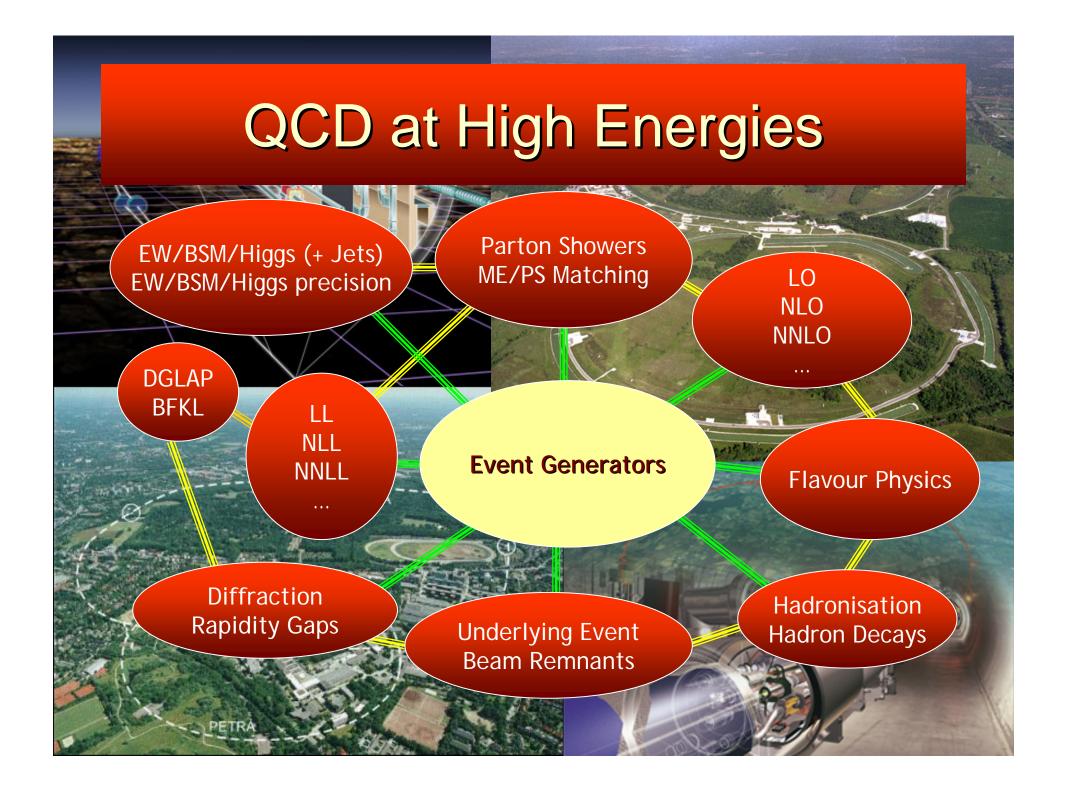
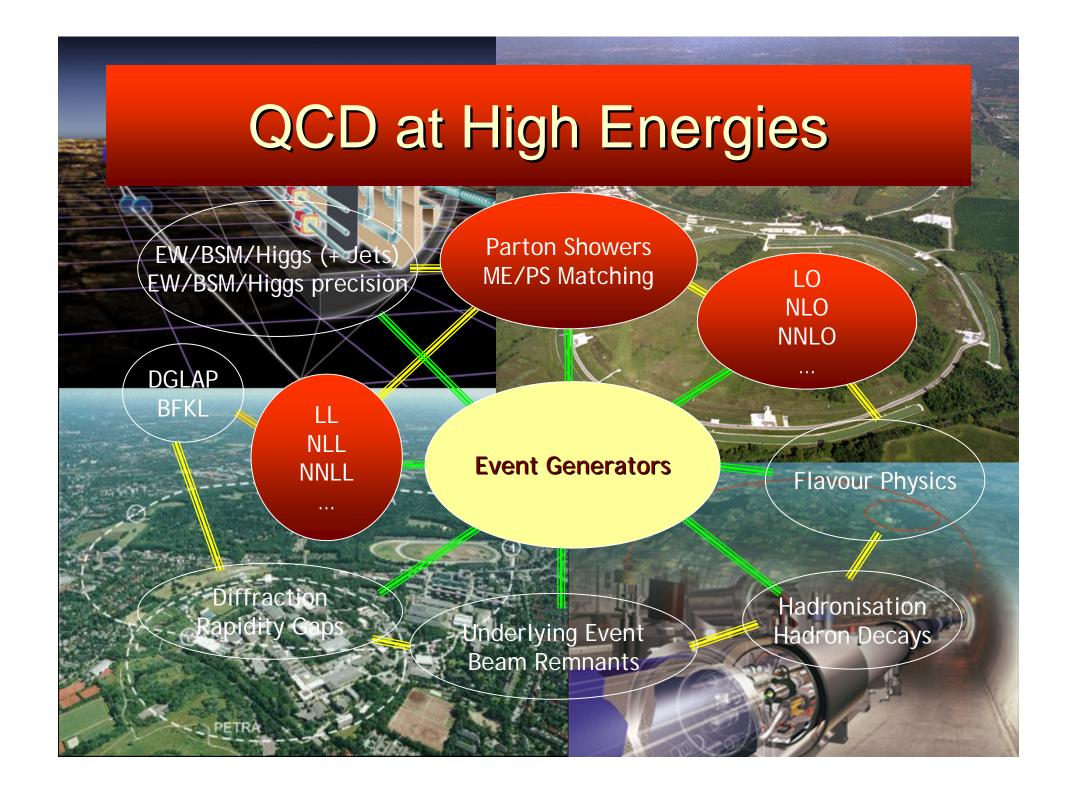
Progress on VirCol

