
Top Threshold Physics ... and Other Applications

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thanks to T. Teubner, A. Manohar, I. Stewart, C. Reisser, C. Farrell, P. Ruiz-Femenia



Outline

- Physics at the Top Threshold at the ILC
- Theoretical Issues
- Effective Theory: stable top quarks
- Effective Theory: unstable top quarks
- Status of Computations $\sigma(e^+e^- \rightarrow t\bar{t})$
- Other Applications:
 - $e^+e^- \rightarrow t\bar{t}H$
 - $e^+e^- \rightarrow \tilde{q}\tilde{q}$



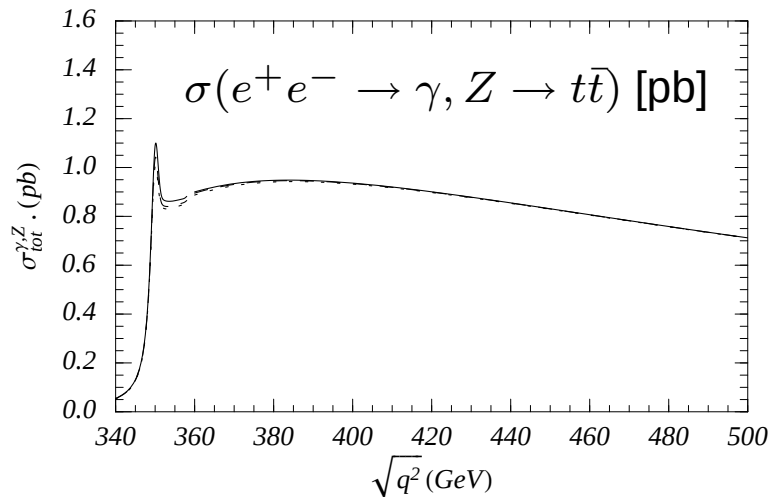
Top Physics and the ILC

- e^+e^- collider: $E_{\text{cm}} = 350 \text{ GeV} - 1 \text{ TeV}$
- Luminosity: $10^{34} - 10^{35} \text{ cm}^{-2} \text{ s}^{-1} \rightarrow 100 - 1000 \text{ fb}^{-1} / \text{year}$
 - LC $\sim 10^5 \text{ } t\bar{t}$ pairs $[\sigma_{\text{tot}} < 1 \text{ pb}] \text{ } (e^+e^- \rightarrow t\bar{t})$
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- Initial state tunable and very well known
 - Centre of mass energy variable $\rightarrow$ threshold & continuum



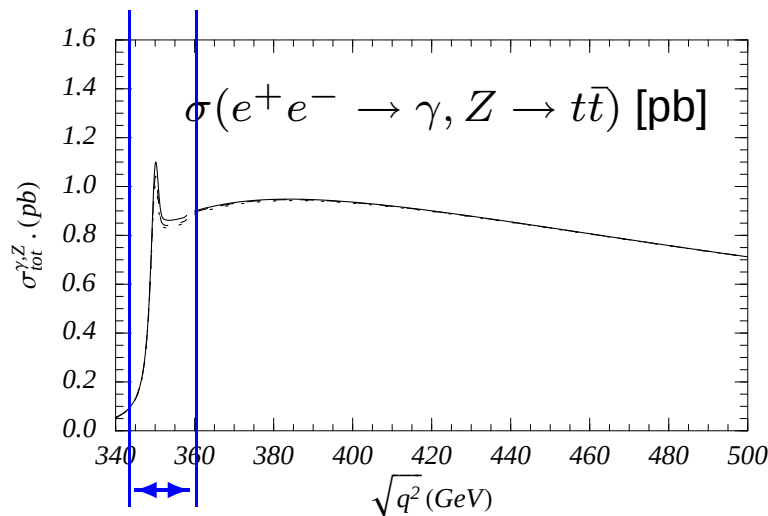
(Luminosity spectrum ignored)

- Electron beam polarizable: $\rightarrow e^-$: 80%
- $\gamma\gamma, \gamma e$ options: $e^+e^- \rightarrow t\bar{t} (^3S_1), \quad \gamma\gamma \rightarrow t\bar{t} (^1S_0)$



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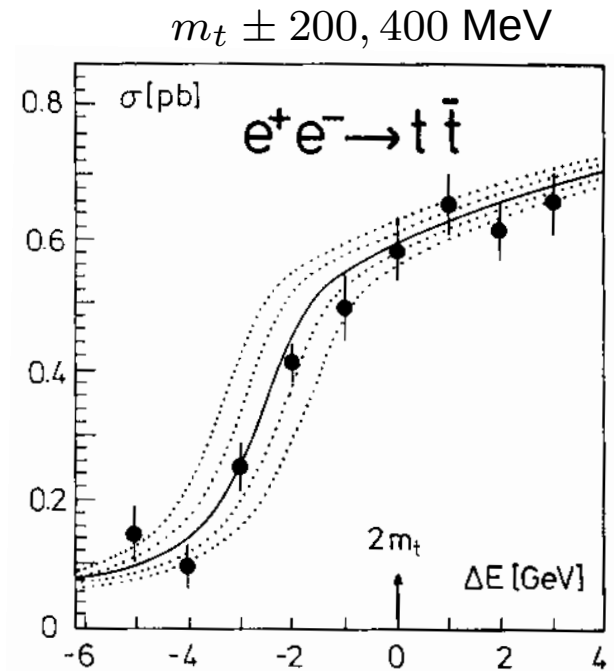


Threshold Measurements

- Top quark mass $\longrightarrow$ threshold scan

- ▷ count number of $t\bar{t}$ events
- ▷ color singlet state
- ▷ background is non-resonant
- ▷ measured mass is unambiguous

$\longrightarrow$ threshold masses (e.g. m_t^{1S})
to avoid artificial perturb. ambiguities

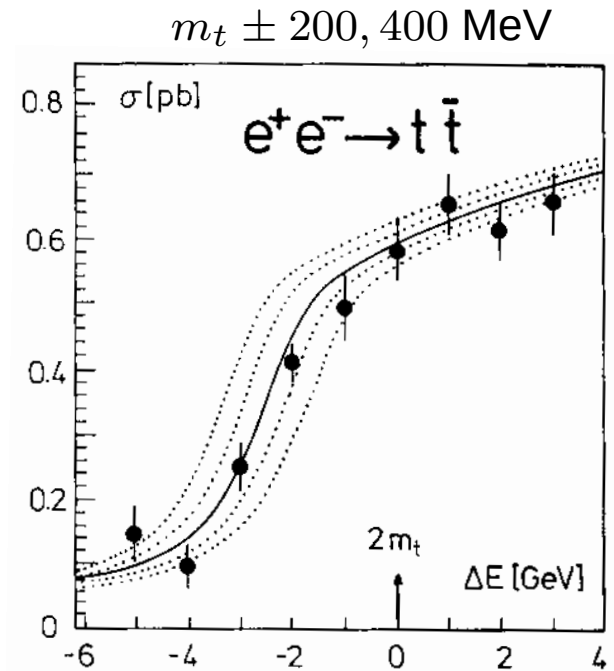


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Simulations

$\mathcal{L} = 300 \text{ fb}^{-1}$, 9 + 1 scan points

Peralta, Martinez, Miquel

$$(\delta m_t)^{\text{stat}} \sim 20 \text{ MeV}$$

$$(\delta \lambda_t / \lambda_t)^{\text{stat}} = 15 - 50\%$$

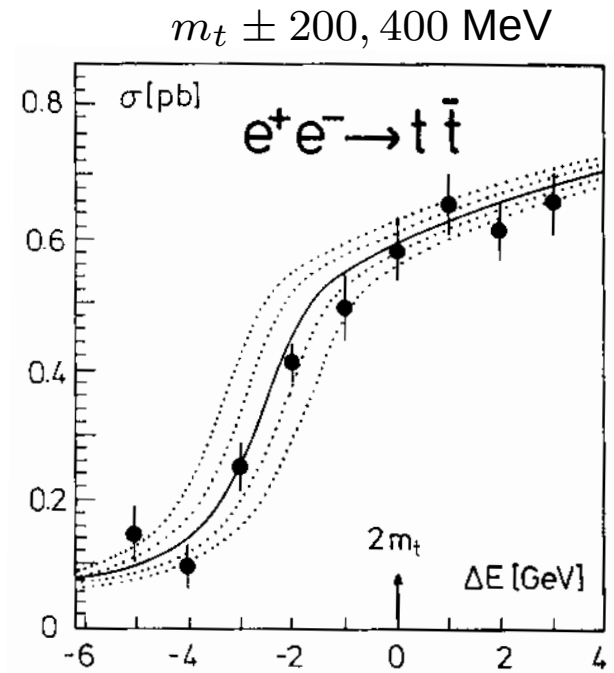
$$(\delta \alpha_s(M_Z))^{\text{stat}} = 0.001$$

