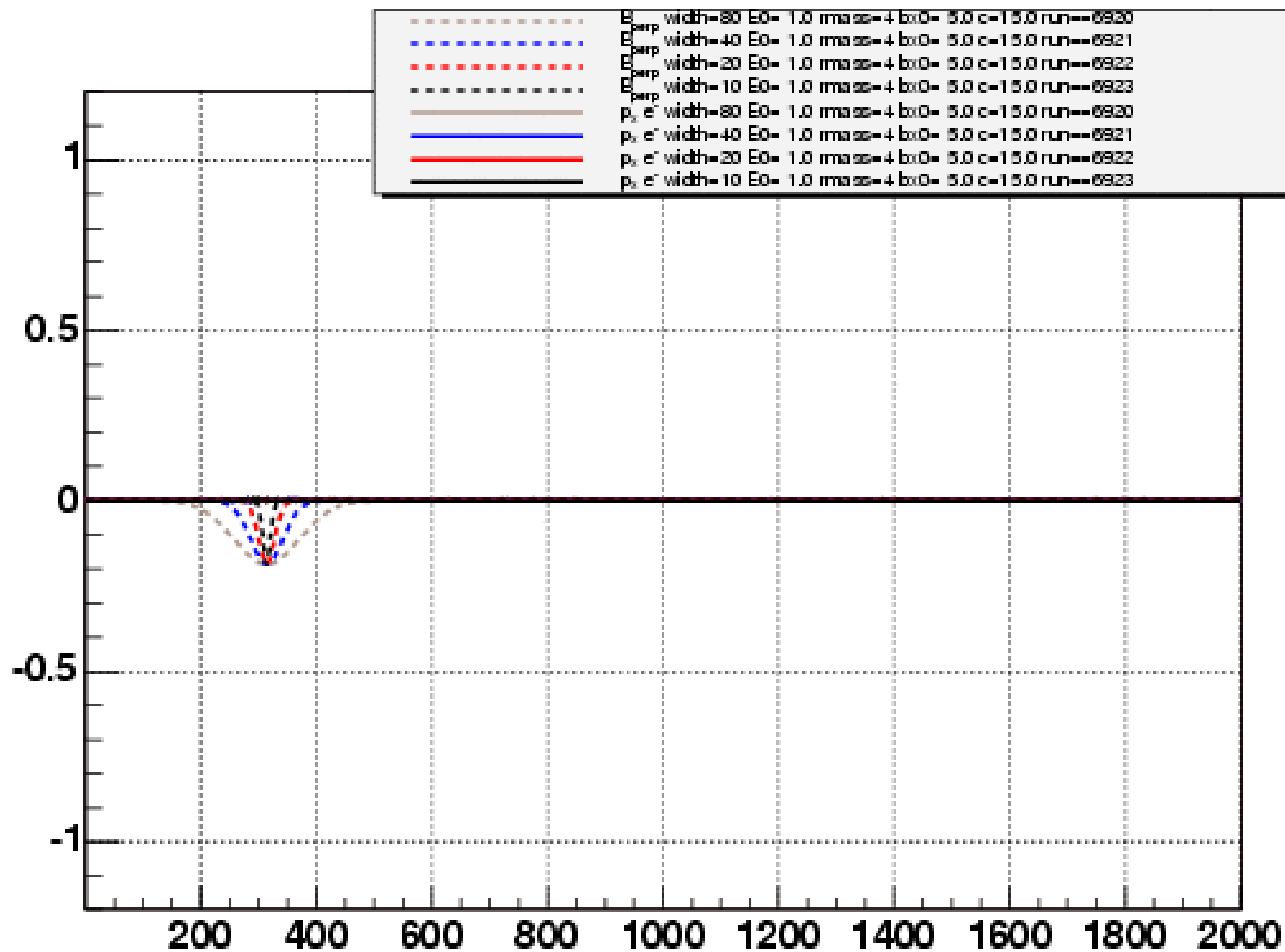


$R_{\text{mass}}=2 \quad T=275 \quad \omega_p^{-1}$ (Zoomed)