A Parton Shower MC based on Antenna Formalism

W. Giele (Fermilab), D. Kosower (Saclay), P. Skands (Fermilab)







Hard & Soft

Matrix Elements (Fixed Order):

- •Fixed order in α -> Exact interference, helicity, loops, ...
- At present can do 2->5/6 (less with loops)
- Perturbative expansion better at higher energy (asymptotic freedom)
- •Multiple soft emissions important for full event structure = exclusive observables
- •Widely separated scales -> big logs / big truncation errors.
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Parton Showers:

- Derived in universal limit of QCD → depend on universal parameters
- •Exponentiate \rightarrow infinite $O(\alpha) \rightarrow$ ideal for widely separated scales (logs resummed)
- Arbitrary number of partons in final state → match to hadronisation descriptions
- •Derived in <u>limit</u> (collinear) of QCD→ approx. for wide-angle / hard emissions.

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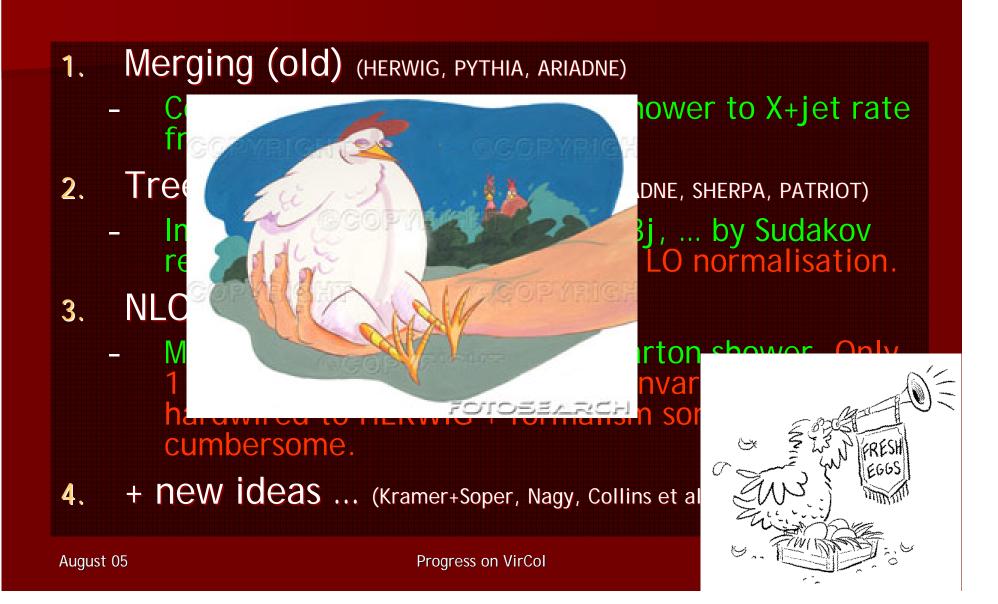
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Marriage desireable!!