$$(\delta \Gamma_t)^{\text{stat}} = 50 \text{ MeV}$$



Threshold Measurements

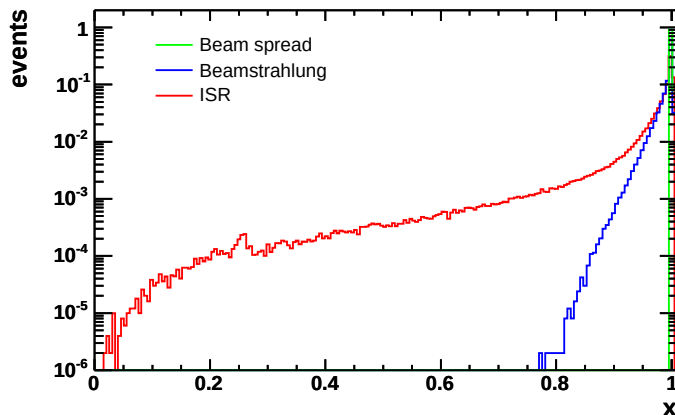
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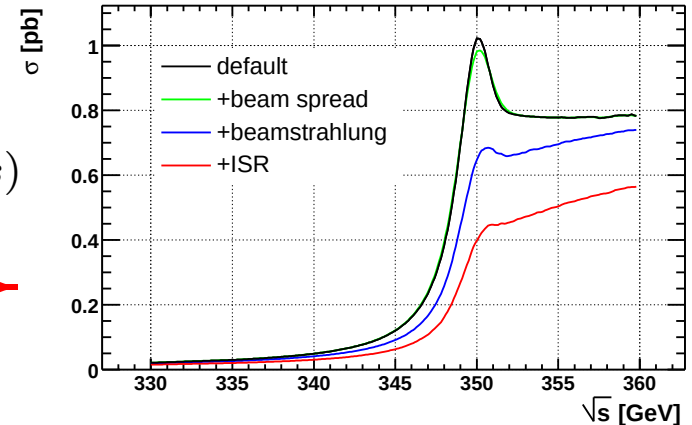
Simulations

Influence of Luminosity spectrum

Boogert



$$\sigma(s) = \int_0^1 dx L(x) \sigma^0(x^2 s)$$

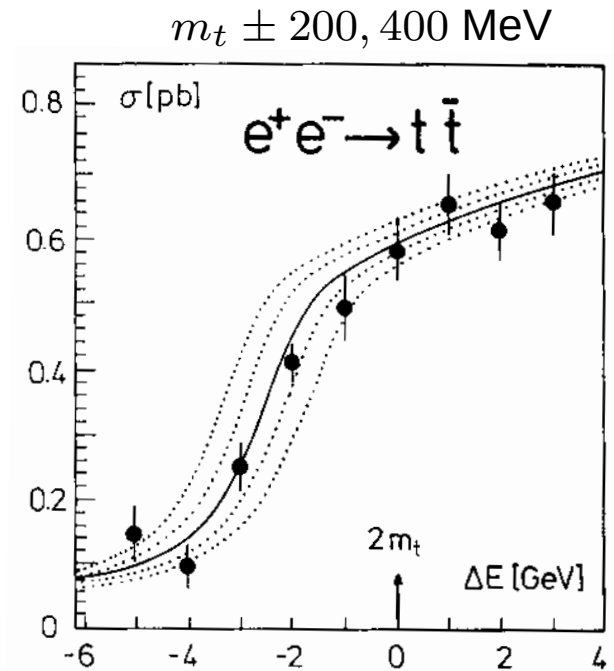


Threshold Measurements

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to avoid artificial perturb. ambiguities



Simulations & theory goal:

$$(\delta\sigma/\sigma)^{\text{theo}} \leq 3\%$$

$$(\delta m_t)^{\text{stat}} \sim 20 \text{ MeV}$$

$$(\delta m_t)^{\text{syst}} = 50 \text{ MeV}$$

$$(\delta m_t)^{\text{theo}} = 100 \text{ MeV}$$

$$(\delta\lambda_t/\lambda_t)^{\text{stat}} = 15 - 50\%$$

$$(\delta\lambda_t/\lambda_t)^{\text{syst}} = ?$$

$$(\delta\lambda_t/\lambda_t)^{\text{theo}} \sim ?$$

$$(\delta\alpha_s(M_Z))^{\text{stat}} = 0.001$$

$$(\delta\alpha_s(M_Z))^{\text{syst}} = 0.002$$

$$(\delta\alpha_s(M_Z))^{\text{theo}} \sim ?$$

$$(\delta\Gamma_t)^{\text{stat}} = 50 \text{ MeV}$$

$$(\delta\Gamma_t)^{\text{syst}} = 15 \text{ MeV}$$

$$(\delta\Gamma_t)^{\text{theo}} \sim ?$$



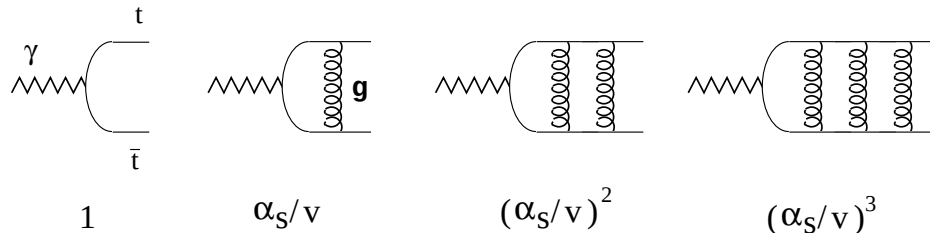
Theory Issues

$$m_t \text{ (hard)} \gg p \sim m_t v \text{ (soft)} \gg E \sim m_t v^2 \text{ (ultrasoft)}$$

“multi-scale problem”

- perturbation theory in α_s breaks down

$$(\alpha_s/v)^n$$



“Coulomb singularities”

→ Schrödinger Equation

- perturbation theory in α_s breaks down → large logs $(\alpha_s \ln v)^n$

$$m_t = 175 \text{ GeV}, \quad p \sim 25 \text{ GeV}, \quad E \sim 4 \text{ GeV} \quad \Rightarrow \quad \ln \left(\frac{m_t^2}{E^2} \right) = 8 \quad \rightarrow \text{RGE's}$$



Theory Issues

- $\Gamma_t \approx 1.5 \text{ GeV} \gg \Lambda_{\text{QCD}} \Rightarrow v = \sqrt{\frac{E}{m}} \rightarrow v_{\text{eff}} = \sqrt{\frac{E + i\Gamma_t}{m}}$
(Fadin, Khoze)

$\Rightarrow m_t \gg p = mv_{\text{eff}} \gg E = mv_{\text{eff}}^2 \gtrsim \Lambda_{\text{QCD}}$ always true !

$\Rightarrow$ top threshold entirely perturbative ! $\rightarrow$ “Schrödinger theory”

- top quarks are always produced off-shell !
 - $\rightarrow$ methods for on-shell production do not apply for
 - theoretical computations
 - experimental analysis

“theory for unstable particles”



EFT - Shopping List

$$\sigma_{t\bar{t}} \sim v \sum_{n,m} \left(\frac{\alpha_s}{v}\right)^n \left(\alpha_s \ln v\right)^m \left[1, \{v, \alpha_s\}, \{v^2, \alpha_s v, \alpha_s^2\}, \dots \right]$$

LL

NLL

NNLL

- Action from **Lagrangian** in terms of field operators
 - ▷ Manifest power counting in v at all times
consistently obey $E \ll p \ll m$ and $\Gamma_t \sim E \sim p^2/m$
 - ▷ Symmetries: gauge, spin, ...
 - ▷ Regulator independence $\rightarrow$ matching in different schemes
- **Renormalization** of the theory $\rightarrow$ **anomalous dimensions,**
summation of $(\alpha_s \ln v)^m$
- **Dimensional regularization**



Different EFT's (for QCD)

Classic "NRQCD": ('92-'95)

Bodwin, Braaten, Lepage

Separate $p > m$ from $p < m$

[Production, Decays, Lattice]

"NRQCD-pNRQCD" ('97-'99)

Pineda, Soto, Brambilla, Vairo

Separate also $p > mv$ from $p < mv$

$[\bar{c}c, \bar{b}b]$

Combined 1- and 2-step matching, running

Pineda ('01)

Power counting not uniform in different parts of the theory

"vNRQCD" ('99)

Luke, Manohar, Rothstein, Stewart, A.H.

