

Graph 5

Graph 6

Graph 7

Graph 8

NLO-QCD calculation in GRACE

- GRACE status -

Graph 9

Y. Kurihara (KEK)

GRACE Group

19/Aug./2005 @ LoopFest IV

Graph 13

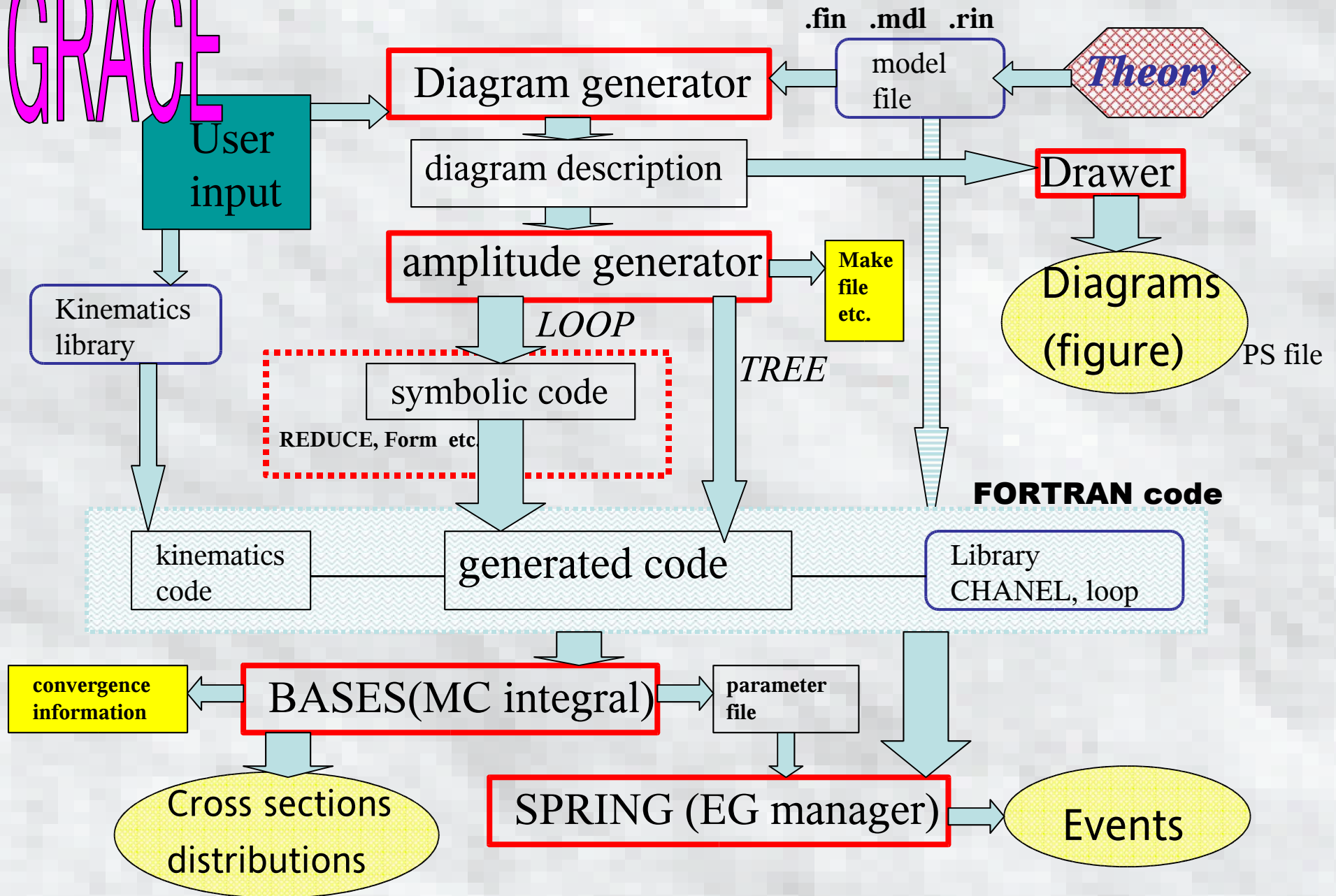
Graph 16

GRACE Author list

A map of Japan is shown in the background. The main islands are colored green, and the surrounding water is light blue. A red dot is placed on the main island of Honshu, specifically in the Kanto region, indicating the location of the authors.

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M. Jimbo, T. Kaneko,
K. Kato, S. Kawabata,
T. Kon, Y. Kurihara,
M. Kuroda, N. Nakazawa,
Y. Shimizu, H. Tanaka,
Y. Yuasa**

GRACE





Physics in LHC

LHC Experimental requirement

New Particle Search/Precision Measurements

LO-QCD Event generator+K-factor

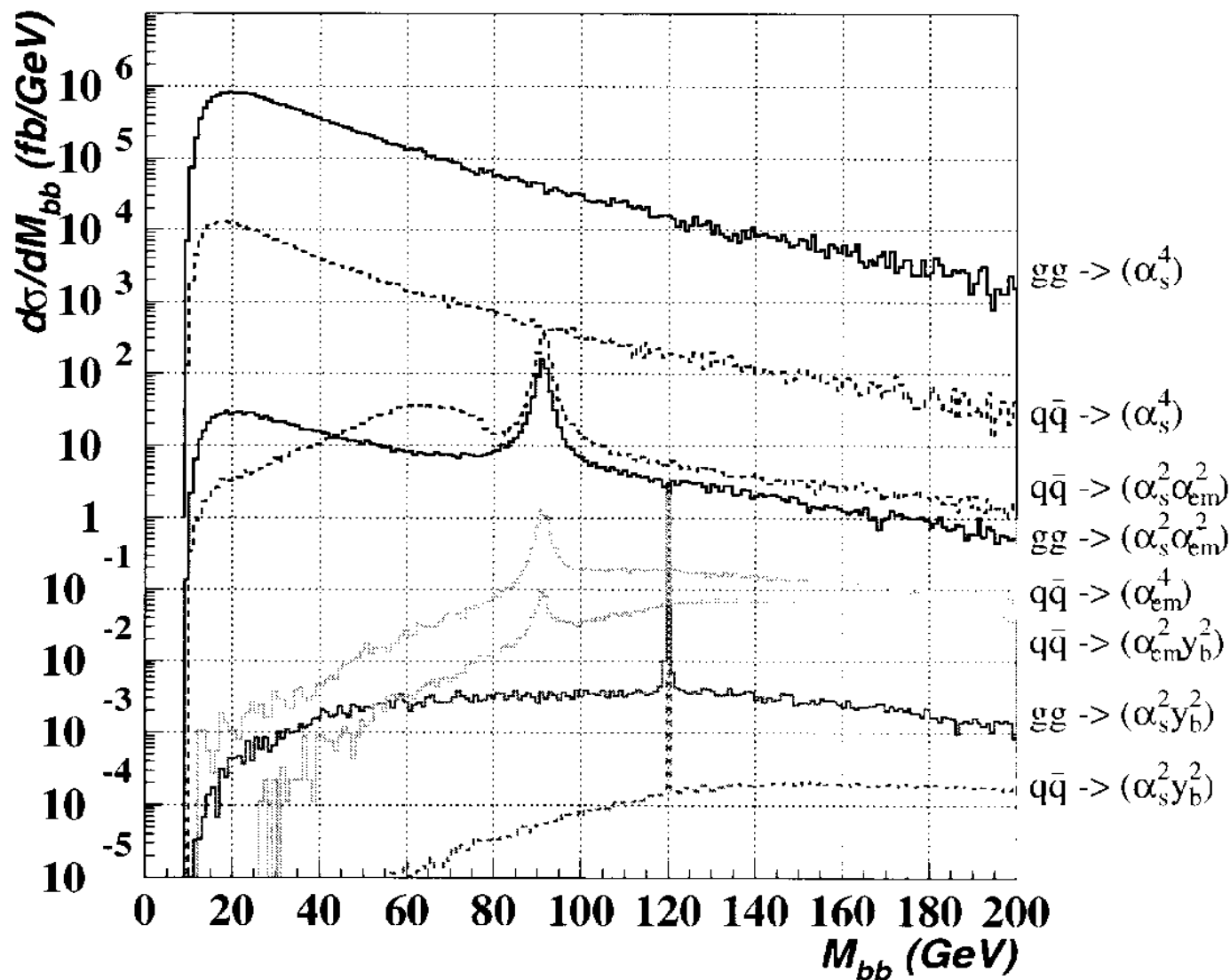


Obviously not enough!

We need
NLO Event generator!

QCD-event generator @ Tree-level

$pp \rightarrow bbbb$, TEVATRON/LHC, GR@PPA_4b, S. Tsuno



S. Tsuno et al., Comput.Phys.Commun.151(2003)216

pp → many, TEVATRON/LHC, GR@PPA_ALL,
S. Tsuno

- W + jets (up to 4 jets) with the subsequent W decay to a fermion pair,
- Z + jets (up to 4 jets) with the subsequent Z decay to a fermion pair,
- Four bottom quarks via Z and Higgs-boson mediated processes as well as those from pure QCD interactions (same as GR@PPA_4b),
- top-quark pair with the subsequent decay to W and b, and the W decay to a fermion pair,
- di-boson (WW, WZ and ZZ) with the subsequent W/Z decay to a fermion pair.