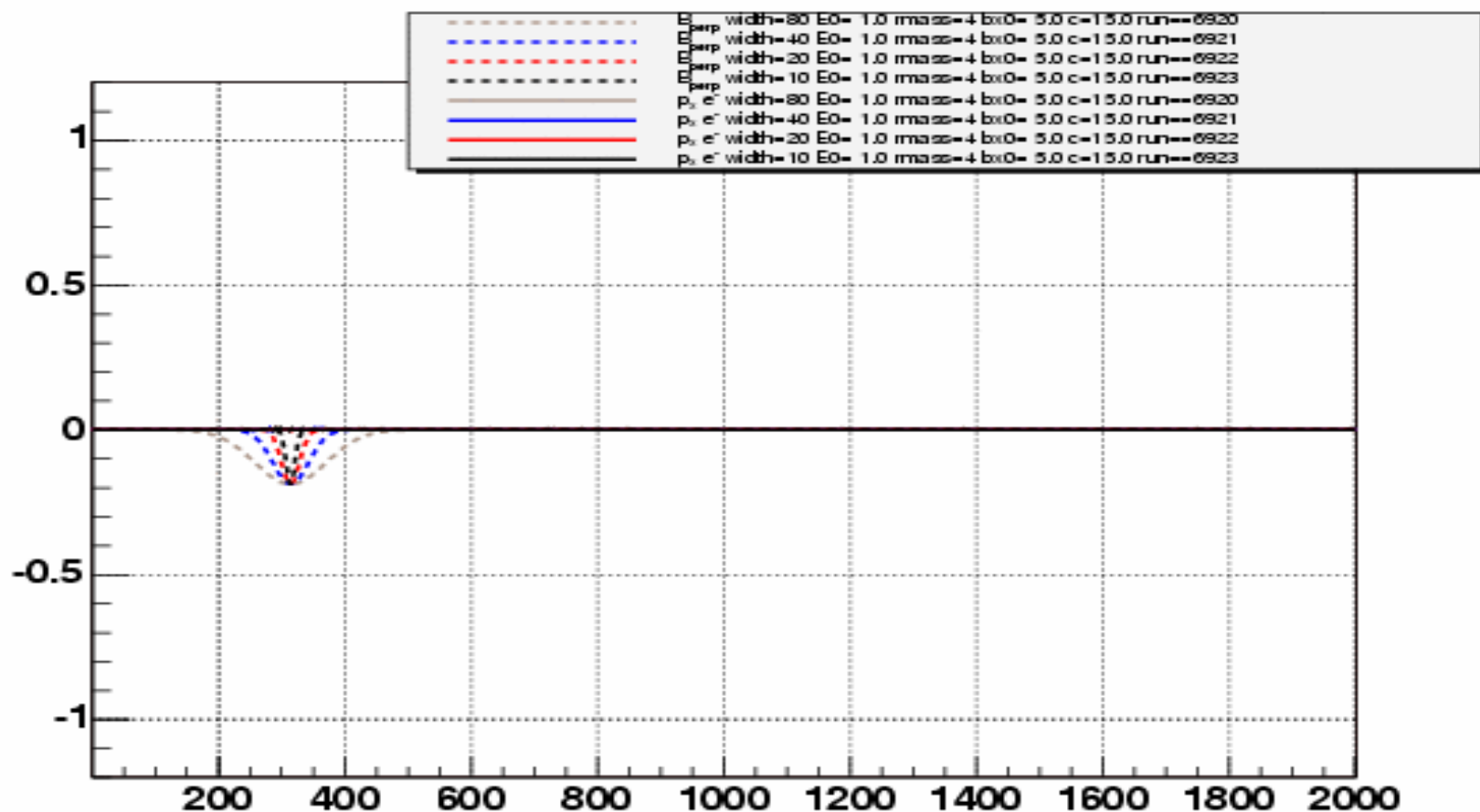


$R_{\text{mass}}=2 \quad T=400 \quad \omega_p^{-1}$ (Zoomed)