Separation of all degrees of freedom at m

$[t\bar{t}]$

1-step matching and running, uniform power counting

Correlation of mv and mv^2

(needed in general for $\Lambda_{\text{QCD}} \ll mv^2$: e. g. c_1 , also in QED)

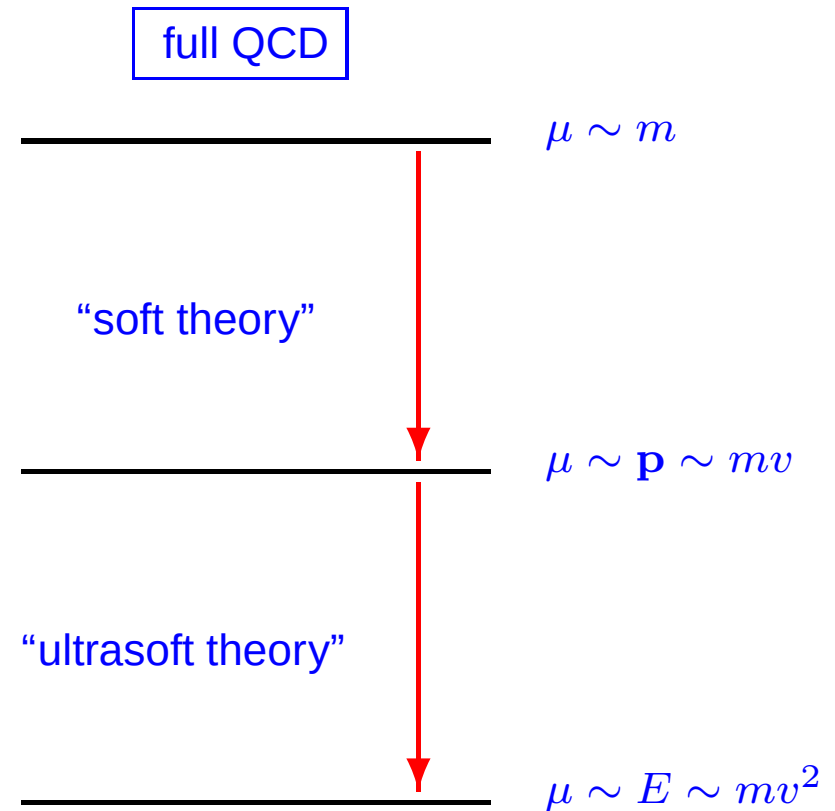


Degrees of Freedom

naive picture:

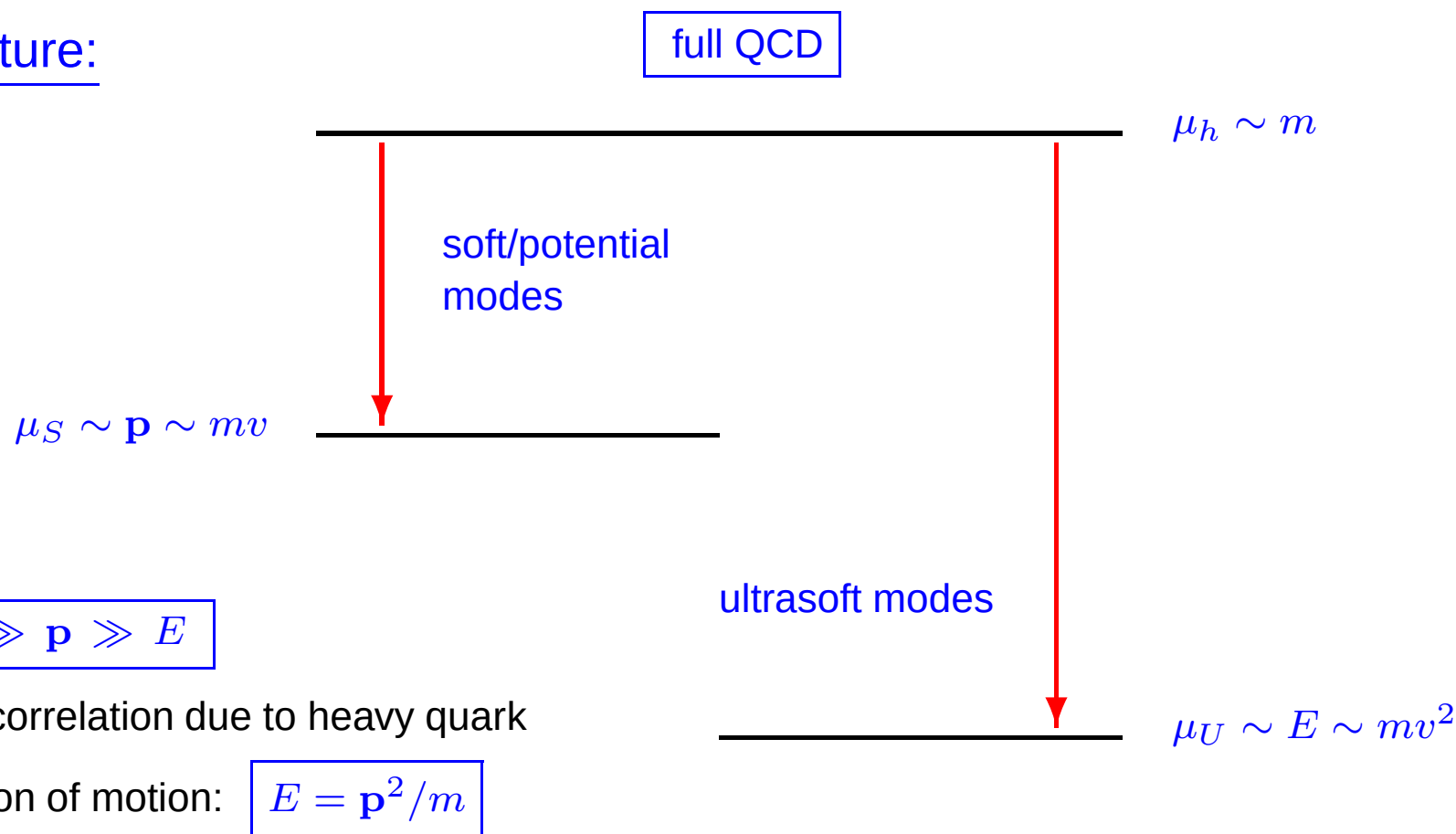
$$m_t \gg \mathbf{p} \gg E$$

- successively integrated out d.o.f.'s
- sequence of effective theories



Degrees of Freedom

proper picture:



- soft and ultrasoft modes exist at the same time
- $\mu_U = \mu_S^2/m \rightarrow \mu_S = mv, \mu_U = mv^2$



Degrees of Freedom

- fields for degrees of freedom that can **resonate** for the quark-antiquark system

potential quarks $\psi_{\mathbf{p}}, \chi_{\mathbf{p}}:$ $(k_0, \mathbf{k}) \sim (mv^2, mv)$

soft gluons $A_q^\mu:$ $(k_0, \mathbf{k}) \sim (mv, mv)$

ultrasoft gluons $A^\mu:$ $(k_0, \mathbf{k}) \sim (mv^2, mv^2)$

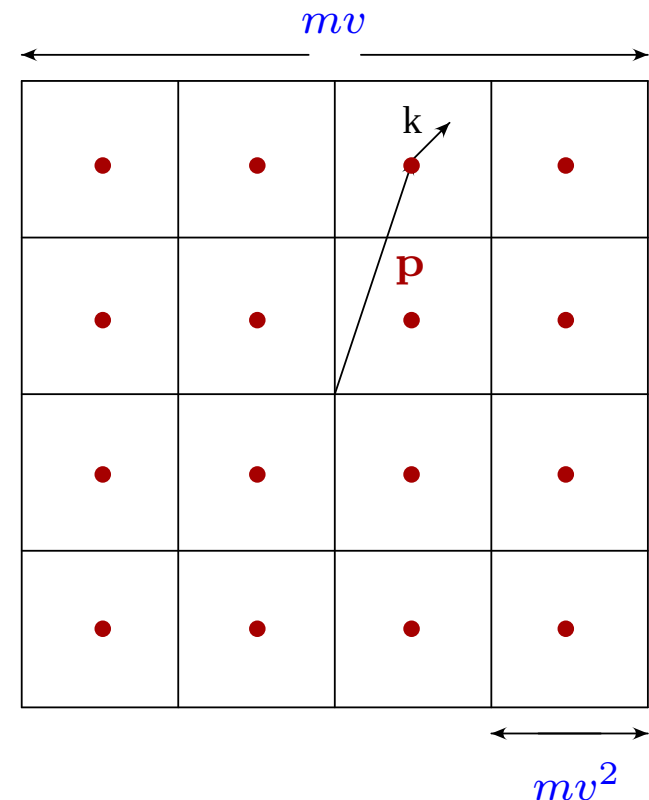
- vNRQCD label formalism:

$$(P^0, \mathbf{P}) = (0, \mathbf{p}) + (k^0, \mathbf{k})$$

soft component
label

ultrasoft component
dynamic variable

$$\psi_{\text{QCD}}(\mathbf{x}) \rightarrow \sum_{\mathbf{p}} e^{i\mathbf{p} \cdot \mathbf{x}} \psi_{\mathbf{p}}(\mathbf{x})$$



vNRQCD (stable quarks)

$$\mathcal{L} = \mathcal{L}_{\text{usoft}} + \mathcal{L}_{\text{potential}} + \mathcal{L}_{\text{soft}}$$