Loop Calc. by GRACE-loop in ELWK proc. (Higgs production)

★ *Single Higgs production*

→ $e^+e^- \rightarrow ZH$ (full number of graphs = 341)

→ $e^+e^- \rightarrow \nu \nu H$ (1,350) *Phys.Lett. B559 (2003) 252-262*
A.Denner et.al. PLB 560(2003)196, NPB 660(2003)289

→ $e^+e^- \rightarrow e^+e^- H$ (4,470) *Phys.Lett.B600 (2004) 65-76*

★ *top Yukawa*

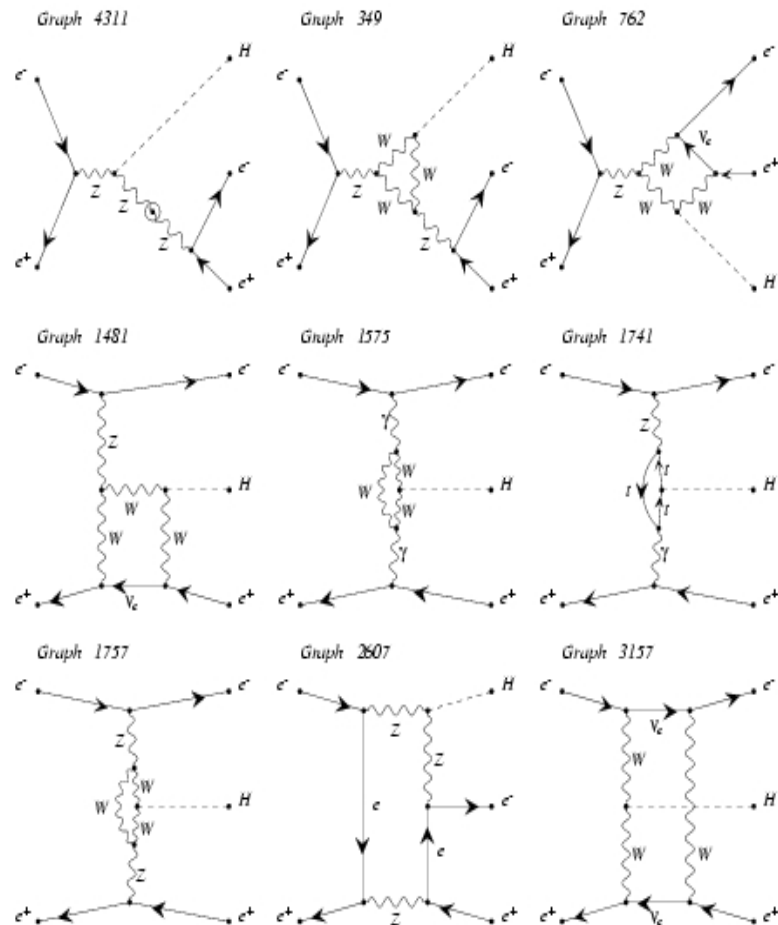
→ $e^+e^- \rightarrow t t H$ (2,327) *Phys.Lett. B571 (2003) 163-172*
Y.You et.al.PLB 571(2003)85
A.Denner et.al. PLB 575(2003)290, NPB 680 (2004)85

★ *Multi Higgs production*

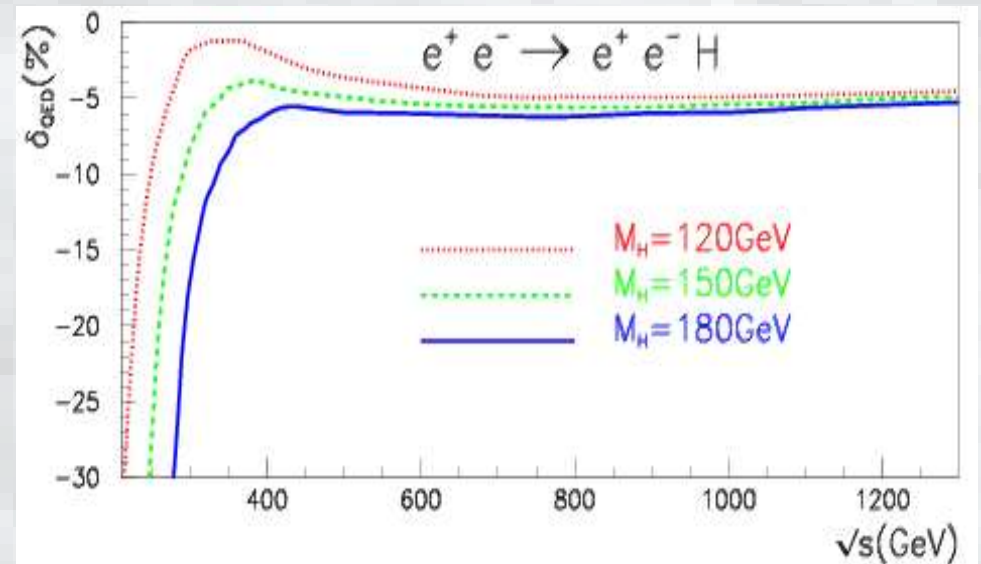
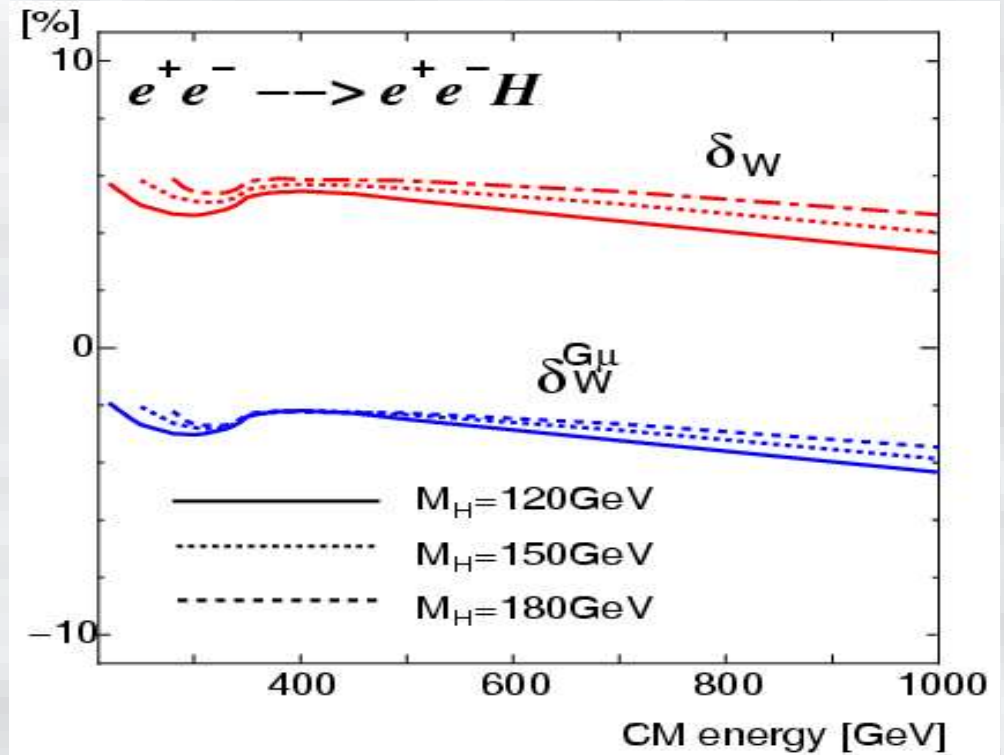
→ $e^+e^- \rightarrow ZHH$ (5,417) *Phys.Lett. B576 (2003) 152-164*
R.Zhang et.al.PLB(2004)349