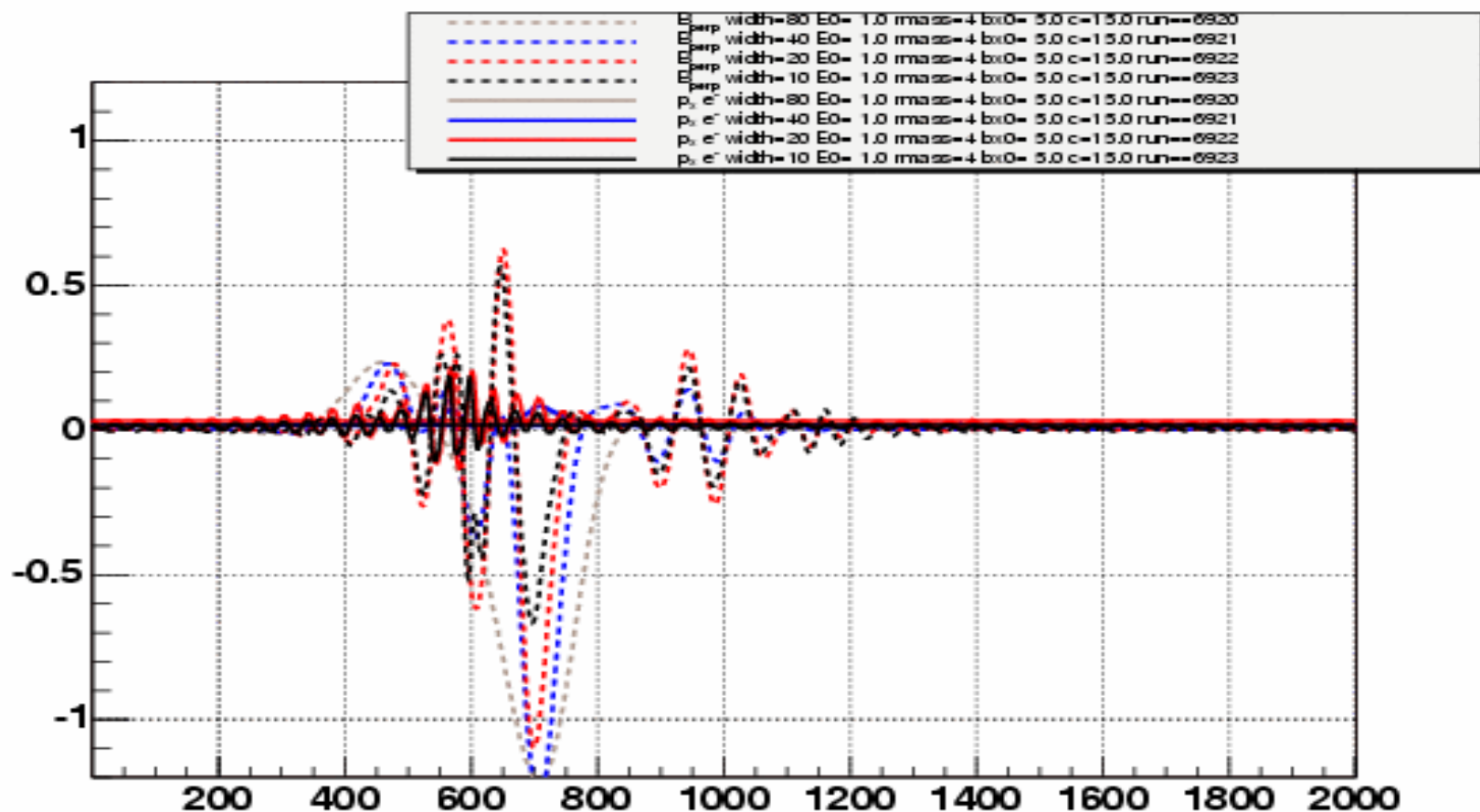




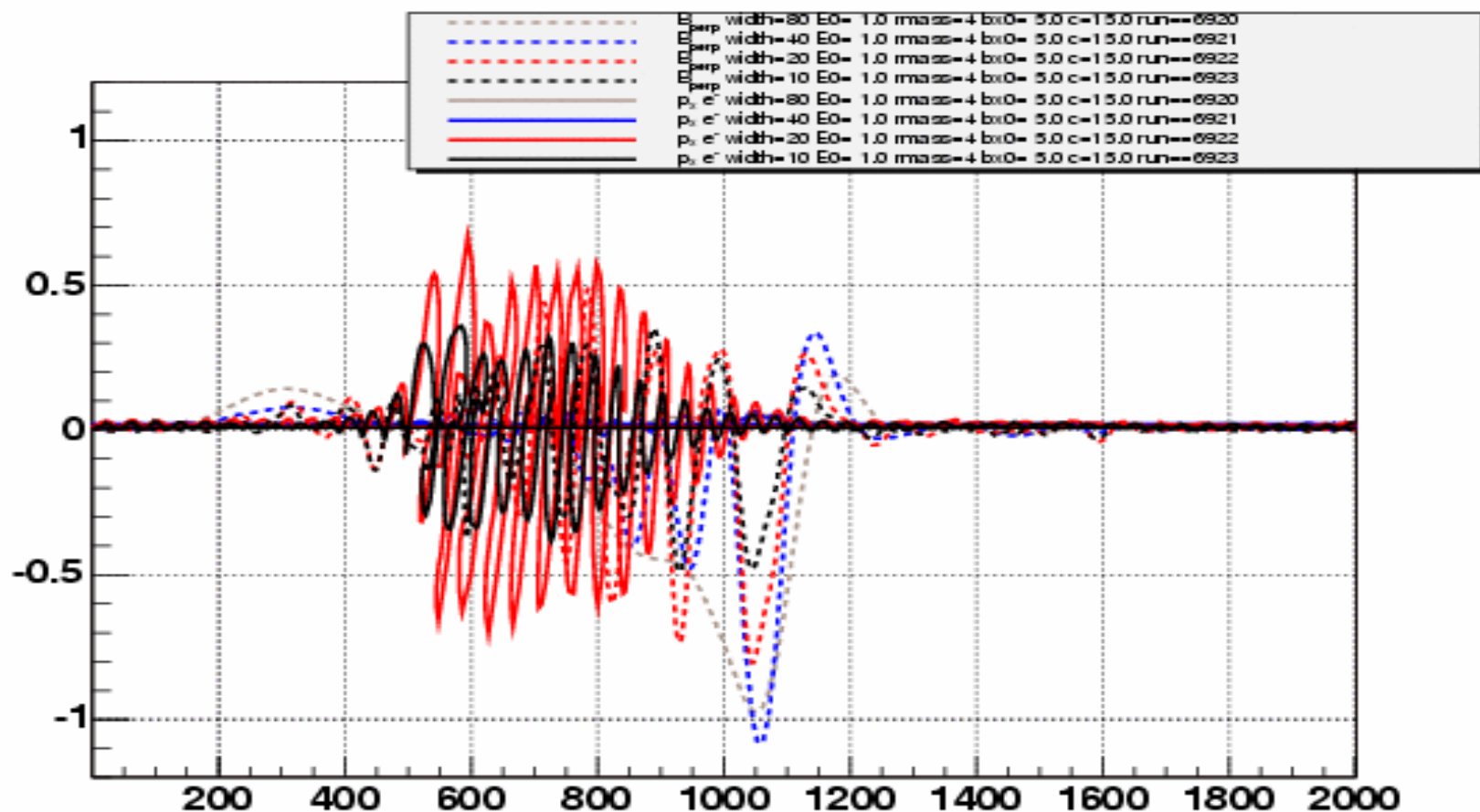
$$R_{\text{mass}}=4 \quad T=25 \quad \omega_p^{-1}$$



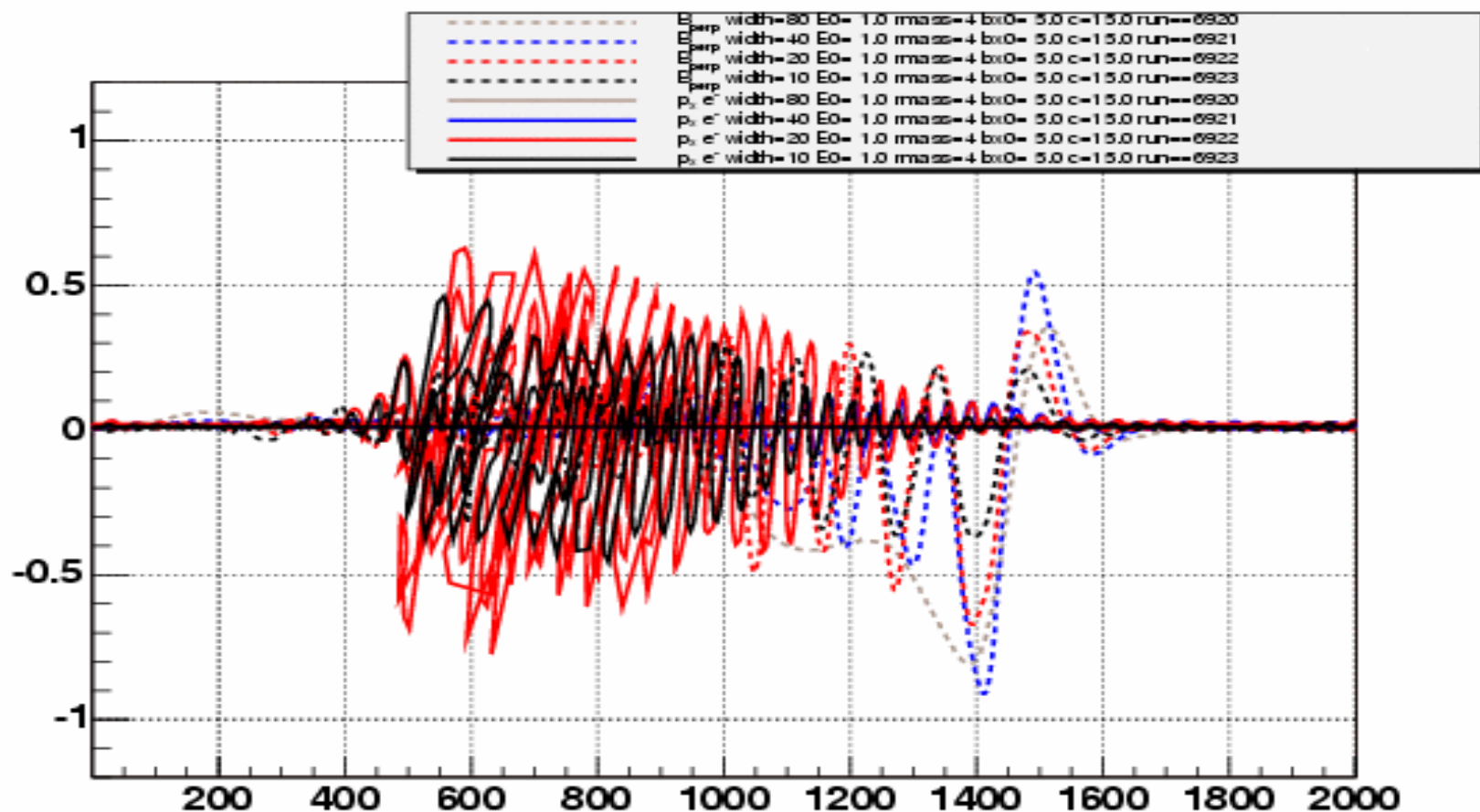
$$R_{\text{mass}}=4 \quad T=150 \quad \omega_p^{-1}$$