Luke, Manohar, Rothstein, Stewart, A.H.

$$\mathcal{L}_{\text{usoft}} : \text{---} \bullet \text{---} \text{---} \text{---} \text{---}$$

$$\psi_{\mathbf{p}}^\dagger(x) \left\{ iD^0 - \frac{(\mathbf{p} - i\mathbf{D})^2}{2m_t} - \delta m_t \right\} \psi_{\mathbf{p}}(x)$$

$$\mathcal{L}_{\text{potential}} : \text{---} \bullet \text{---}$$

$$\left\{ \frac{V_c(\nu)}{(\mathbf{p} - \mathbf{p}')^2} + \dots \right\} \psi_{\mathbf{p}'}^\dagger \psi_{\mathbf{p}} \chi_{-\mathbf{p}'}^\dagger \chi_{-\mathbf{p}}$$

$$\mathcal{L}_{\text{soft}} : \text{---} \bullet \text{---}$$

$$U_{\mu\nu}(\nu) \psi_{\mathbf{p}'}^\dagger A_q^\mu A_{q'}^\nu \psi_{\mathbf{p}}$$

$$V = \left[\frac{\mathcal{V}_c(\nu)}{\mathbf{k}^2} + \frac{\mathcal{V}_k(\nu)\pi^2}{m|\mathbf{k}|} + \frac{\mathcal{V}_r(\nu)(\mathbf{p}^2 + \mathbf{p}'^2)}{2m^2\mathbf{k}^2} + \frac{\mathcal{V}_2(\nu)}{m^2} + \frac{\mathcal{V}_s(\nu)}{m^2} \mathbf{S}^2 + \frac{\mathcal{V}_\Lambda(\nu)}{m^2} \Lambda + \frac{\mathcal{V}_t(\nu)}{m^2} T \right]$$

$$\mathbf{k} \equiv \mathbf{p} - \mathbf{p}'$$



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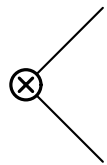
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$$\mathcal{L}_{\text{soft}} : \text{---} \bullet \text{---}$$

$$U_{\mu\nu}(\nu) \psi_{\mathbf{p}'}^\dagger A_q^\mu A_{q'}^\nu \psi_{\mathbf{p}}$$

external currents: (production & annihilation)



$$\mathbf{O}_{\mathbf{p}} = C_{3S_1}(\nu) \cdot (\psi_{\mathbf{p}}^\dagger \boldsymbol{\sigma} \tilde{\chi}_{-\mathbf{p}}^*) + \cdots \quad t\bar{t} (^3S_1)$$



Cross Section at NNLL Order

Schematic:

$$\begin{aligned}\sigma_{\text{tot}} &\propto \text{Im} \left[\int d^4x e^{-i\hat{q}x} \left\langle 0 \left| T \mathbf{O}_{\mathbf{p}}^\dagger(0) \mathbf{O}_{\mathbf{p}'}(x) \right| 0 \right\rangle \right] \\ &\propto \text{Im} \left[C(\nu)^2 G(0, 0, \sqrt{s}) \right]\end{aligned}$$

$$\left(-\frac{\nabla^2}{m} - \frac{\nabla^4}{4m^3} + V(\mathbf{r}) - (\sqrt{s} - 2m - 2\delta m) - i\Gamma_t \right) G(\mathbf{r}, \mathbf{r}') = \delta^{(3)}(\mathbf{r} - \mathbf{r}')$$



potentials fully known
at NNLL order

Pineda,Soto '00-'01
Manohar,Stewart, AHH '99-'02



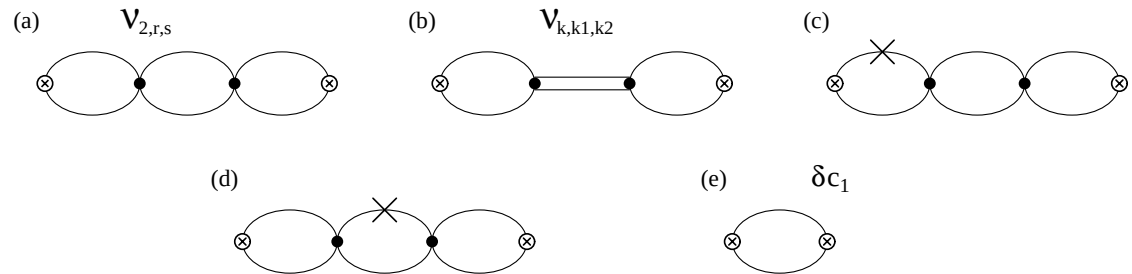
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$$\propto \text{Im} \left[C(\nu)^2 G(0, 0, \sqrt{s}) \right]$$

NLL running of $C(\nu)$



$$\frac{d}{d \ln \nu} \ln C(\nu) = -\frac{\mathcal{V}_c(\nu)}{16\pi^2} \left(\frac{\mathcal{V}_c(\nu)}{4} + \mathcal{V}_2(\nu) + \mathcal{V}_r(\nu) + \mathbf{s}^2 \mathcal{V}_s(\nu) \right) + \frac{\mathcal{V}_k(\nu)}{2}$$

Luke, Manohar, Rothstein '99, Pineda '01
Stewart, AHH '02



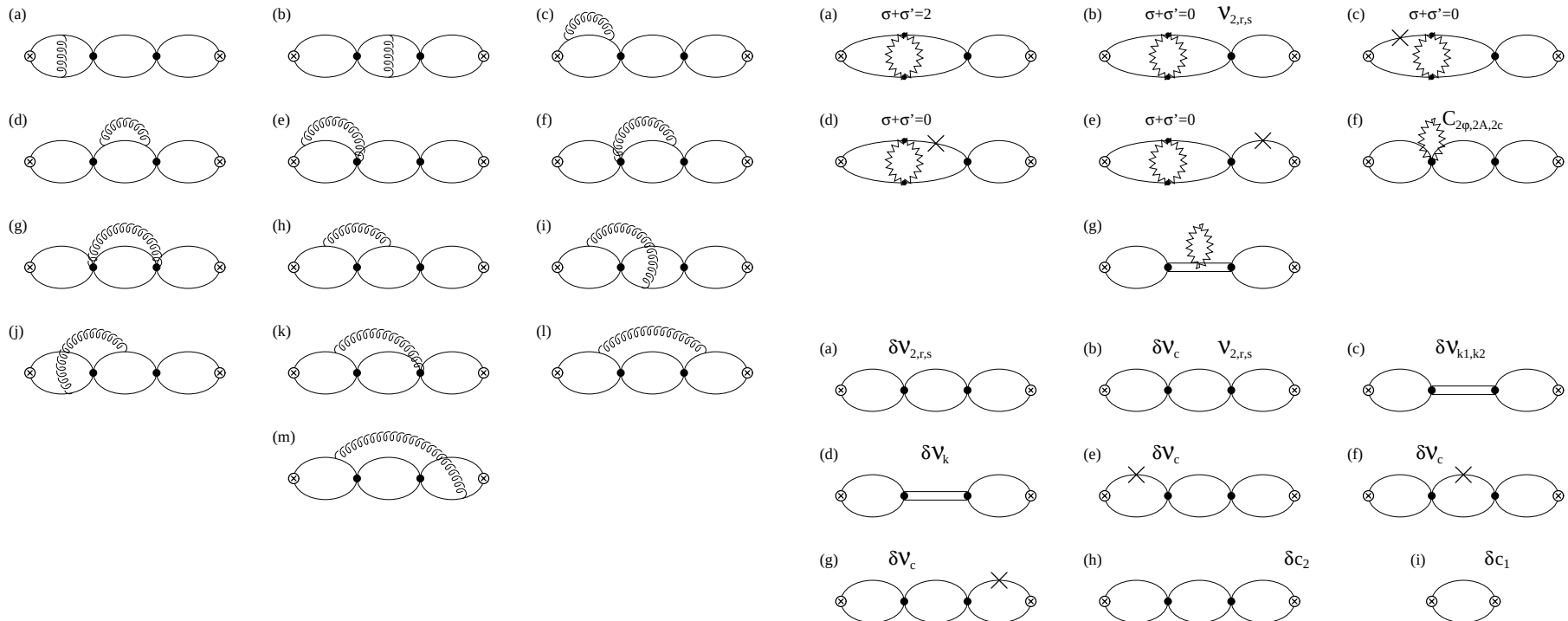
Cross Section at NNLL Order

NNLL running of $C(\nu)$

(a) next-order running of couplings in NLL anom. dim. $\rightarrow$?