→ $e^+e^- \rightarrow \nu_e \nu_e HH$ (19,638) ⇒ *Preliminary*

$$e^+e^- \rightarrow e^+e^-H$$



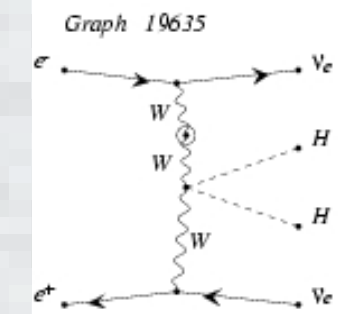
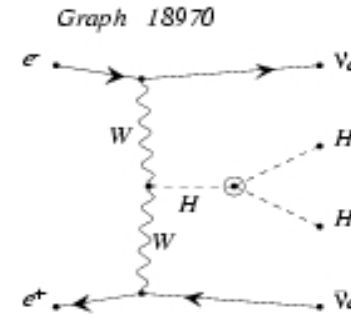
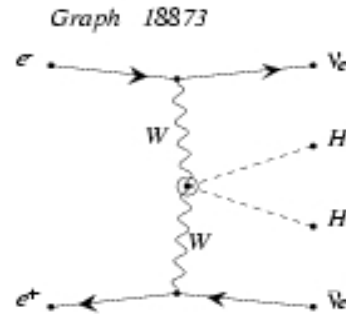
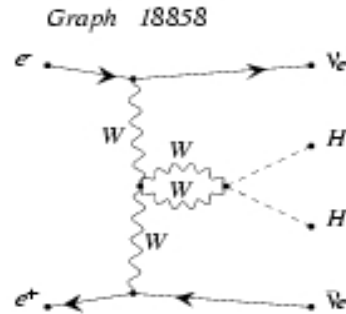
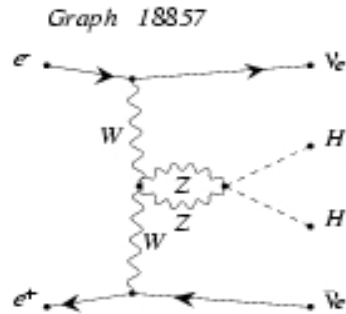
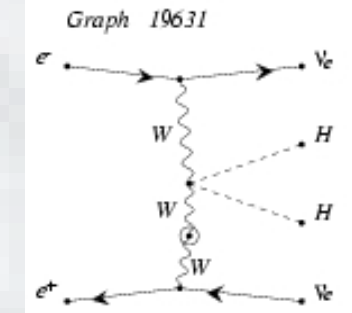
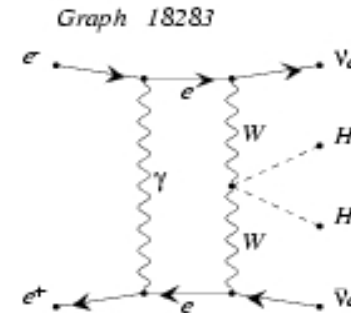
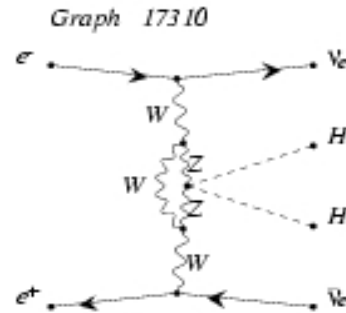
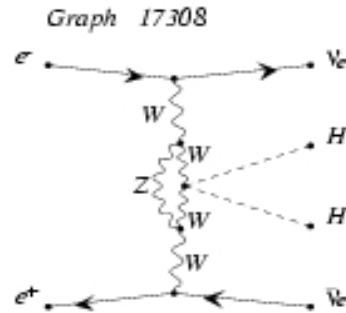
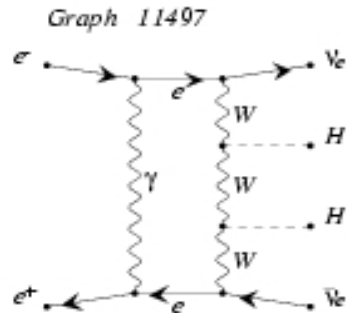
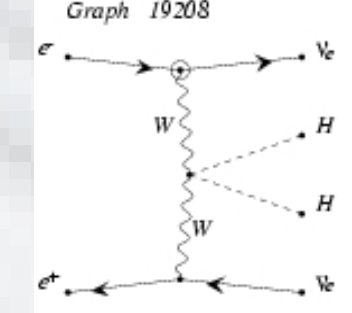
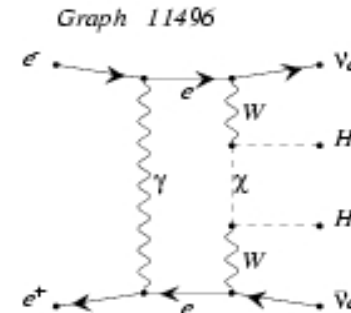
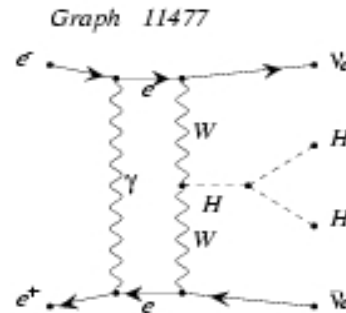
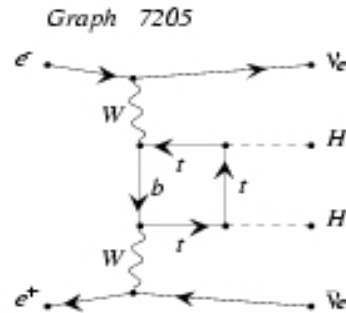
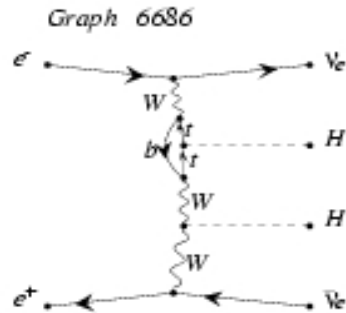
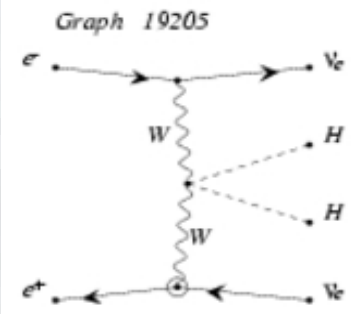
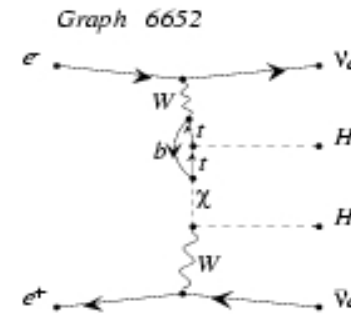
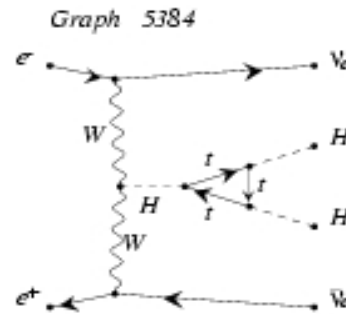
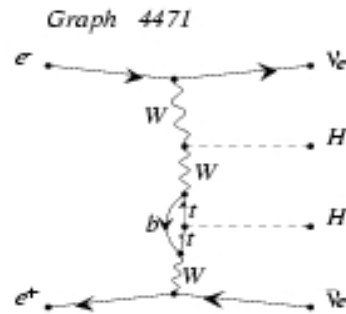
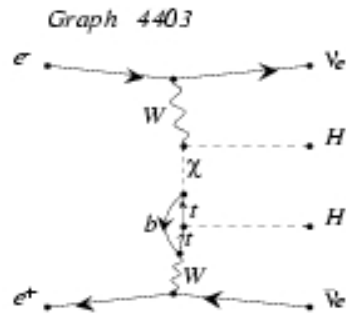
produced by GRACEFIG



GRACE *Phys.Lett. B559 (2003) 252-262*

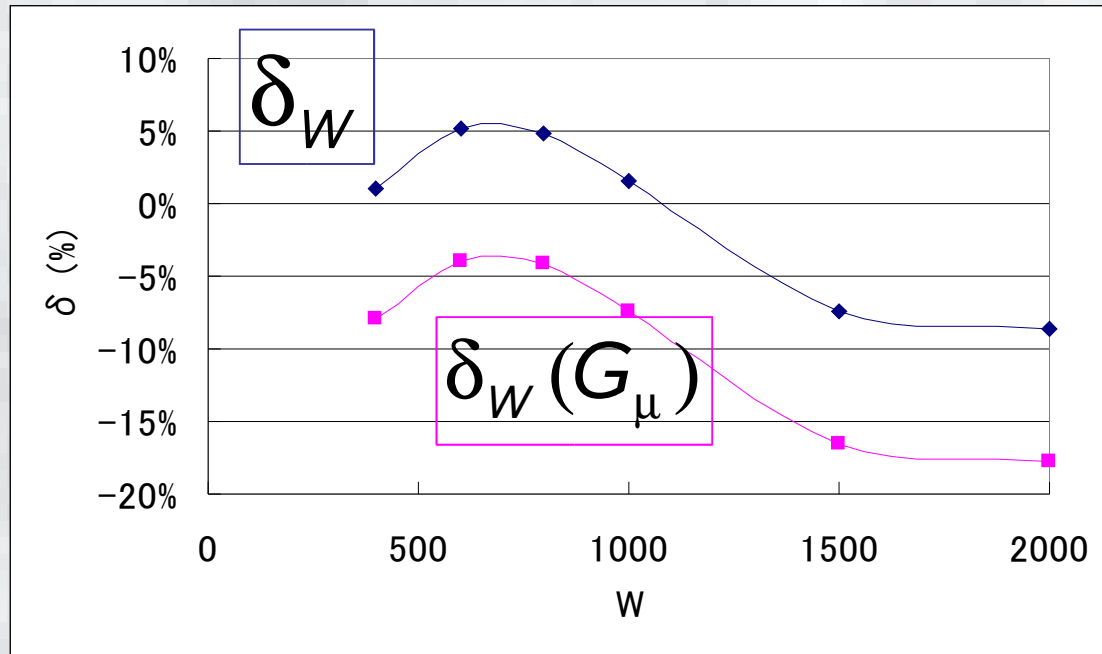


$e^+e^- \rightarrow \nu_e \bar{\nu}_e HH$ (final 4-body process)