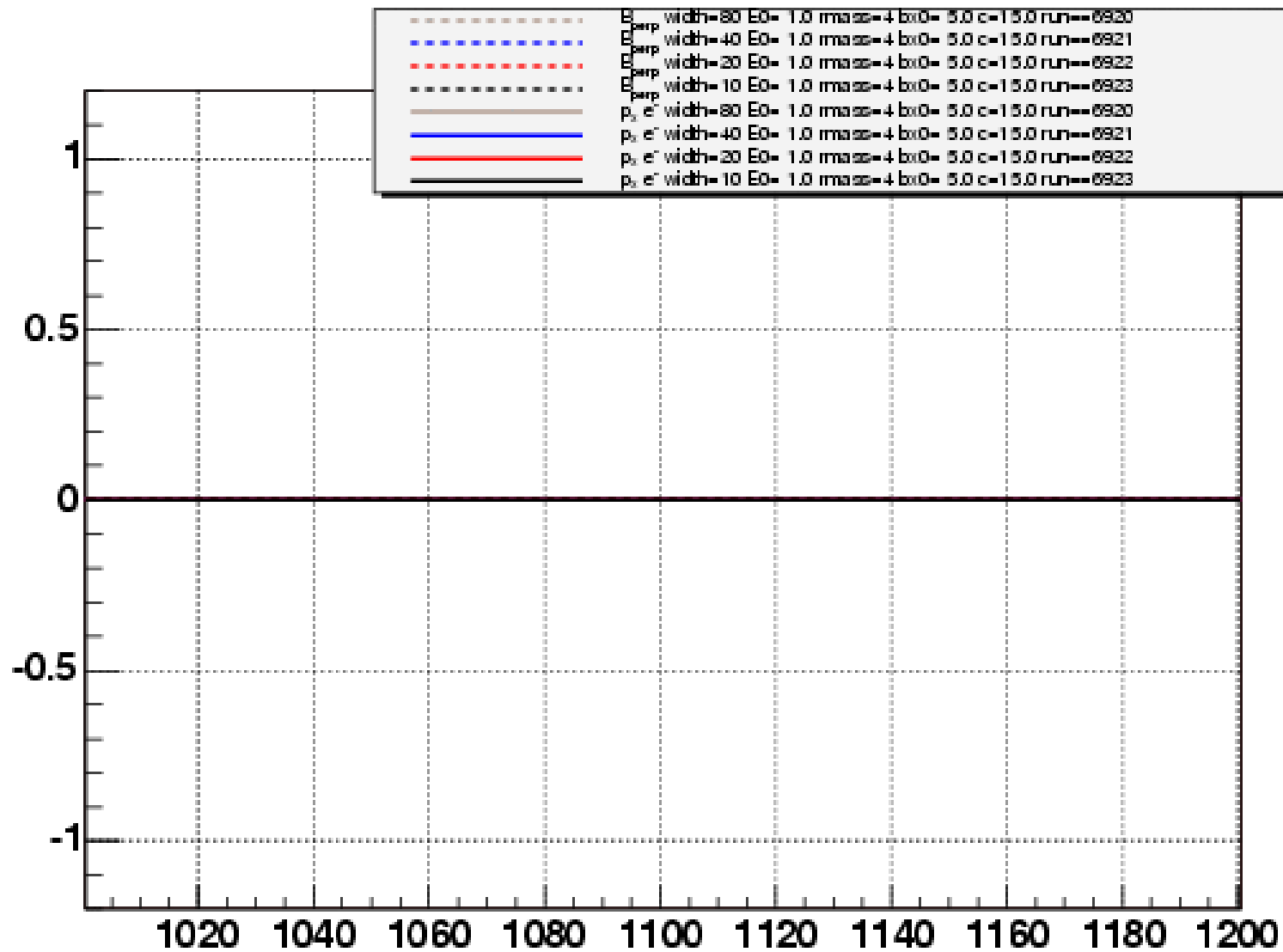


$$R_{\text{mass}} = 4 \quad T = 275 \quad \omega_p^{-1}$$

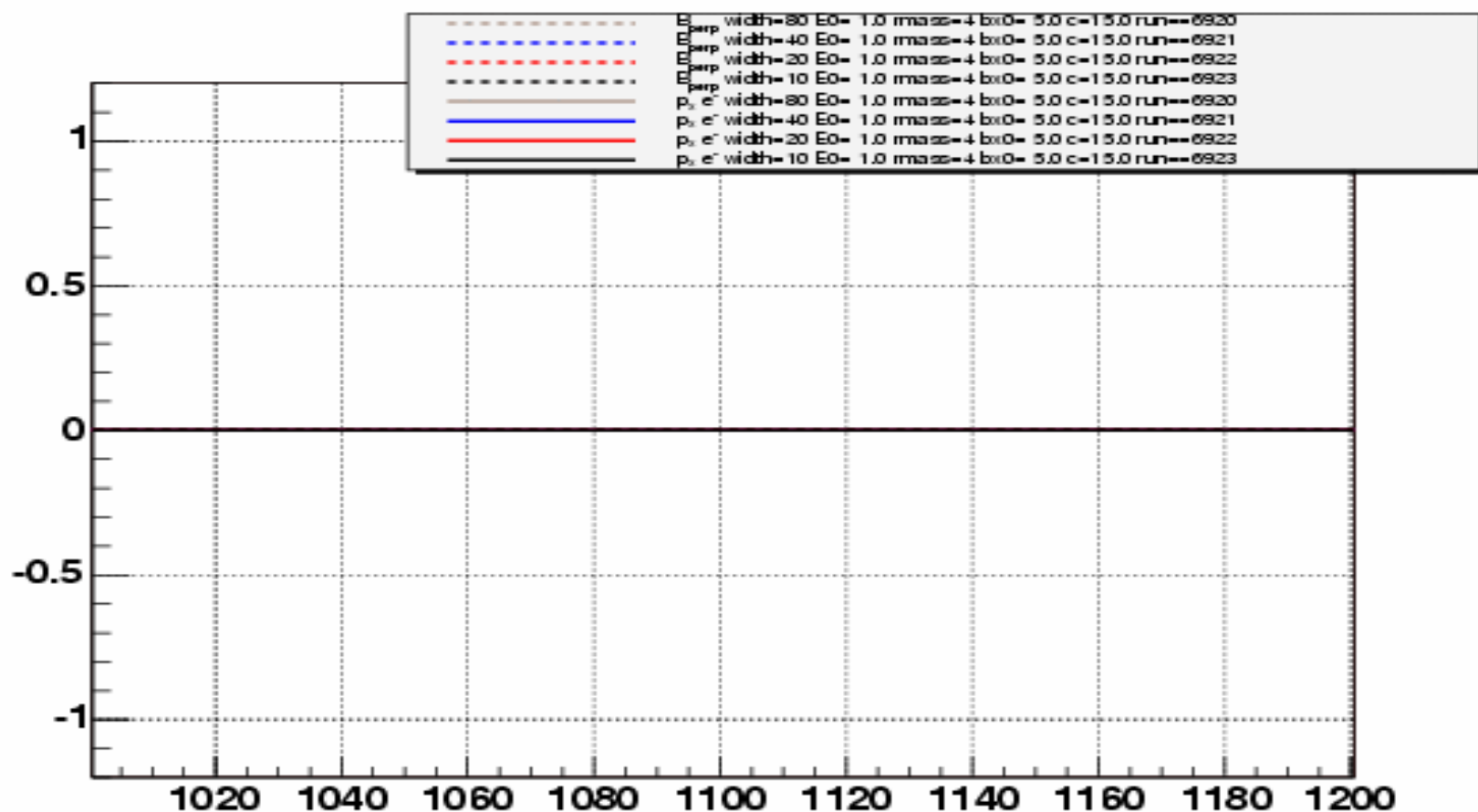


$$R_{\text{mass}}=4 \quad T=400 \quad \omega_p^{-1}$$