(b) genuine next-order diagrams

AHH '03



Cross Section at NNLL Order

NNLL running of $C(\nu)$

- (a) next-order running of couplings in NLL anom. dim. $\rightarrow$?
- (b) genuine next-order diagrams AHH '03

$$\frac{\delta z_c^{\text{NNLL},1}}{\epsilon} = \frac{\text{const}}{\epsilon} + \frac{\text{const}}{\epsilon} \times \ln \left(\frac{m_t \mu_U}{\mu_S^2} \right)$$

$\Rightarrow$ consistency under renormalization requires:

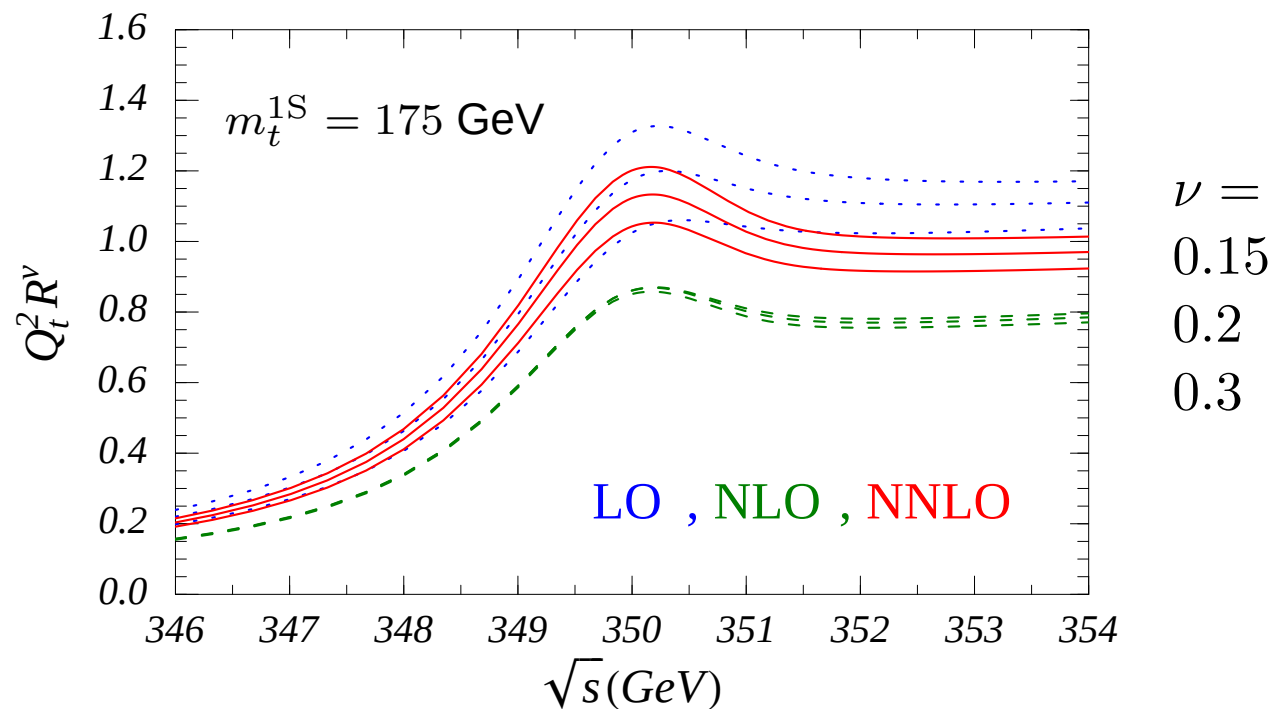
$$\mu_U = \mu_S^2 / m_t$$



Cross Section at NNLL Order

A.H.

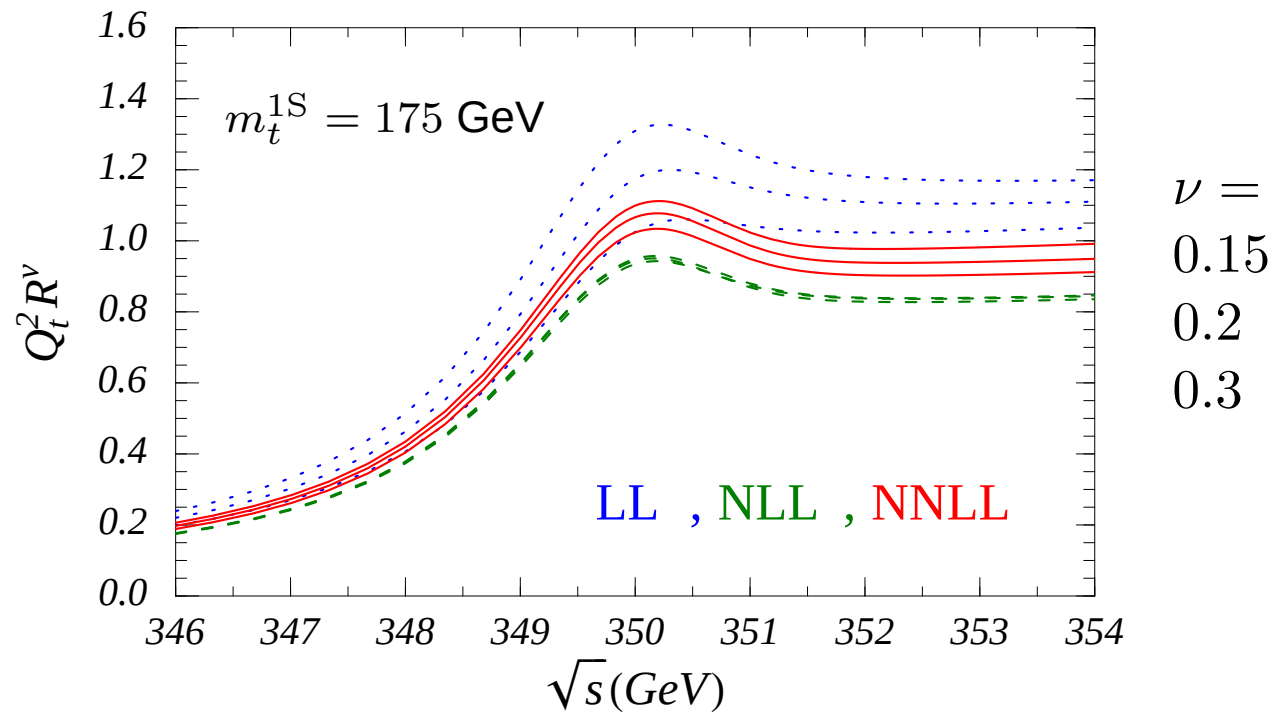
1S mass - fixed order



Cross Section at NNLL Order

1S mass - RG-improved, with NNLL non-mixing terms

A.H.



- expansion shows better convergence
 - theory error: $\delta\sigma_{t\bar{t}}/\sigma_{t\bar{t}} \sim \pm 6\%$
- full NNLL running of $C(\nu)$ required → w.i.p.



vNRQCD (unstable quarks)

“inclusive treatment”

- ⇒ Optical Theory: effective complex indices of refraction for absorptive processes
- ⇒ vNRQCD: contributions from *real Wb final states* included in EFT matching conditions to QCD+ew. theory (=SM)
- complex matching conditions
 - effective Lagrangian non-hermitian
 - total rates through the optical theorem
 - power counting maintained

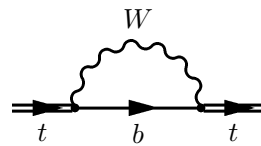


vNRQCD (unstable quarks)

quark bilinears: $(\rightarrow \mathcal{L}_{\text{usoft}})$

time dilatation
correction

$$\text{Im}\Sigma_t = \frac{1}{2}\Gamma_t$$



$$= i\Sigma_t \quad \Rightarrow \quad \delta\mathcal{L} = \psi_{\mathbf{p}}^\dagger \left[i\frac{\Gamma_t}{2} \left(1 - \frac{\mathbf{p}^2}{2m_t^2} \right) \right] \psi_{\mathbf{p}}$$

—•—

- power counting: $\Gamma_t \propto m_t g^2 \sim m_t v^2 \Rightarrow \boxed{g \sim g' \sim v \sim \alpha_s}$
- finite lifetime is LL order, $E \rightarrow E + i\Gamma_t$ Fadin, Khoze