$$e^+e^- \rightarrow \bar{\nu}_\mu \nu_\mu HH$$

$$\delta = \sigma(O(\alpha))/\sigma(\text{tree}) - 1$$



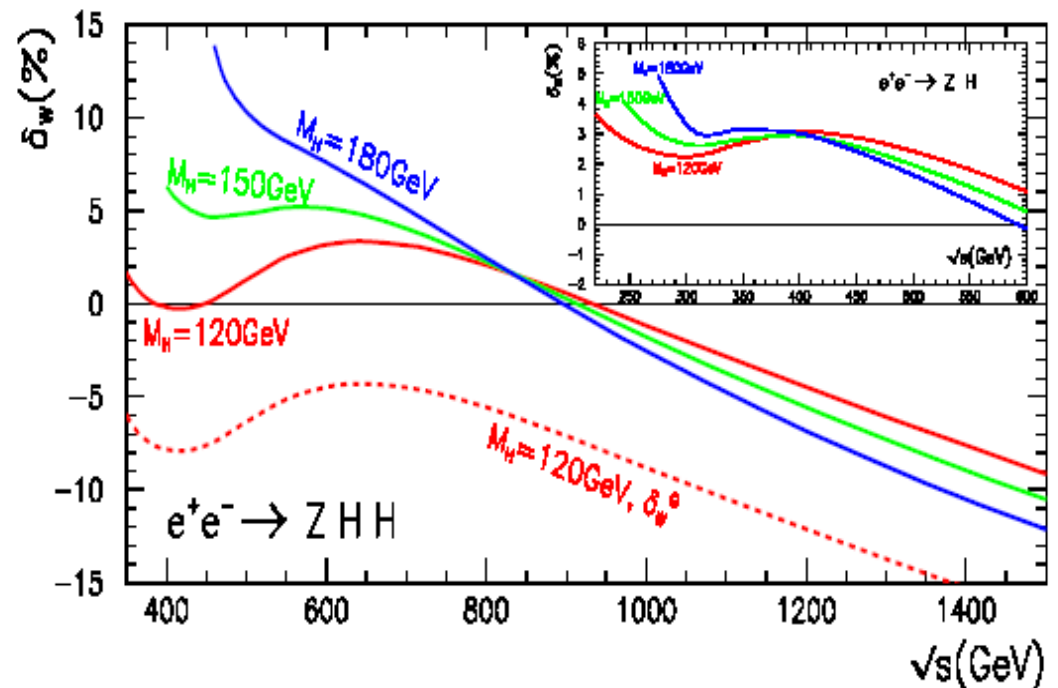
$M_H = 120 \text{ GeV}$

$m_t = 180 \text{ GeV}$

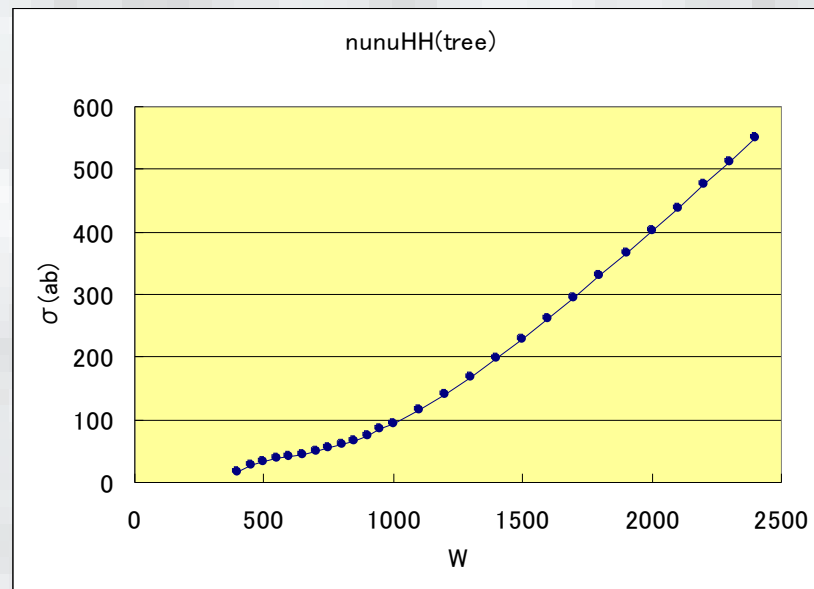
“s-channel”

$$e^+e^- \rightarrow ZHH$$

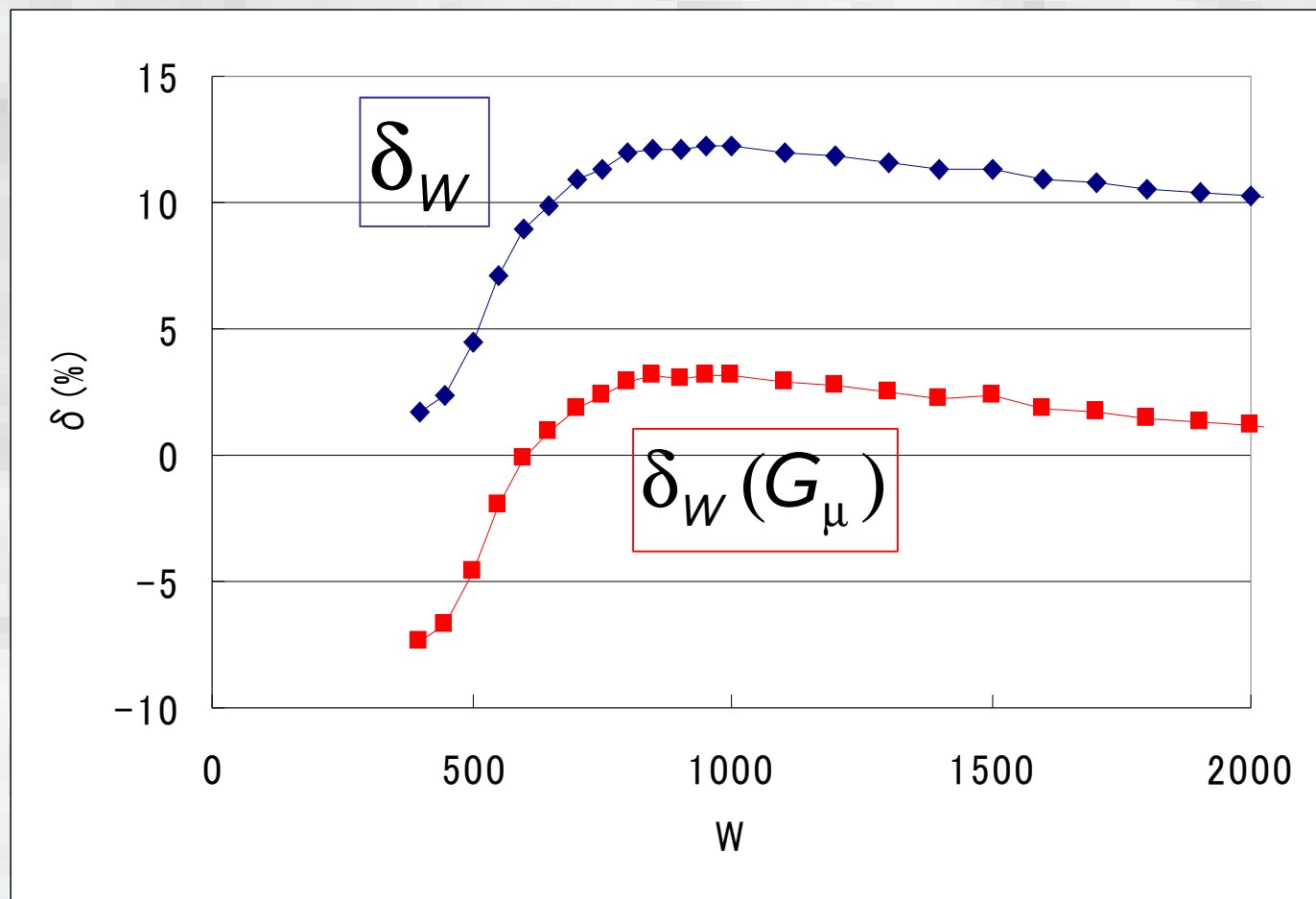
Phys.Lett.B 576(2003)152.
 $m_t = 174 \text{ GeV}$



$$e^+e^- \rightarrow \nu\bar{\nu}HH$$



MH=120GeV



Internal consistency check

For EW part:

System passes successfully the usual checks:

- ultraviolet finiteness (better than 20 digits)
- infrared finiteness (better than 20 digits)
- NLG dependence (better than 20 digits)
- k_c dependence consistent with MonteCarlo statistical error (0.02%)