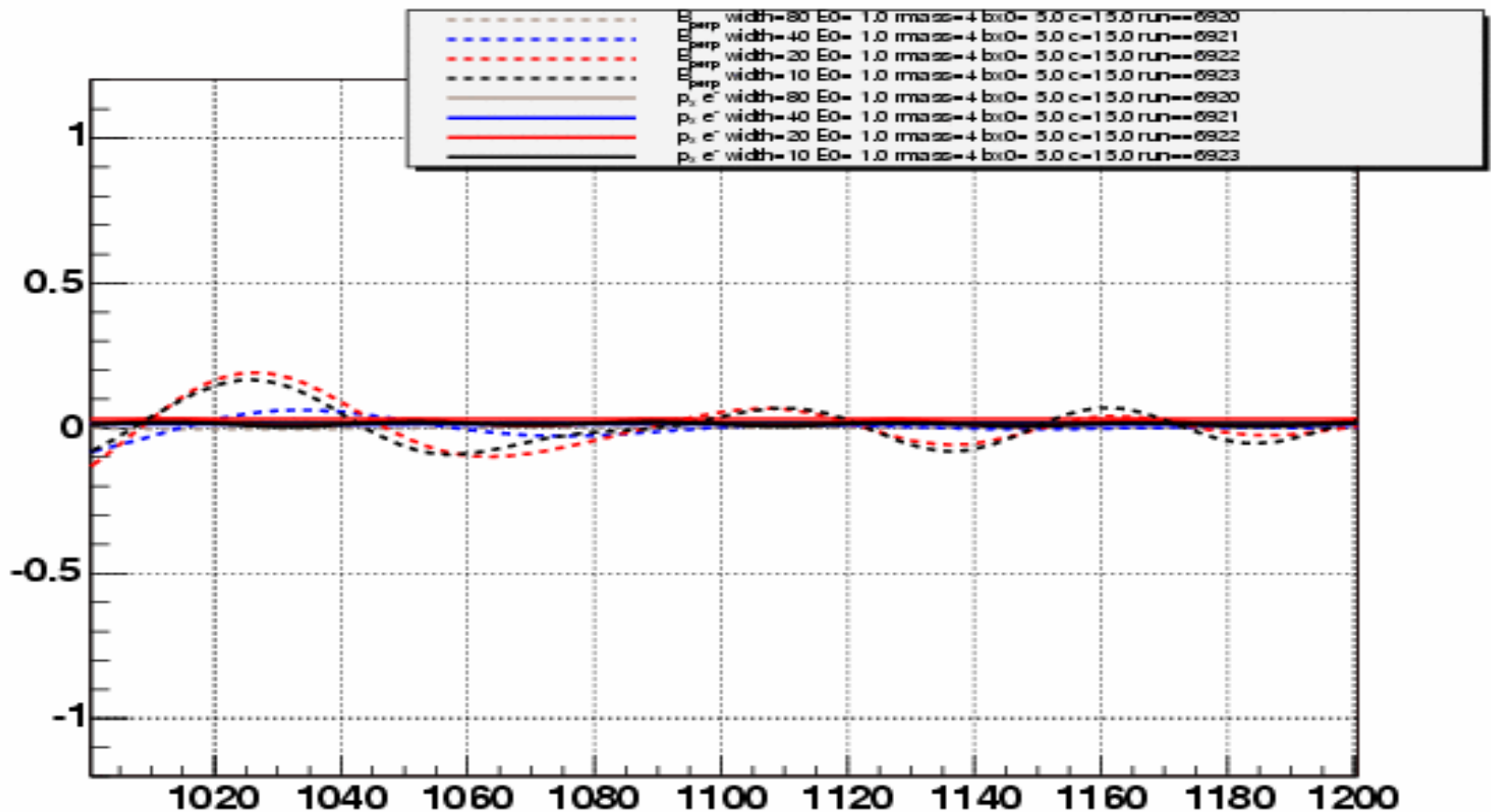




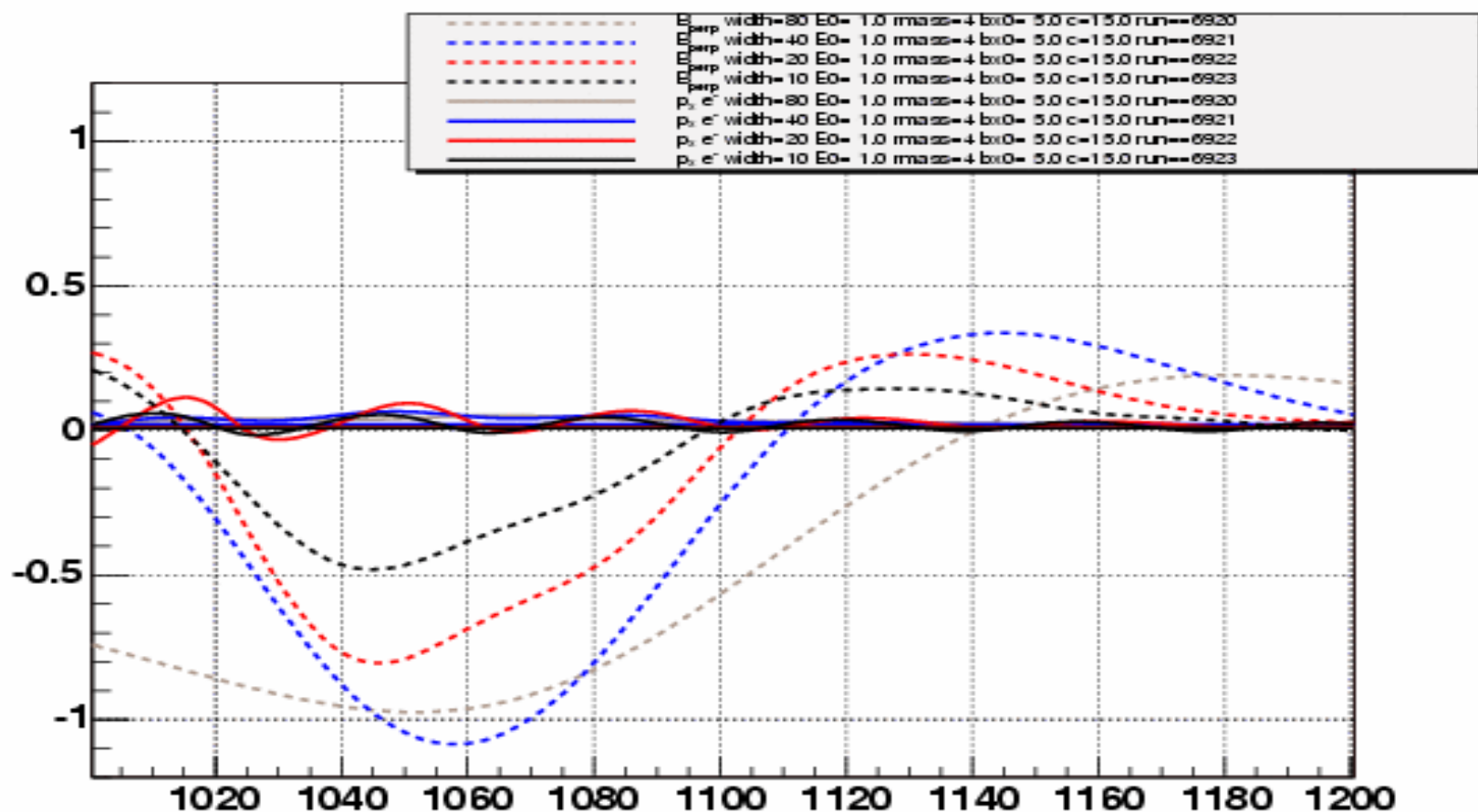
$R_{\text{mass}}=4 \quad T=25 \quad \omega_p^{-1}$ (Zoomed)