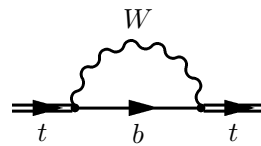


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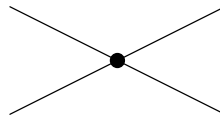
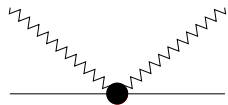
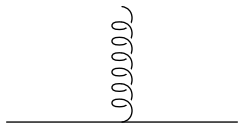
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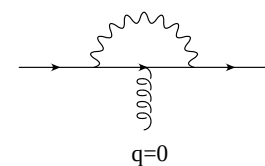


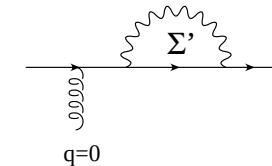
$$= i\Sigma_t \quad \Rightarrow \quad \delta\mathcal{L} = \psi_{\mathbf{p}}^\dagger \left[i\frac{\Gamma_t}{2} \left(1 - \frac{\mathbf{p}^2}{2m_t^2} \right) \right] \psi_{\mathbf{p}}$$



gluon interactions & potentials:





$$+ \text{[diagram with } \Sigma' \text{]} = 0$$


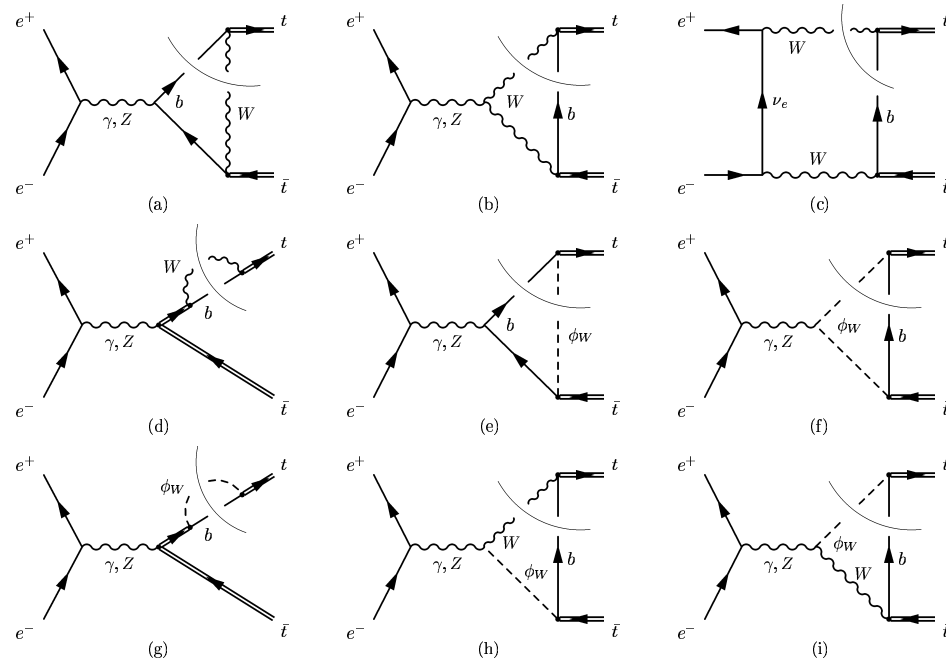
- electroweak corrections either beyond NNLL order or vanish due to gauge cancellations

→ ultrasoft gluon interference effects vanish at NLL [Khoze etal., Melnikov etal.]
and NNLL order (new !)



ν NRQCD (unstable quarks)

Currents:



real virtual electroweak
corrections

bW^+ and $\bar{b}W^-$ cuts

$$\mathbf{O}_p = \left[C^{\text{LL}} + C^{\text{NLL}} + C^{\text{NNLL}} + iC_{\text{abs}}^{\text{NNLL}} \dots \right] \cdot \left(\begin{array}{cc} e^+ & t \\ e^- & \bar{t} \end{array} \right) + \dots$$



vNRQCD (unstable quarks)

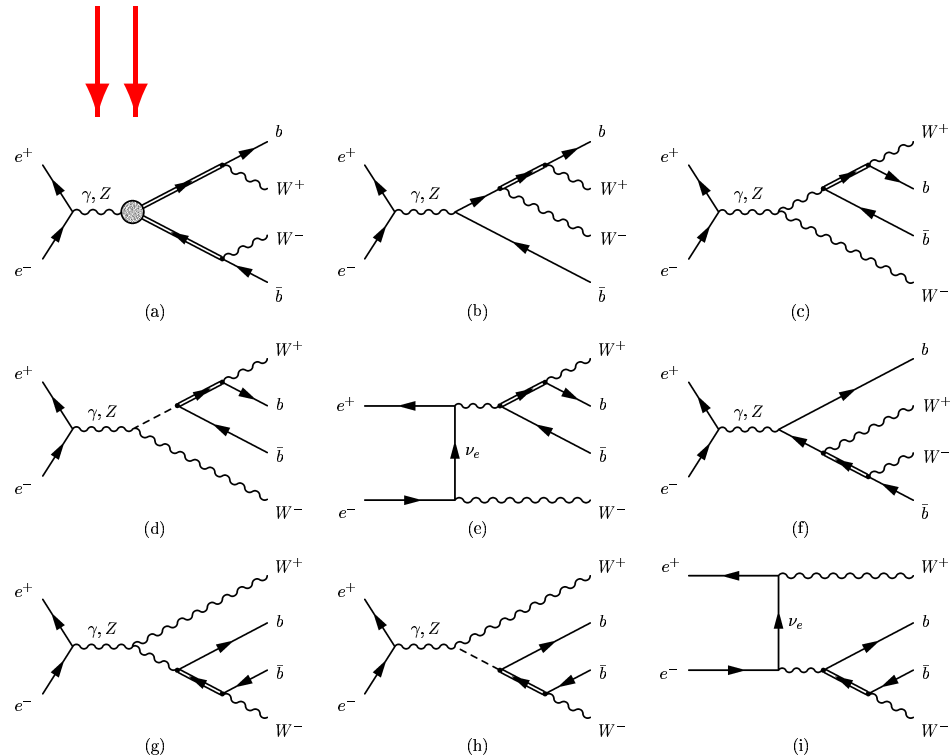
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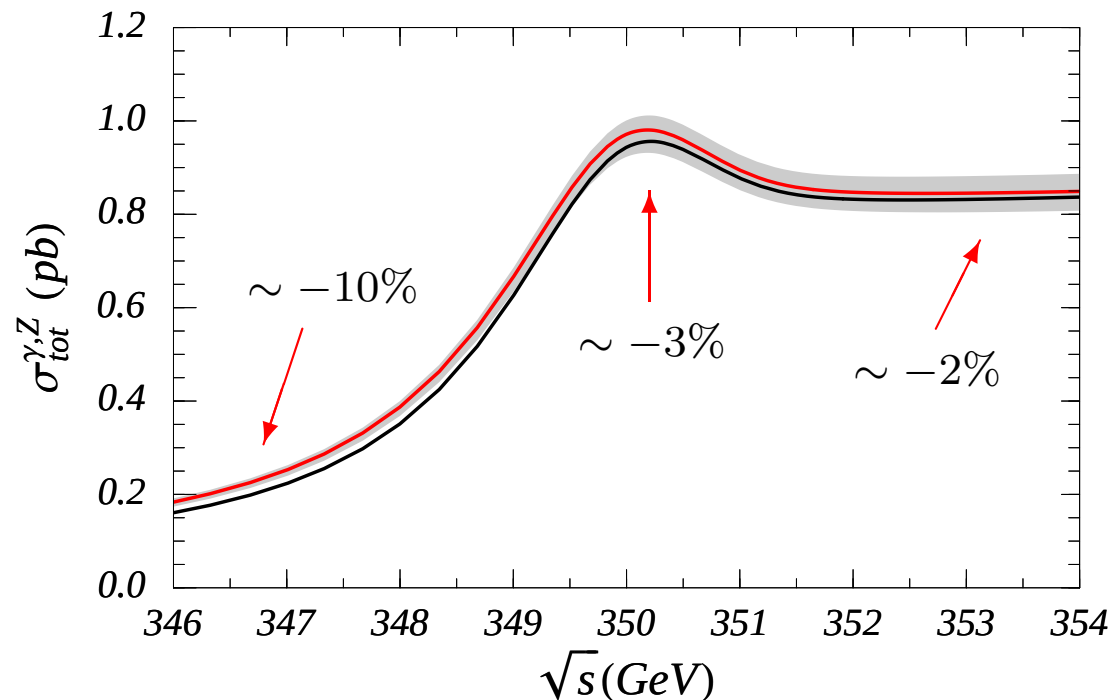
$$\sigma_{\text{tot}} \propto \text{Im} [C(\nu)^2 G(0, 0, \sqrt{s} + i\Gamma_t)]$$

- accounts for **irreducible interference** contributions:

resonant $\leftrightarrow$ non-resonant
 $W^+ W^- b \bar{b}$ final states



Total Cross Section



Reisser, AHH

- also included summation of phase space logs $\sim (\alpha_s \ln v)^n$
- finite lifetime corrections comparable to NNLL QCD corrections
- shift in the peak position: 30 – 50 MeV $(\delta m_t^{\text{ex}} \approx 50 \text{ MeV})$



What's left to do . . .