1. One random phase space point
2. Full set of diagrams
3. Quadruple precision

• New feature of NLO-QCD issues in GRACE

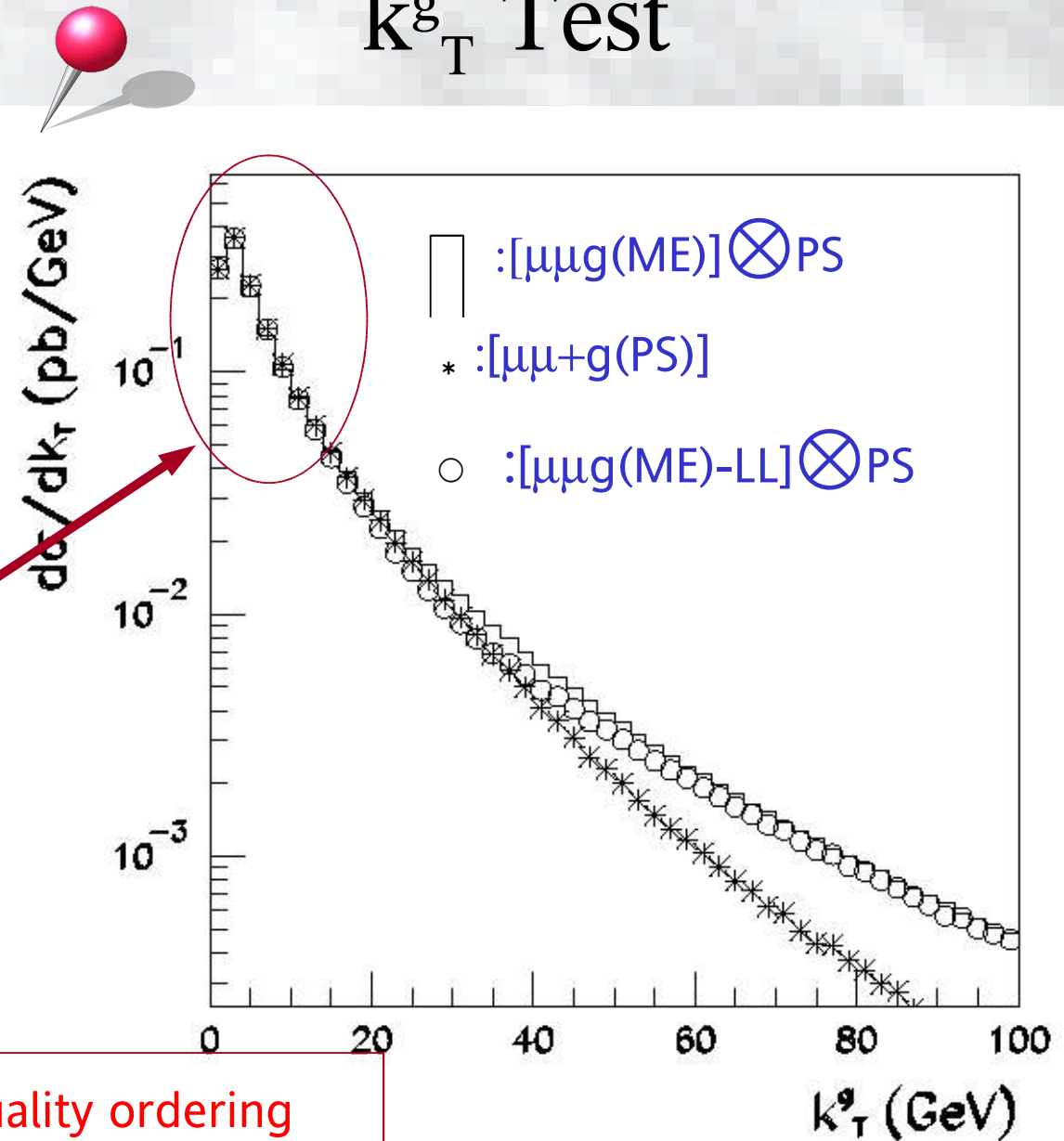
QCD-tree : OK
ELWK 1-loop : OK

- PDF/PS $\leftrightarrow$ Real emission Double counting
 - Virtuality ordering/LL-subtraction
(YK et al. Nucl. Phys. B654 (2003) 301)
- Dimensional regularization in loop integrals for IR
(fictitious mass in photon in ELWK)
- IR (soft/collinear) approximation terms

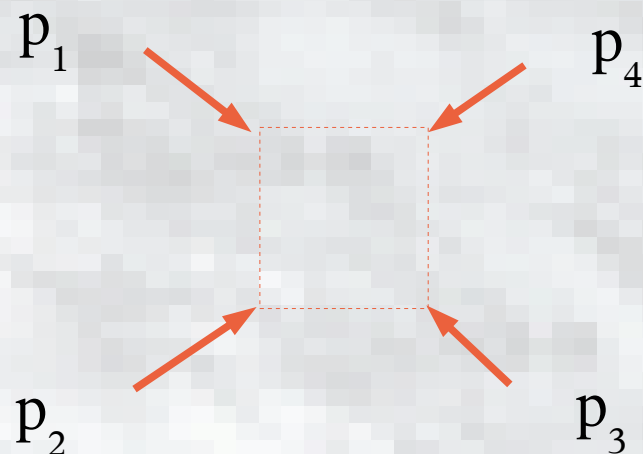
Drell-Yan process

- Process :
 $u\bar{u} \rightarrow \mu^+\mu^- (+\text{gluon})$
in $p\bar{p}$ collision
- Cuts:
 $\sqrt{s}_{\mu\mu} > 40 \text{ GeV}$
 $k_T^g > 1 \text{ GeV}$

k_T^g Test



Box Integral



$$J_{(4)}(s, t; p_1^2, p_2^2, p_3^2, p_4^2; n_x, n_y, n_z) = \frac{\Gamma(2 - \varepsilon_{IR})}{(4\pi)^2 (4\pi\mu_R^2)^{\varepsilon_{IR}}} \int_0^1 dx \int_0^{1-x} dy \int_0^{1-x-y} dz \frac{x^{n_x} y^{n_y} z^{n_z}}{D^{2-\varepsilon_{IR}}},$$

$$D = -s xz - t yw - p_1^2 xy - p_2^2 yz - p_3^2 zw - p_4^2 xw - i0,$$

$$w = 1 - x - y - z,$$

$$s = (p_1 + p_2)^2,$$

$$t = (p_1 + p_4)^2.$$

All on-shell (massless) external legs

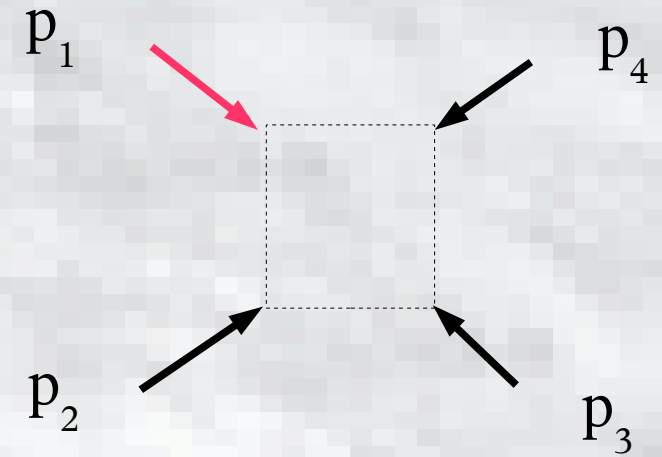
$$\begin{aligned}
 J_4(s, t; 0, 0, 0, 0; n_x, n_y, n_z) &= \frac{1}{(4\pi)^2 s t} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) n_x! \Gamma(\varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR}) \\
 &\times \left[\left(\frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \left(\frac{-t}{s} \right)^{n_x} \frac{B(1 + n_z, n_x + n_y + \varepsilon_{IR})}{\Gamma(n_x + \varepsilon_{IR})} \right. \\
 &\times {}_2F_1 \left(1 + n_x, n_x + n_y + \varepsilon_{IR}, 1 + n_x + n_y + n_z + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right) \\
 &+ \left(\frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \sum_{l=0}^{n_x} \left(\frac{-s}{t} \right)^l \frac{(-1)^l}{\Gamma(l + \varepsilon_{IR}) (n_x - l)!} B(1 + n_y, l + n_z + \varepsilon_{IR}) \\
 &\times \left. {}_2F_1 \left(1 + l, l + n_z + \varepsilon_{IR}, 1 + l + n_y + n_z + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}} \right) \right],
 \end{aligned}$$

Scalar Integral

$$\begin{aligned}
 J_{(4)}(s, t; 0, 0, 0, 0; 0, 0, 0) &= \frac{1}{(4\pi)^2 s t} \frac{B(\varepsilon_{IR}, \varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR})}{\varepsilon_{IR}} \\
 &\times \left[\left(\frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left(1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}} \right) + \left(\frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left(1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right) \right]
 \end{aligned}$$

This result is compared with G. Duplanić, B. Nižić, Eur. Phys. J. C **20**, 357 (2001)

One off-shell box integral



$$\begin{aligned}
 J_4(s, t; p_1^2, 0, 0, 0; n_x, n_y, n_z) = & \frac{\Gamma(2 - \varepsilon_{IR})}{(4\pi)^2 (4\pi\mu_R^2)^{\varepsilon_{IR}}} \int_0^1 dx \int_0^{1-x} dy \int_0^{1-x-y} dz \frac{x^{n_x} y^{n_y} z^{n_z}}{(-xzs - y(1-x-y-z)t - p_1^2 xy - i0)^{2-\varepsilon_{IR}}} \\
 = & \frac{1}{(4\pi)^2 s t} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) n_x! \Gamma(\varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR}) \\
 & \times \left[\left(\frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \left(\frac{-t}{s} \right)^{n_x} \frac{B(1 + n_z, n_x + n_y + \varepsilon_{IR})}{\Gamma(n_x + \varepsilon_{IR})} \mathcal{I}^{(1)} \right. \\
 & \left. + \left(\frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \sum_{l=0}^{n_x} \frac{(-1)^l B(1 + n_y, l + n_z + \varepsilon_{IR})}{\Gamma(l + \varepsilon_{IR}) (n_x - l)!} \mathcal{I}_l^{(2)} \right]
 \end{aligned}$$

$$\mathcal{I}^{(1)} = B(1+n_z, n_x+n_y+\varepsilon_{IR}) {}_2F_1\left(1+n_x, n_x+n_y+\varepsilon_{IR}, 1+n_x+n_y+n_z+\varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}}\right)$$