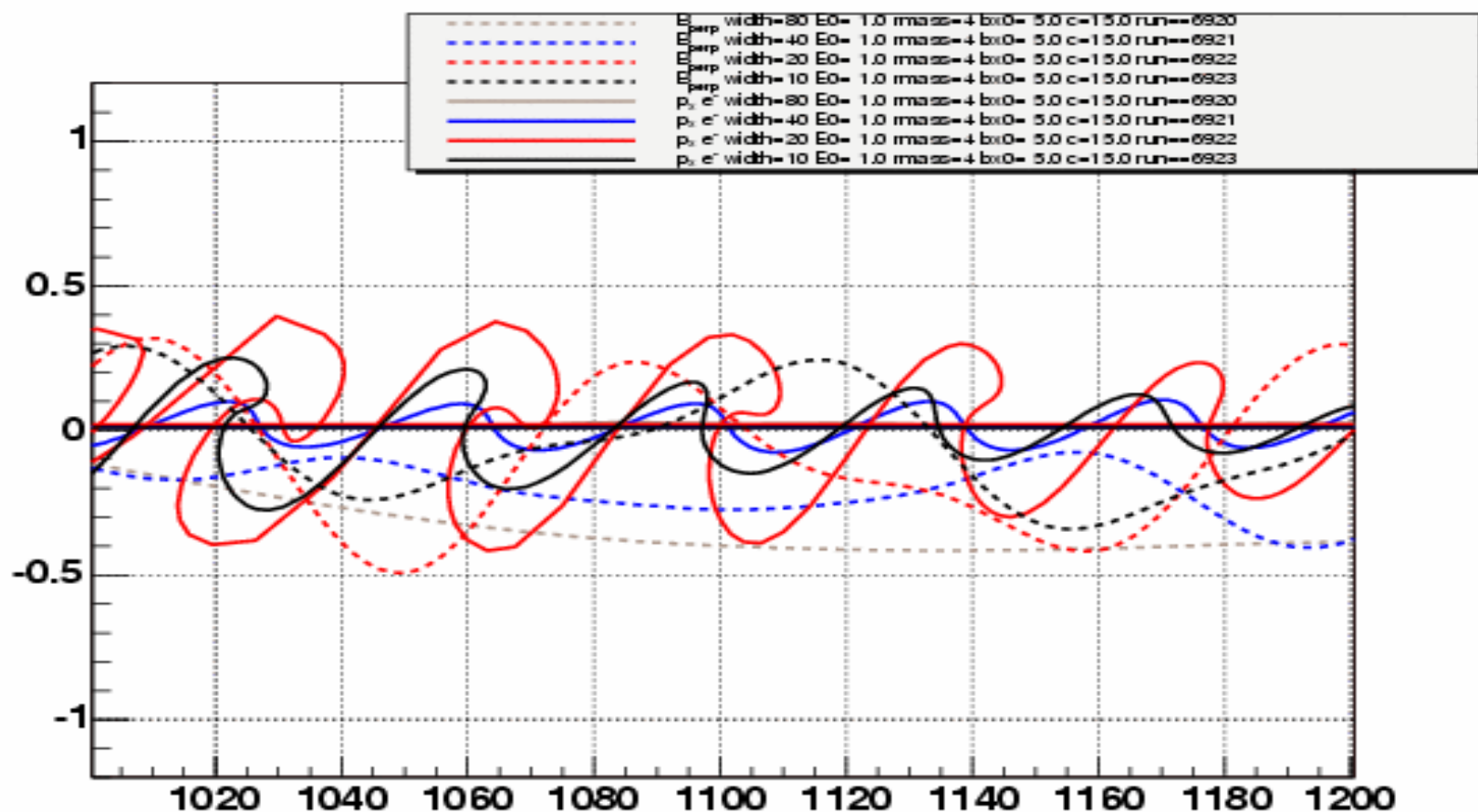
$R_{\text{mass}}=4 \quad T=150 \quad \omega_p^{-1}$ (Zoomed)



$R_{\text{mass}}=4 \quad T=275 \quad \omega_p^{-1}$ (Zoomed)

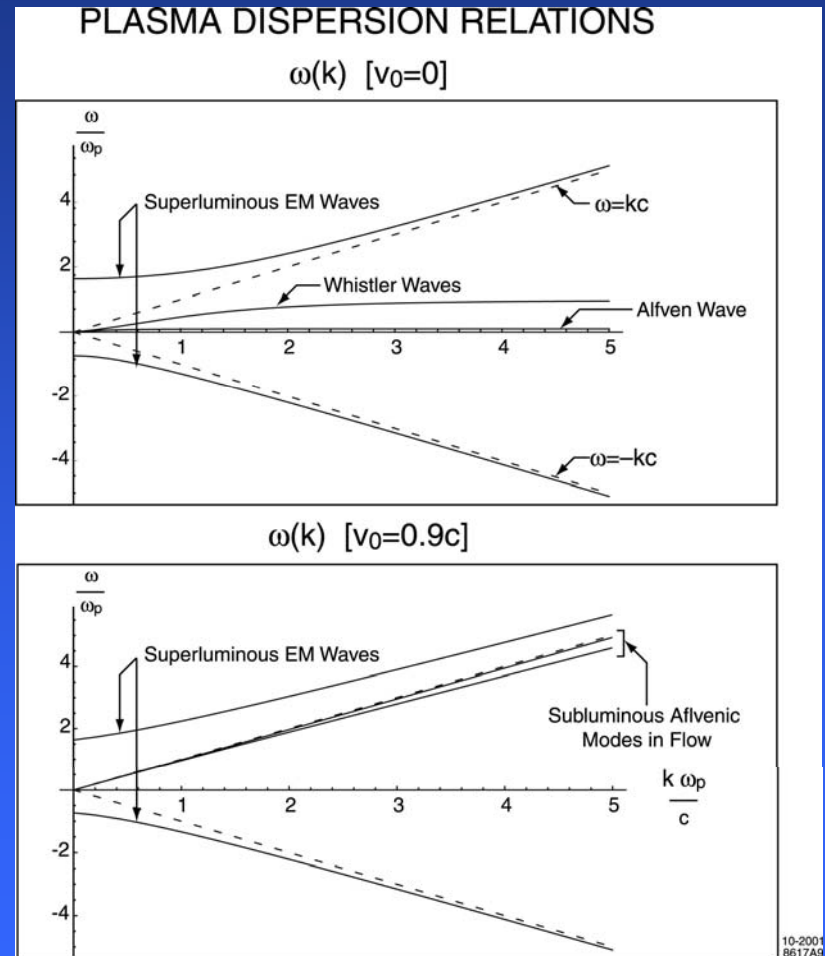


$R_{\text{mass}}=4 \quad T=400 \quad \omega_p^{-1}$ (Zoomed)