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Possible Ceremonies



Vircol – Basic SKETCH

- Perturbative expansion for some observable J, $d\sigma = \Sigma_{m=0} d\sigma_m \quad ; \qquad d\sigma_m = d\Pi_m |M|^2 \delta(J J(k_1, k_2, ..., k_m))$
- Assume we know some Matrix Elements $d\sigma_0$, $d\sigma_1$, ... $d\sigma_n$ (w or w/o loops)
- And we have some approximation $T_{n \to n+1}$, so that $d\sigma_{n+1} \sim T_{n \to n+1} d\sigma_n$ (~ parton shower)
- A 'best guess' cross section is then:

$$d\sigma \sim d\sigma_0 + d\sigma_1 + \dots + d\sigma_n \left(1 + T_{n \to n+1} + T_{n \to n+1} T_{n+1 \to n+2} + \dots \right)$$

$$d\sigma \sim d\sigma_0 + d\sigma_1 + \dots + d\sigma_n S_n \qquad ; S_n = 1 + T_{n \to n+1} S_{n+1}$$

- For this to make sense, the $T_{n \to n+1}$ have to at least contain the correct singularities (in order to correctly sum up all logarithmically enhanced terms), but they are otherwise arbitrary.
- We will now reorder this series in a useful way ...

Reordering example: h→gg

■ Assume we know ME for H→gg and H→ggg. Then reorder:

- I.e shower off gg and modified ggg matrix element.
- Double counting avoided since singularities/shower subtracted in $\mathrm{d}\chi_{\mathrm{ggg}}$.

What IS THE Difference?

CKKW (& friends) in a nutshell:

- 1. Generate a n-jet Final State from n-jet (singular) ME.
- 2. Construct a "fake" PS history.
- Apply Sudakov weights on each "line" in history → from inclusive n-jet ME to exclusive n-jet (i.e. probability that n-jet FS remains n-jet above cutoff) → gets rid of double counting when mixed with other ME's (Sudakov wt dampens singularity).
- 4. Apply PS with no emissions above cutoff.

VirCol in a nutshell:

- 1. Subtract PS singularities from n-jet ME (antenna subtraction)
- 2. Generate a n-jet Final State from the subtracted (finite) ME.
- 3. Apply PS → Leading Logs resummed.
- + full NLO: divergent part already there = unitarity of shower assumption \rightarrow just include extra finite contribution in $d\sigma_0$: $d\sigma = d\sigma_0^{(0)} + d\sigma_1^{(0)} + sing[d\sigma_0^{(1)}] + F^{(1)} + \dots$
- + now NNLO/NLL possible → talks by Gehrmann, Gehrmann-De Ridder

The ANTENNA Shower

■ So far, we have written a C++ code that (for the moment) generates a pure gluon cascade ordered in:

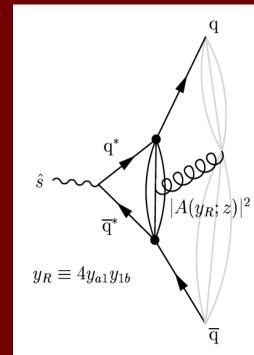
$$y_R = 4s_{a1}s_{1b}/s_{a1b}^2 = 4p_{T:ARIADNE}^2/s_{a1b}$$

...with the antenna / subtraction function:

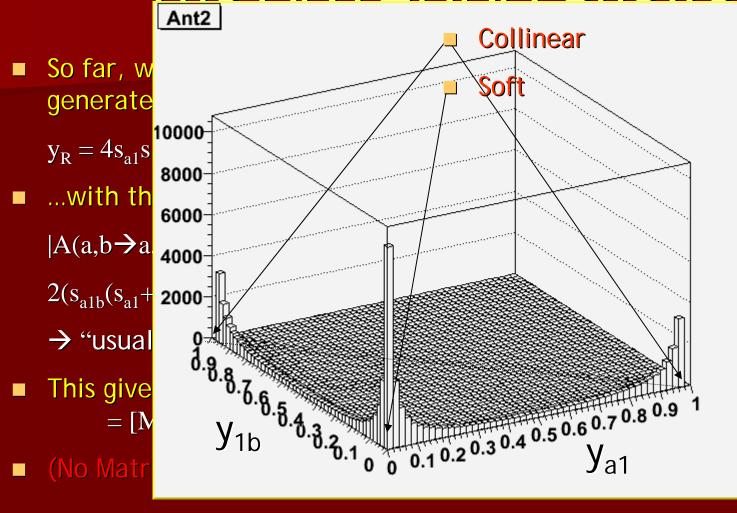
$$|A(a,b\to a,1,b)|^2 =$$

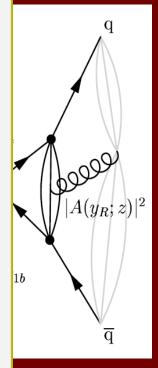
$$2(s_{a1b}(s_{a1}+s_{1b})+s_{ab}^{2})^{2}/(s_{a1}s_{1b}(s_{a1b}s_{ab}+s_{a1}s_{1b})s_{a1b})$$

- → "usual" collinear limit, but different outside.
- This gives an analytical Sudakov integral = [Mathematica output].
- (No Matrix Elements yet ... but work in progress)



The ANTENNA Shower





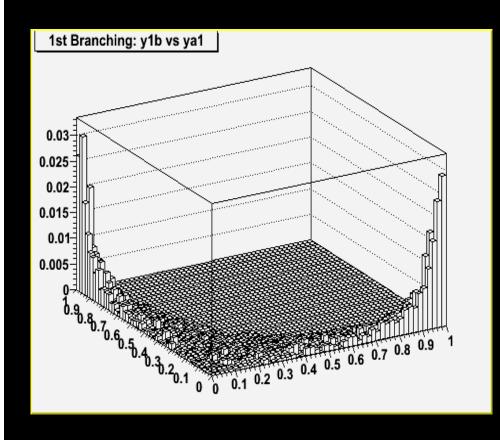
- Sudakov \rightarrow y_R for next branch \rightarrow select phase space point along iso-y_R contour:
 - 1. Rewrite Antenna function in terms of $y^2=y_R$;

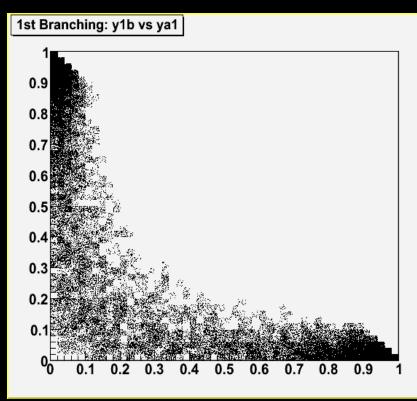
$$\xi = (s_{a1} + s_{1b})/R_1;$$

$$0 dy \frac{1}{y} R_{1=y} dw \frac{(wy + (1i wy)^2)^2}{1 + y^2 = 4i wy} P \frac{1}{w^2i}$$

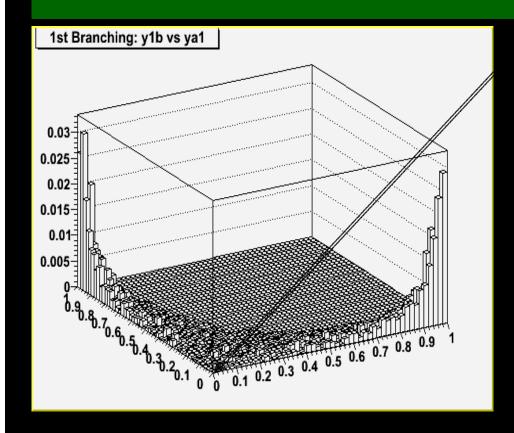
2. Partial-fraction singular structure + overestimate

numerators
$$\rightarrow$$
 generate uniform R and solve for ξ_R :
$$R = \begin{cases} R \\ R \end{cases} \quad d \Rightarrow \quad \frac{A(y; x)}{1; \quad xy + y^2 = 4} + \frac{B(y; x)}{x^2; \quad 1}$$