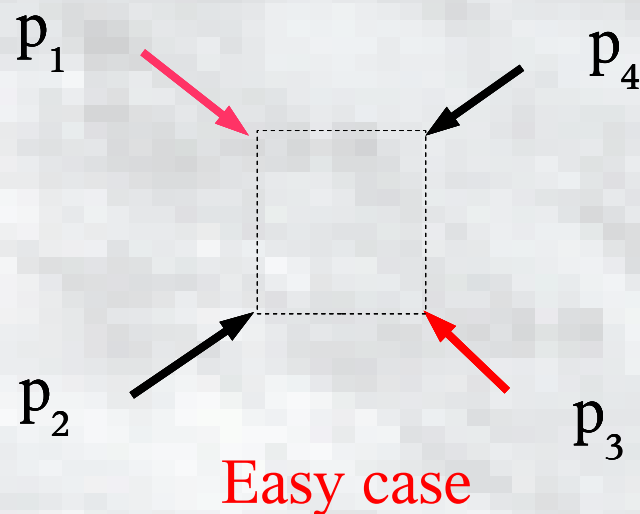
$$\begin{aligned} \mathcal{I}_l^{(2)} &= \sum_{k_1=0}^{n_z} n_z C_{k_1} \left(\frac{s}{p_1^2-s}\right)^{n_y+k_1} \sum_{k_2=0}^{n_y+k_1} n_{y+k_1} C_{k_2} (-1)^{n_y+k_2} \left(\frac{-t}{s}\right) \\ &\times \int_0^1 dw \left(1 + \frac{\tilde{u}}{\tilde{s}} w\right)^{-(l+1)} \left(1 + \frac{\tilde{t} + \tilde{u}}{\tilde{s}} w\right)^{k_2+l-1+\varepsilon_{IR}} \\ &= \sum_{k_1=0}^{n_z} \sum_{k_2=0}^{n_y+k_1} n_z C_{k_1} n_{y+k_1} C_{k_2} (-1)^{k_1+k_2} \left(\frac{s}{p_1^2-s}\right)^{n_y+k_1} \frac{1}{l+k_2+\varepsilon_{IR}} \left(1 + \frac{u}{t}\right)^l \\ &\times \left[{}_2F_1\left(1+l, l+k_2+\varepsilon_{IR}, 1+l+k_2+\varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}}\right) \right. \\ &\left. - \left(\frac{\tilde{p}_1^2}{\tilde{s}}\right)^{l+k_2+\varepsilon_{IR}} {}_2F_1\left(1+l, l+k_2+\varepsilon_{IR}, 1+l+k_2+\varepsilon_{IR}, -\frac{\tilde{u}\tilde{p}_1^2}{\tilde{t}\tilde{s}}\right) \right], \end{aligned}$$

Scalar Integral

$$J_4(s, t; p_1^2, 0, 0, 0; 0, 0, 0) = \frac{1}{(4\pi)^2 s t} \frac{B(\varepsilon_{IR}, \varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR})}{\varepsilon_{IR}} \\ \times \left[\left(\frac{-\tilde{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left(1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{t}} \right) + \left(\frac{-\tilde{t}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left(1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}}{\tilde{s}} \right) \right. \\ \left. - \left(\frac{-\tilde{p}_1^2}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} {}_2F_1 \left(1, \varepsilon_{IR}, 1 + \varepsilon_{IR}, -\frac{\tilde{u}\tilde{p}_1^2}{\tilde{t}\tilde{s}} \right) \right],$$

This result is compared with G. Duplanić, B. Nižić, Eur. Phys. J. C **20**, 357 (2001)

Two off-shell box integral



$$\begin{aligned}
 J_4(s, t; p_1^2, 0, p_3^2, 0; n_x, n_y, n_z) &= \frac{\Gamma(2 - \varepsilon_{IR})}{(4\pi)^2 (4\pi\mu_R^2)^{\varepsilon_{IR}}} \\
 &\int_0^1 dx \int_0^{1-x} dy \int_0^{1-x-y} dz \frac{x^{n_x} y^{n_y} z^{n_z}}{(-xzs - y(1-x-y-z)t - p_1^2 xy - p_3^2 z(1-x-y-z) - i0)^{2-\varepsilon_{IR}}} \\
 &= \frac{1}{(4\pi)^2 (s - p_3^2)(t - p_3^2)} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) n_x! \Gamma(\varepsilon_{IR}) \Gamma(1 - \varepsilon_{IR}) \\
 &\times \left[\left(-\frac{\bar{t} - p_3^2}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \left(-\frac{t - p_3^2}{s - p_3^2} \right)^{n_x} \frac{1}{\Gamma(n_x + \varepsilon_{IR})} \mathcal{I}^{(1)} + \left(\frac{-\bar{s}}{4\pi\mu_R^2} \right)^{\varepsilon_{IR}} \sum_{l=0}^{n_x} \frac{(-1)^l}{\Gamma(l + \varepsilon_{IR}) (n_x - l)!} \mathcal{I}_l^{(2)} \right]
 \end{aligned}$$

$$\begin{aligned}
\mathcal{I}^{(1)} &= \frac{1}{n_x + \varepsilon_{IR}} \sum_{k_1=0}^{n_x} \sum_{k_2=0}^{n_y+k_1} n_x C_{k_1} n_{y+k_1} C_{k_2} (-1)^{k_1+k_2} \left(\frac{p_3^2 - s}{u} \right)^{n_y+k_1} (1-\alpha)^{k_2-n_x-1} \\
&\times \left[\left(1 + \frac{\tilde{p}_3^2}{\tilde{t} - \tilde{p}_3^2} \right)^{n_x+\varepsilon_{IR}} {}_2F_1 \left(1+n_x-k_2, n_x+\varepsilon_{IR}, 1+n_x+\varepsilon_{IR}, \frac{\tilde{u}/(\tilde{s} - \tilde{p}_3^2) + \alpha}{\alpha - 1} \right) \right. \\
&\left. - \left(\frac{\tilde{p}_3^2}{\tilde{t} - \tilde{p}_3^2} \right)^{n_x+\varepsilon_{IR}} {}_2F_1 \left(1+n_x-k_2, n_x+\varepsilon_{IR}, 1+n_x+\varepsilon_{IR}, \frac{\alpha}{\alpha - 1} \right) \right]
\end{aligned}$$

$$\begin{aligned}
\mathcal{I}_l^{(2)} &= \sum_{k_1=0}^{n_x} \sum_{k_2=0}^{n_y+k_1} n_x C_{k_1} n_{y+k_1} C_{k_2} (-1)^{k_1+k_2} \left(\frac{s}{s - p_1^2} \right)^{n_y+k_1} \frac{1}{l+k_2+\varepsilon_{IR}} \left(\frac{1}{1-\beta} \right)^{l+1} \left(\frac{t - p_3^2}{t+u - p_3^2} \right) \left(\frac{s}{s - p_3^2} \right)^l \\
&\times \left[{}_2F_1 \left(1+l, l+k_2+\varepsilon_{IR}, 1+l+k_2+\varepsilon_{IR}, \frac{\beta}{\beta - 1} \right) \right. \\
&\left. - \left(\frac{\tilde{p}_1^2}{\tilde{s}} \right)^{l+k_2+\varepsilon_{IR}} {}_2F_1 \left(1+l, l+k_2+\varepsilon_{IR}, 1+l+k_2+\varepsilon_{IR}, \frac{\beta}{\beta - 1} \frac{\tilde{p}_1^2}{\tilde{s}} \right) \right]
\end{aligned}$$

$$\alpha = \frac{\tilde{p}_3^2}{\tilde{t} - \tilde{p}_3^2} \frac{\tilde{u}}{\tilde{s} - \tilde{p}_3^2}, \quad \beta = \frac{\tilde{u}}{\tilde{s} - \tilde{p}_3^2} \frac{\tilde{s}}{\tilde{t} + \tilde{u} - \tilde{p}_3^2}$$

Numerical calculation (IR finite case)

For $l, m, n \in \mathcal{N}$,

$$\begin{aligned} & {}_2F_1(l, m+1, n+m+2; z) \\ &= \frac{1}{B(m+1, n+1)} \int_0^1 \tau^m (1-\tau)^n (1-z\tau)^{-l} d\tau \\ &= \sum_{k_1=0}^n \sum_{k_2=0}^{m+k_1} (-1)^{k_1+k_2} \frac{n C_{k_1} \quad m+k_1 C_{k_2}}{B(m+1, n+1)} \frac{1}{z^{m+k_1}} \int_0^1 (1-z\tau)^{-l+k_2} d\tau. \\ & \int_0^1 (1-z\tau)^{-l+k_2} d\tau = \begin{cases} -\frac{\ln(1-z)}{z} & k_2 - l + 1 = 0, \\ \frac{1}{k_2 - l + 1} \frac{(1-z)^{k_2 - l + 1} - 1}{-z} & k_2 - l + 1 \neq 0. \end{cases} \end{aligned}$$