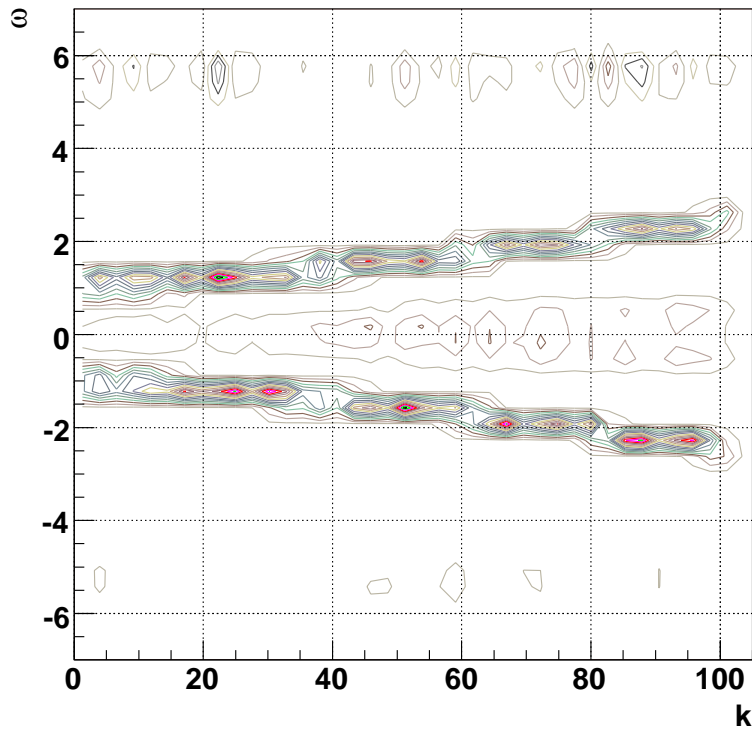
Dispersion Relation

- Study dispersion relation to ensure what we are observing is Alfvén branch:
- Excitation increases sharply with shorter pulse width (higher frequency).
- May be Whistler as opposed to Alfvén branch.
- Work in progress....