NNLL Total Cross Section:

(A) QCD corrections (RGE-improved)

- Current uncertainty: $(\delta\sigma_{t\bar{t}}/\sigma_{t\bar{t}})^{\text{QCD}} \sim \pm 6\%$
→ full NLL running of potentials

(B) electroweak corrections

- experimental cuts
- QED effects: beam effects $\leftrightarrow$ QED dynamics

Differential Cross Sections, (multi-quark final states): → efficiencies

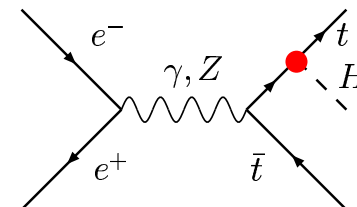
→ more involved treatment of unstable particles



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ top-Yukawa coupling



- Theory Status: $\sigma(e^+e^- \rightarrow t\bar{t}H)$

Born ✓

1-loop ew. ✓

$\mathcal{O}(\alpha_s)$ fixed-order ✓

NLL large- E_H QCD endpoint corrections ✓

[Gaemers et al., Djouadi et al.]

[Denner et al., Belanger et al., You et al.]

[Dittmaier et al., Dawson et al.]

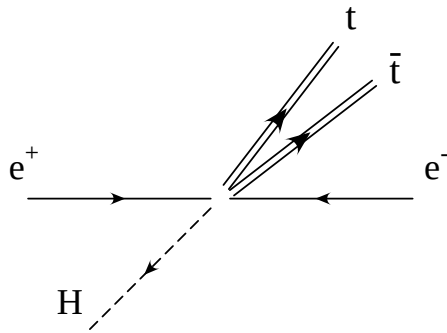
[C. Farrell, AHH]



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ region of large Higgs energy



→ $t\bar{t}$ collinear

→ QCD effects localized in $t\bar{t}$ system

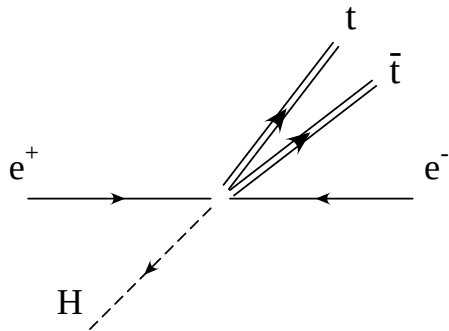
⇒ $t\bar{t}$ dynamics non-relativistic



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

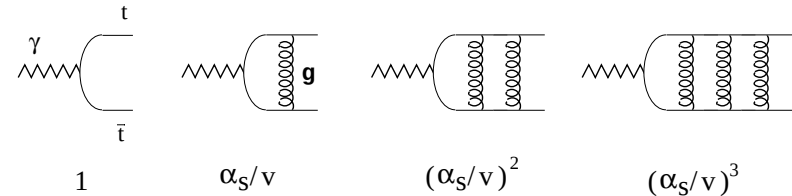
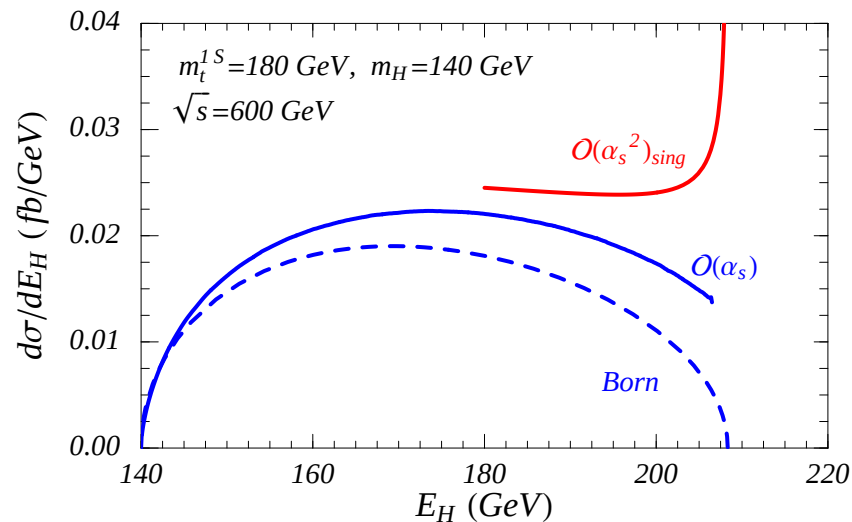
→ region of large Higgs energy



→ $t\bar{t}$ collinear

→ QCD effects localized in $t\bar{t}$ system

⇒ $t\bar{t}$ dynamics non-relativistic



→ singularities: $\sim (\alpha_s/v)^n$,
 $\sim (\alpha_s \ln v)^n$

→ fixed order expansion breaks down

⇒ summation of singular terms

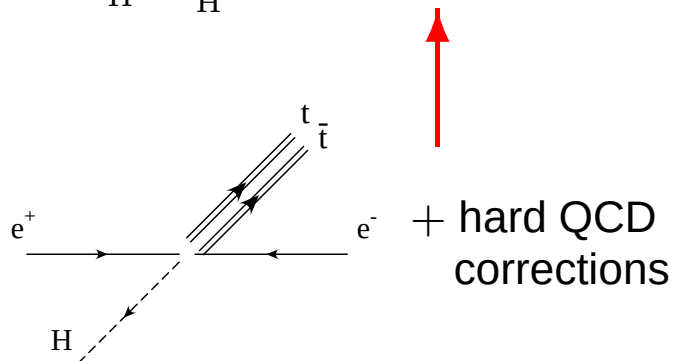


Other Applications

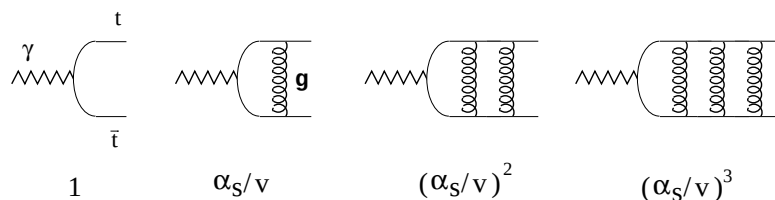
$$e^+e^- \rightarrow t\bar{t}H$$

→ factorization formula

$$\left(\frac{d\sigma}{dE_H}\right)_{E_H \approx E_H^{\max}} \sim C^2(\mu, \sqrt{s}, m_t, m_H) \times \text{Im}[G(0, 0, v, \mu)]$$