Numerical check: mathematica $\leftrightarrow$ our FORTRAN program
more than ten digit agreement



Numerical calculation (IR divergent case)

$$\begin{aligned}
 \mathcal{I}_{l,m,n} &\equiv \int_0^1 \tau^{l+n-1+\varepsilon_{IR}} (1-\tau)^m (1-z\tau)^{-(l+1)} d\tau, \\
 &= B(1+m, l+n+\varepsilon_{IR}) \, \underline{{}_2F_1(1+l, l+n+\varepsilon_{IR}, 1+l+n+m+\varepsilon_{IR}, z)} \\
 &= \sum_{j=-1}^{\infty} \mathcal{F}_{l,m,n}^{(j)}(z) \varepsilon_{IR}^j,
 \end{aligned}$$

Expansion w.r.t. ε_{IR}

$$\tilde{F}_{j_1, j_2}^{(n)}(z) \equiv \frac{(-1)^n}{n!} \int_0^1 d\tau \, \tau^{j_1-1} (1-z\tau)^{-(j_2+1)} \ln^n \tau,$$

When $n=1$

$$\tilde{F}_{1,0}^{(1)}(z) = \frac{\text{Li}_2(z)}{z},$$

$$\tilde{F}_{1,1}^{(1)}(z) = -\frac{\ln(1-z)}{z}$$

$$\tilde{F}_{1,j_2+1}^{(1)}(z) = \frac{j_2}{j_2+1} \tilde{F}_{1,j_2}^{(1)} + \frac{(1-z)^{-j_2} - 1}{j_2(j_2+1)z}$$

$$\tilde{F}_{j_1,j_2}^{(1)}(z) = \frac{1}{z^{j_1-1}} \sum_{k=0}^{j_1-1} (-1)^k {}_{j_1-1}C_k \tilde{F}_{1,j_2-k}^{(1)}(z)$$

When $n=2$

$$\tilde{F}_{1,0}^{(2)}(z) = \frac{\text{Li}_3(z)}{z}$$

$$\tilde{F}_{1,1}^{(2)}(z) = \frac{\text{Li}_2(z)}{z}$$

$$(j_2 + 1)\tilde{F}_{1,j_2+1}^{(2)}(z) - j_2\tilde{F}_{1,j_2}^{(2)}(z) - \tilde{F}_{1,j_2}^{(1)}(z) = 0$$

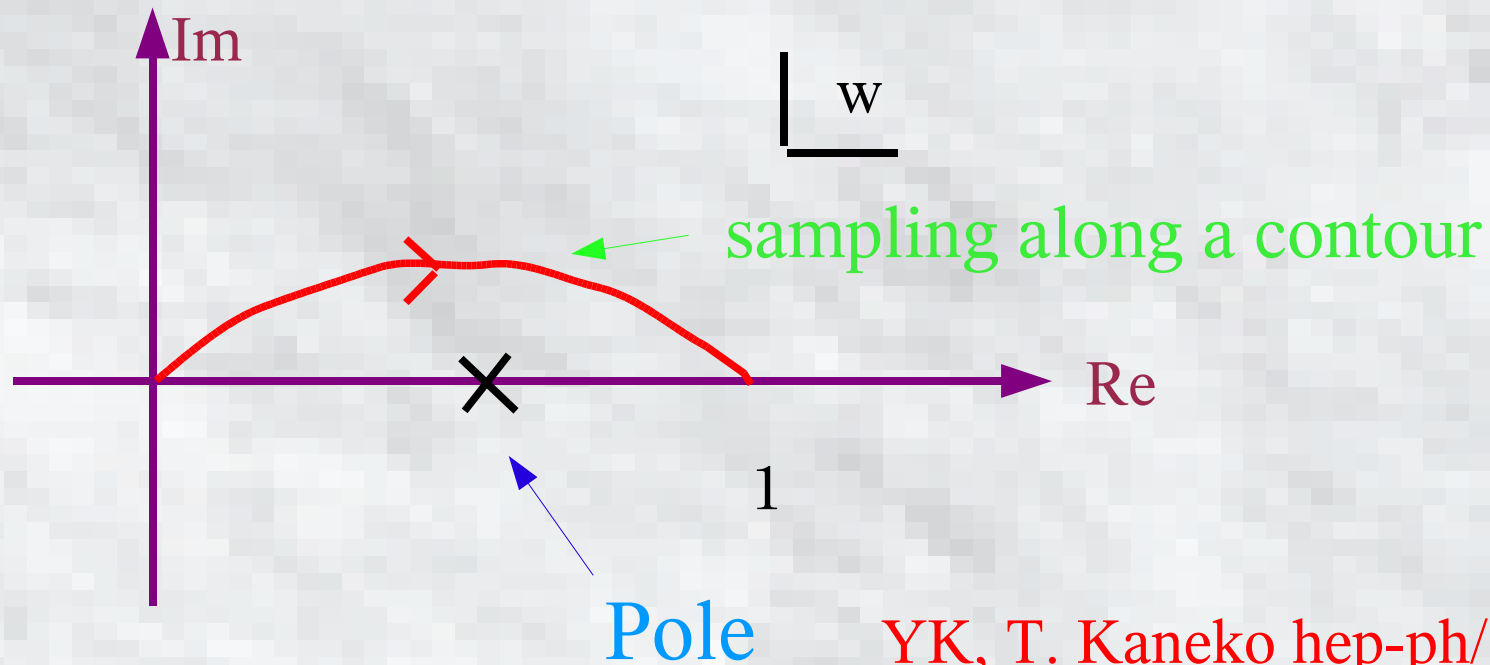
$$\tilde{F}_{j_1,j_2}^{(2)}(z) = \frac{1}{z^{j_1-1}} \sum_{k=0}^{j_1-1} (-1)^k {}_{j_1-1}C_k \tilde{F}_{1,j_2-k}^{(2)}(z)$$



Numerical test for Infrared finite case

$$\begin{aligned} J_4 &= \left(4\pi\mu^2\right)^{-\varepsilon_{IR}} \int_0^1 dr \, r^{-1+n_x+\varepsilon_{IR}} (1-r)^{-1+n_y+n_z+\varepsilon_{IR}} \\ &\times \int_0^1 dv \int_0^1 dw \frac{w^{n_y} (1-w)^{n_z} v^{n_x}}{(-s \, v(1-w) - t \, (1-v)w - i0)^{2-\varepsilon_{IR}}} \\ &= \left(4\pi\mu^2\right)^{-\varepsilon_{IR}} B(n_x + \varepsilon_{IR}, n_y + n_z + \varepsilon_{IR}) \\ &\times \int_0^1 dv \int_0^1 dw \frac{w^{n_y} (1-w)^{n_z} v^{n_x}}{(-s \, v(1-w) - t \, (1-v)w - i0)^{2-\varepsilon_{IR}}} \end{aligned}$$

Numerical contour integral



$$J_4(s, t; 0, 0, 0, 0; n_x, n_y, n_z)$$

n_x	n_y	n_z	real/imag.	analytic	NCI
1	2	3	real	-2.15298×10^{-9}	-2.15297×10^{-9}
			imag.	-2.78647×10^{-9}	-2.78650×10^{-9}
2	0	2	real	9.74570×10^{-9}	9.74572×10^{-9}
			imag.	-3.22229×10^{-8}	-3.22230×10^{-8}

$$J_4(s, t; p_1^2, 0, 0, 0; n_x, n_y, n_z)$$

n_x	n_y	n_z	real/imag.	analytic	NCI
1	2	3	real	-7.88683×10^{-10}	-7.88689×10^{-10}
			imag.	-1.95176×10^{-9}	-1.95176×10^{-9}
2	0	2	real	1.48133×10^{-8}	1.48133×10^{-8}
			imag.	-2.04318×10^{-8}	-2.04318×10^{-8}

$$s=123, t=-200, p_1^2=80$$

Numerical test for IR divergent case

A sector decomposition method can be used.