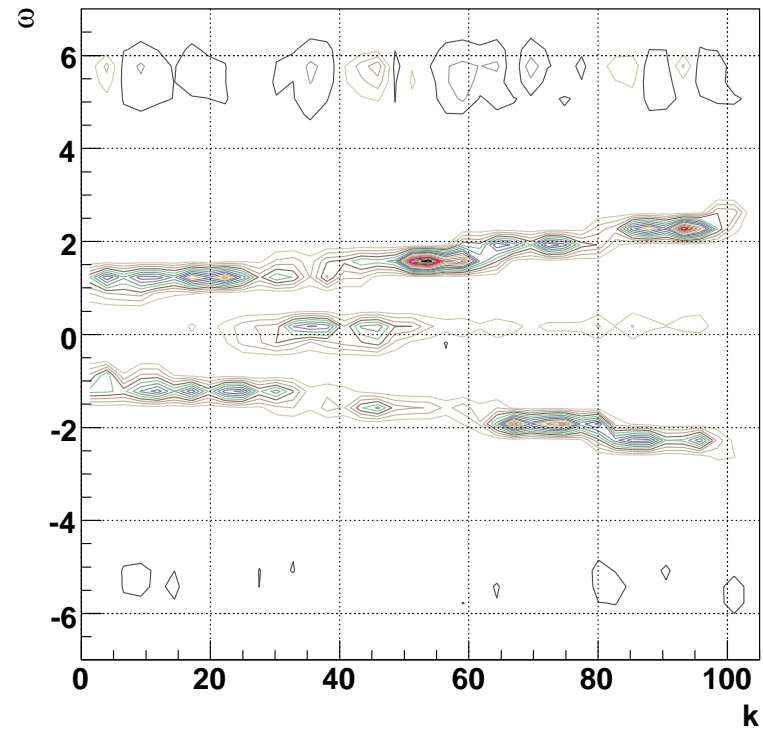


Dispersion Relation

15. 10 7000 7000 1.1 5.0 0.0 0.1 c 4.0 c 200 500 12 1000 XXX 6.0 0.0

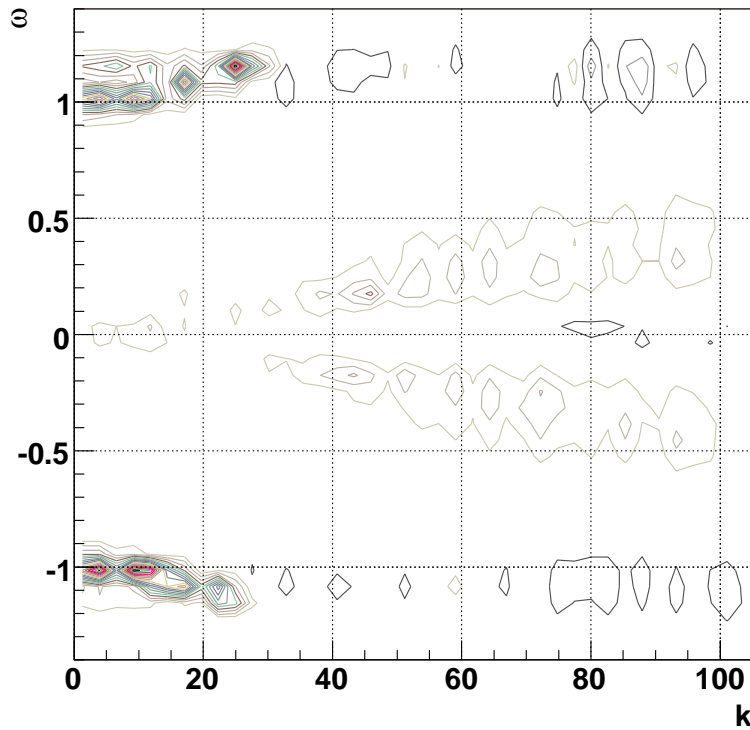


15. 10 7000 7000 1.1 5.0 1.0 0.1 c 4.0 c 200 500 12 1000 XXX 6.0 0.0

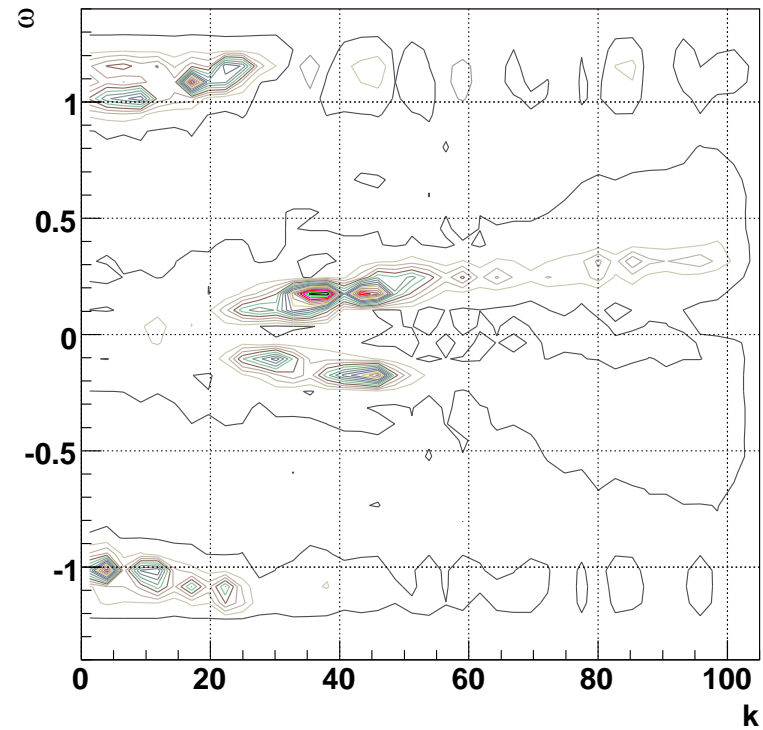


Dispersion Relation

15. 10 7000 7000 1.1 5.0 0.0 0.1 c 4.0 c 200 500 12 1000 XXX 1.2 0.0

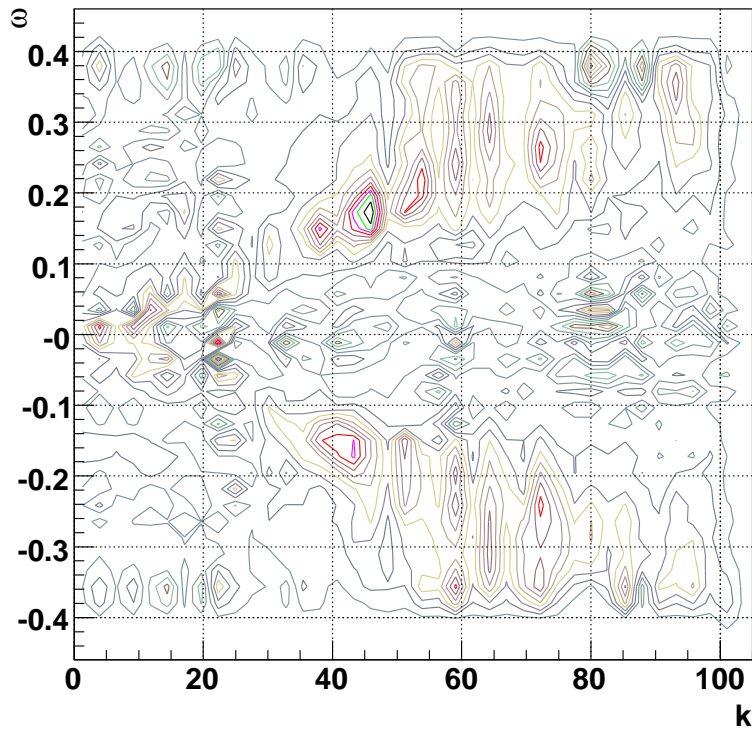


15. 10 7000 7000 1.1 5.0 1.0 0.1 c 4.0 c 200 500 12 1000 XXX 1.2 0.0

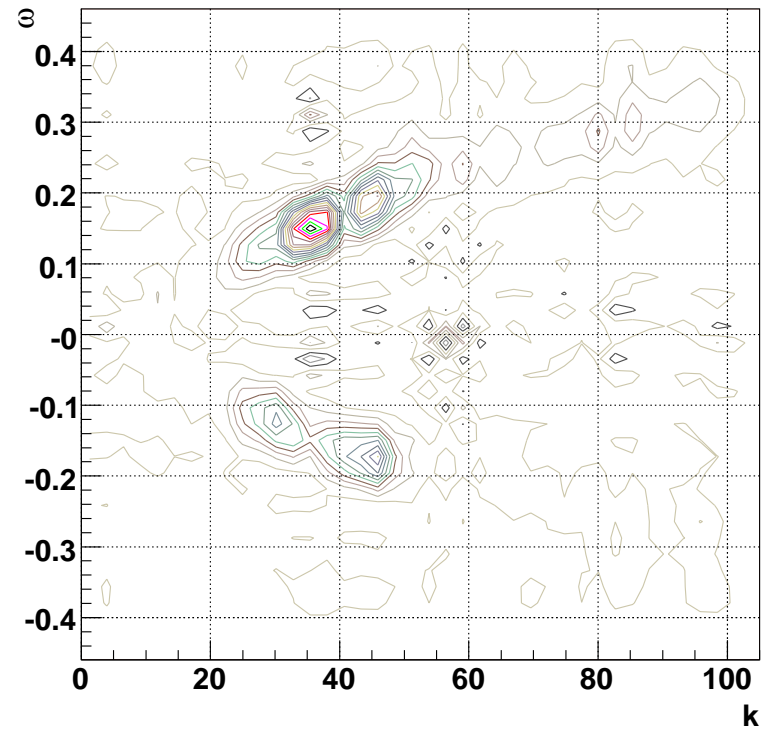


Dispersion Relation

15. 10 7000 7000 1.1 5.0 0.0 0.1 c 4.0 c 200 500 12 1000 XXX 0.4 0.0



15. 10 7000 7000 1.1 5.0 1.0 0.1 c 4.0 c 200 500 12 1000 XXX 0.4 0.0



Summary

- Plasma wakefields induced by Alfvén shocks can in principle efficiently accelerate UHECR particles.
- Preliminary simulation results support the existence of this mechanism, but more investigation needed.
- In addition to GRB, there exist abundant astrophysical sources that carry relativistic plasma outflows/jets.
- Other electromagnetic sources, for example GRB prompt signals, filamentation of e^+e^- jets, intense neutrino outburst, etc., can also excite plasma wakefields.