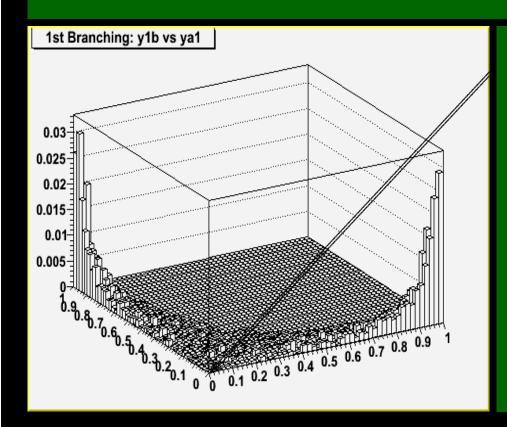




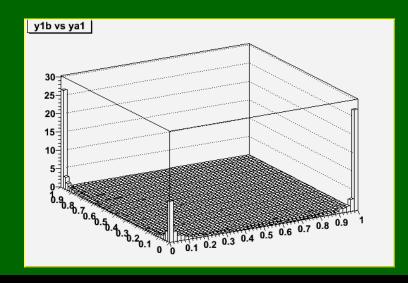
Q1: Where's the soft singularity?



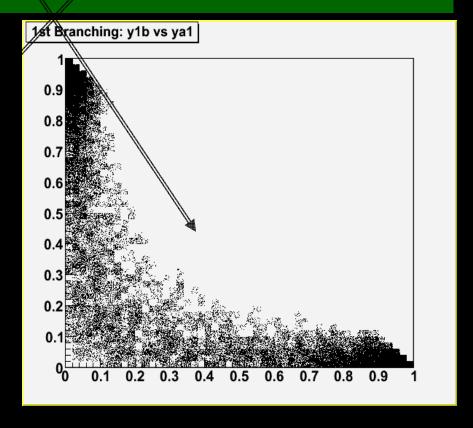
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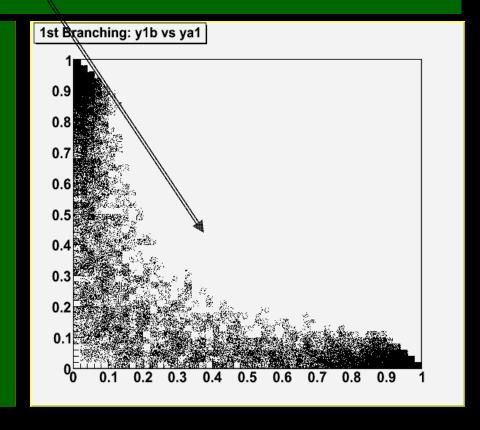
■ A1: Sudakov suppressed, (sum of ordered branching probs → unordered probability)



- Q1: Where's the soft singularity?
- Q2: Is that a 'dead region' ?

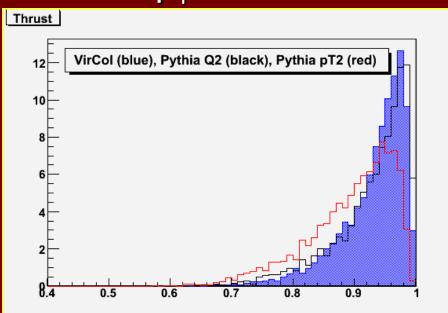


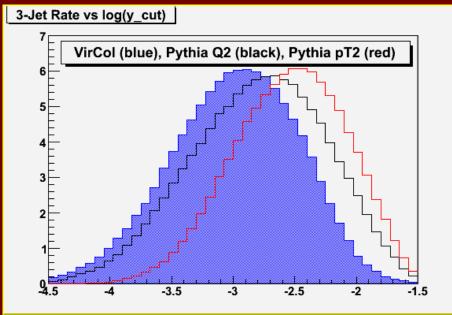
- Q1: Where's the soft singularity?
- Q2: Is that a 'dead region' ?
- A2: Yes, it is cut out by 'unresolution criterion', i.e. that neighbour dipoles remain resolved after branching.
- Due to H→gg in colour singlet state! (pathological)
- Eventually, could be filled by Matrix Elements and/or by changing evolution var.



Preliminary Results

Thrust and 3-jet rate, compared to Q²-ordered and p_T²-ordered PYTHIA showers.





Preliminary: matrix elements should be added and parton-level matched to hadronisation models eventually + all showers only include gluons here...

Conclusion & Outlook

- Construction of VIRCOL shower monte carlo:
 - gluons shower MC (based on LO, done!)
 - gluons shower MC (based on NLO)
 - parton shower MC (LO/NLO(/NNLO))
 - parton shower MC (NLL + NLO/NNLO)
 - Hadron collider shower MC's
 - Higher order Sudakov factor calculations(this will reduce a lot of implicit and explicit uncertainties: e.g. renormalization scale, choice of subtraction function,...)