$$J_4(s, t; 0, 0, 0, 0; 0, 0, 0)$$

	real/imag.	analytic	SD
$1/\varepsilon_{IR}^2$	real	$-1.029686826 \times 10^{-6}$	$-1.029686826 \times 10^{-6}$
	imag.	0	$\mathcal{O}(10^{-16})$
$1/\varepsilon_{IR}$	real	$-5.205325212 \times 10^{-6}$	$-5.205325212 \times 10^{-6}$
	imag.	$1.617428283 \times 10^{-6}$	$1.617428283 \times 10^{-6}$
ε_{IR}^0	real	$-9.739160873 \times 10^{-6}$	$-9.739160872 \times 10^{-6}$
	imag.	$8.569648363 \times 10^{-6}$	$8.569648363 \times 10^{-6}$

$$s=123, t=-200$$

IR cancellation test

ex. Prompt photon production

@ one phase point

$$\delta = a_2/\epsilon_{\text{IR}}^2 + a_1/\epsilon_{\text{IR}} + a_0$$

BOX

$$a_1 = -18601.9993715016$$

$$a_2 = -3793.95539013131$$

S/C+vertex-(terms included in PDF)

$$a_1 = 18601.9993714494$$

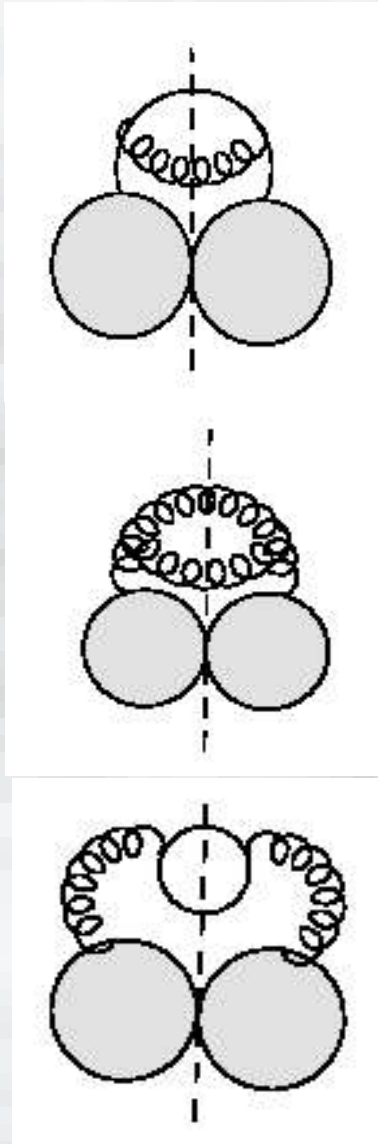
$$a_2 = 3793.95539013130$$

$$1/\epsilon_{\text{IR}}^2 : \text{Soft/Coll.} + \text{Loop} \sim O(10^{-15})$$

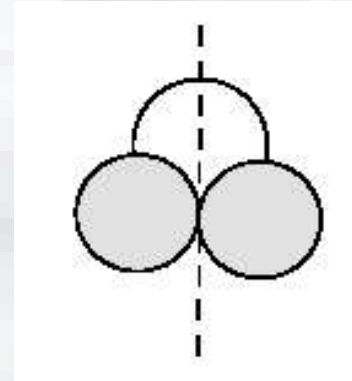
$$1/\epsilon_{\text{IR}} : \text{Soft/Coll.} + \text{Loop} \sim O(10^{-12})$$

Soft/Collinear Approximation

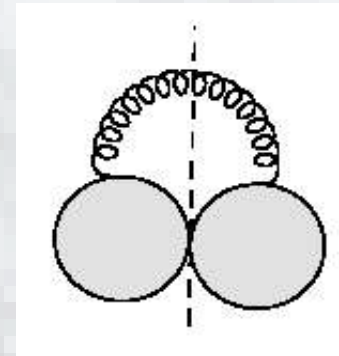
In axial gauge,



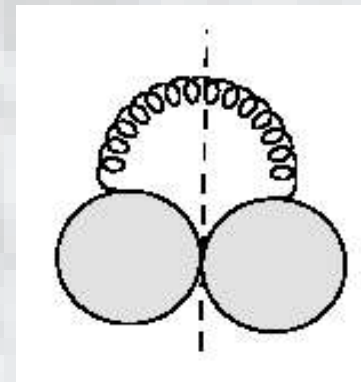
$$= f_{q \rightarrow qg} \times$$



$$= f_{g \rightarrow gg} \times$$

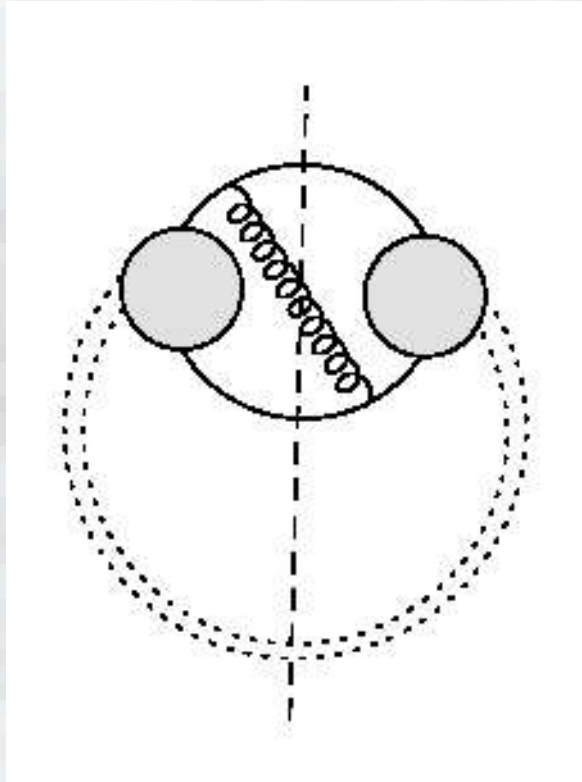


$$= f_{g \rightarrow qq} \times$$

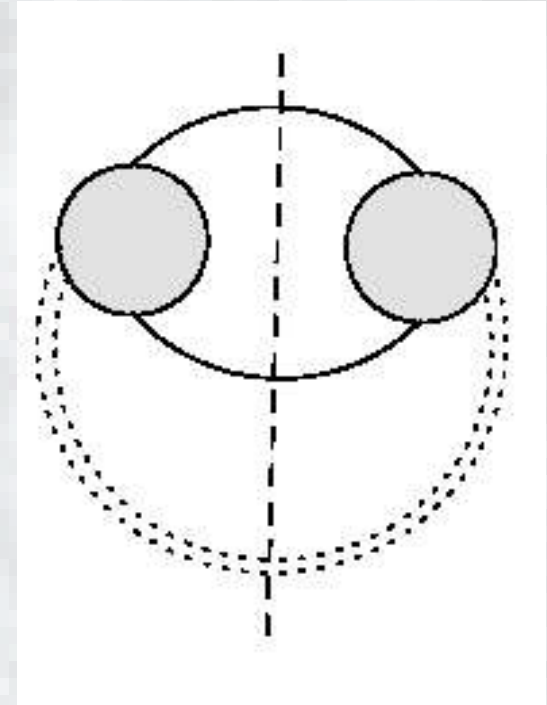


Soft/Collinear Approximation

In axial gauge,



$$= f_{q \rightarrow qg}^{int} \times$$



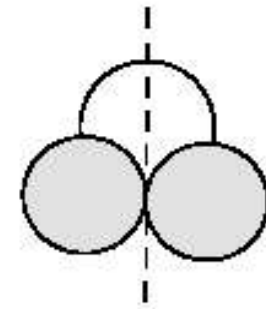
$$\begin{array}{ll} f_{q \rightarrow qg}^{int} = 0 & \text{in collinear region} \\ \neq 0 & \text{in soft region} \end{array}$$

All f 's will be implemented in the system soon.

Soft/Collinear Approximation

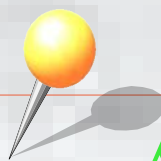
$$f_{q_{in} \rightarrow q_{in} g_{out}}$$

$$\begin{aligned} \sigma_{coll} &= \frac{1}{(2p_1^0)(2p_2^0)v_{rel}} \int_{\Omega_{full}} d\Phi_{N+1}^{(d)} f_{q \rightarrow qg} \times \left[\text{Diagram} \right] \\ &= \left(\frac{s}{4\pi\mu^2} \right)^{\varepsilon_{IR}} \frac{B(\varepsilon_{IR}, \varepsilon_{IR})}{2\Gamma(1+\varepsilon_{IR})} f_c \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) P(x) \left(\frac{1-x}{x} \right)^{2\varepsilon_{IR}} \\ &= \sigma_0(s) \frac{\alpha_s}{2\pi} f_c \left[\frac{1}{\varepsilon_{IR}^2} + \frac{2L-3}{2\varepsilon_{IR}} - \frac{\pi^2}{4} + \frac{L^2}{2} \right] \\ &\quad + \int_0^1 dx \sigma_0(xs) \phi(x, \varepsilon_{IR}) \\ &\quad + f_c \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) \left[L \frac{1+x^2}{(1-x)_+} + \frac{(1+x^2) \ln(1-x)}{(1-x)_+} - \frac{1+x^2}{1-x} \ln x \right] \end{aligned}$$



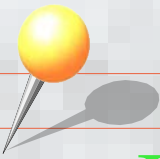
$$\begin{aligned} \phi(x, \varepsilon_{IR}) &= \frac{1}{\varepsilon_{IR}} f_c \frac{\alpha_s}{2\pi} P(x) = \frac{1}{\varepsilon_{IR}} f_c \frac{\alpha_s}{2\pi} \frac{1+x^2}{(1-x)_+} \\ L &= \ln(s/\mu^2). \end{aligned}$$

● Status of GRACE NLO Generator



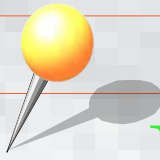
Already tested

- + Drell-Yan process
- + W production
- + Prompt photon



Under development

- + V+1 jet



Future plan

- + V+2 jets,.....
- + VV+jet, 2jets,.....

Summary

- (1) Automatic loop calculation in ELWK is established.
- (2) QCD-NLO Matrix Elements
 - Automatic generation by GRACE
- (3) Loop integral
 - Numerical loop-integration library
- (4) Soft/Collinear treatment
 - Basic tools of NLL-PS is available (under implementation)
 - LL-subtraction method
 - General calculation recipe
- (5) Application
 - Drell-Yan process/w production
 - prompt photon
 - $V+1\text{jet}, 2\text{jet}, VV+1\text{ jet}$