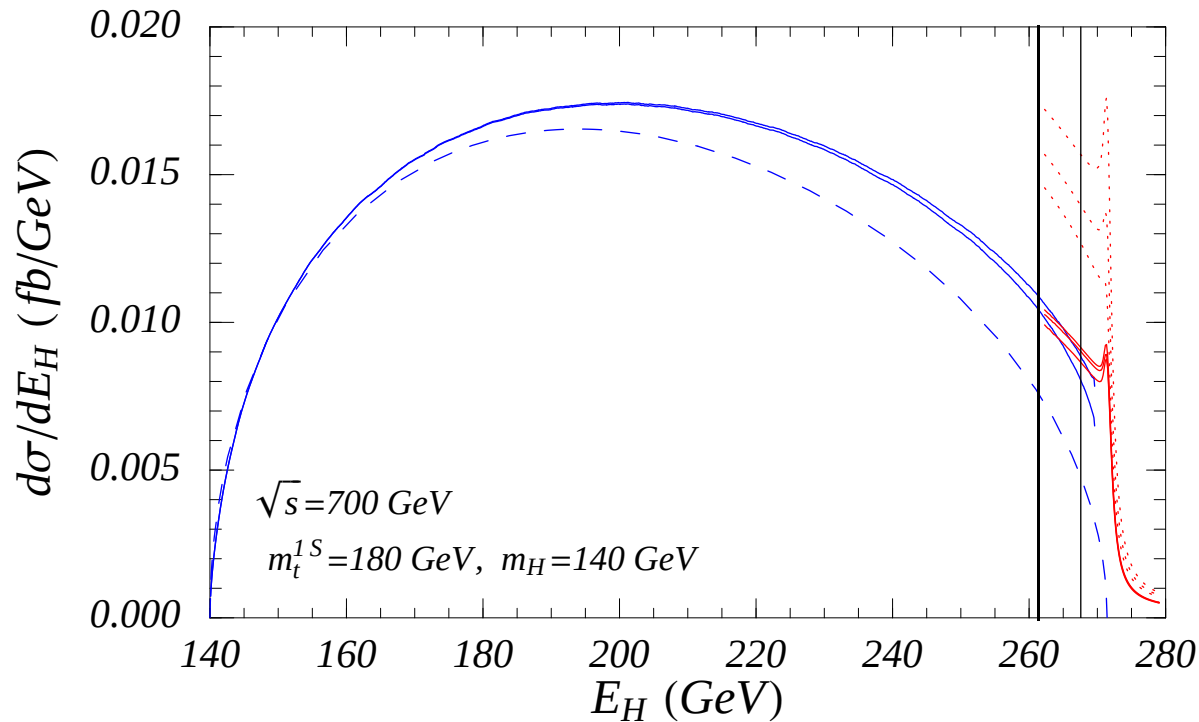
soft & ultrasoft QCD corrections



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ NLL Higgs energy spectrum



Farrell, AHH

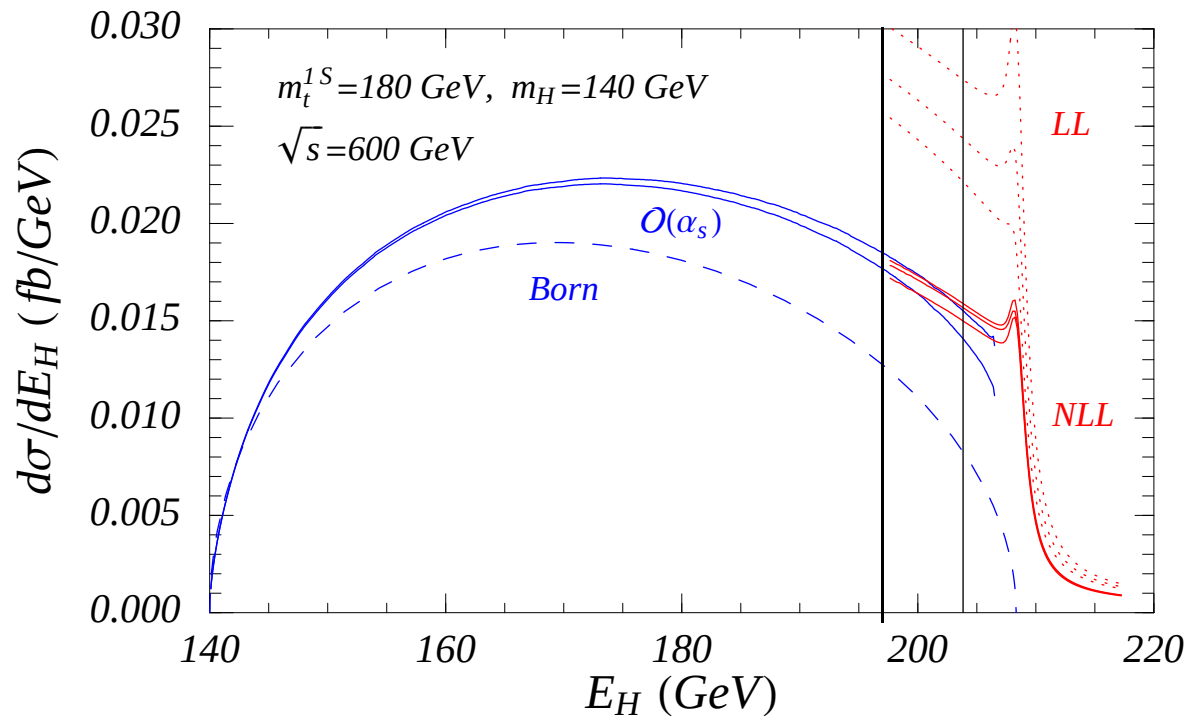
- large E_H endpoint regions increases for smaller $\sqrt{s}$ / larger m_H



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ NLL Higgs energy spectrum



Farrell, AHH

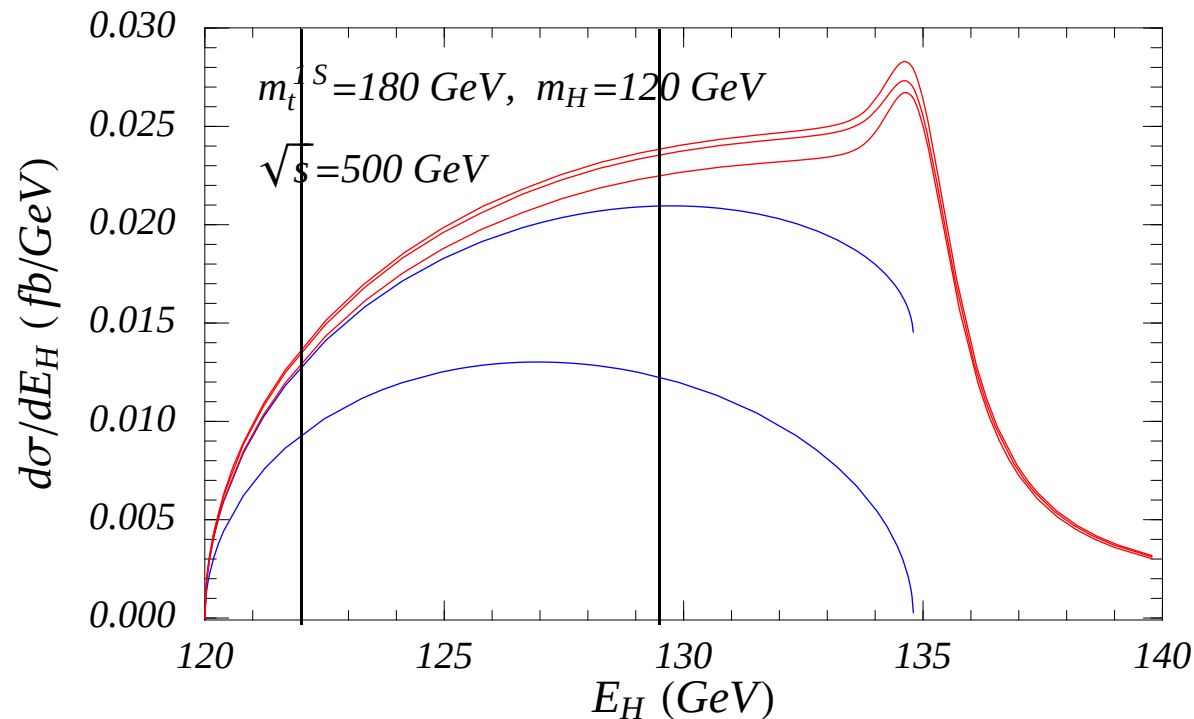
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Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ NLL Higgs energy spectrum



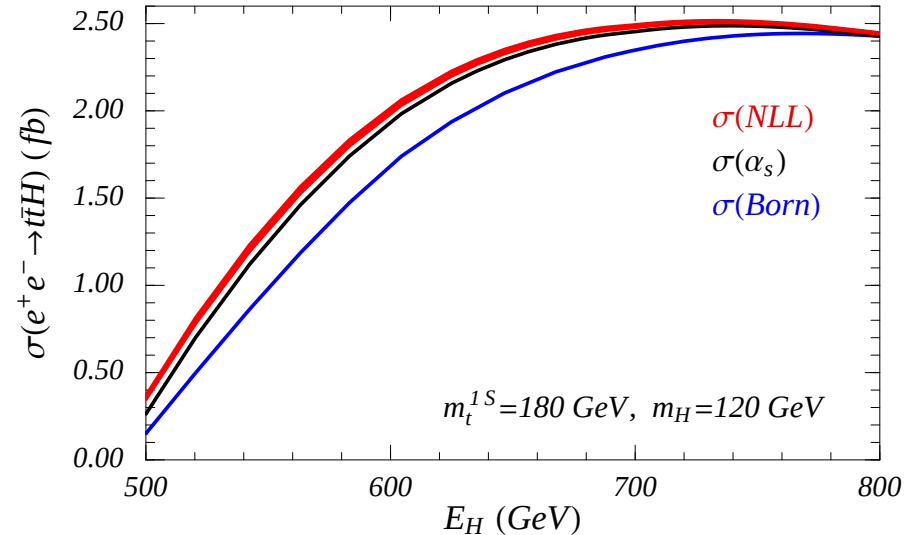
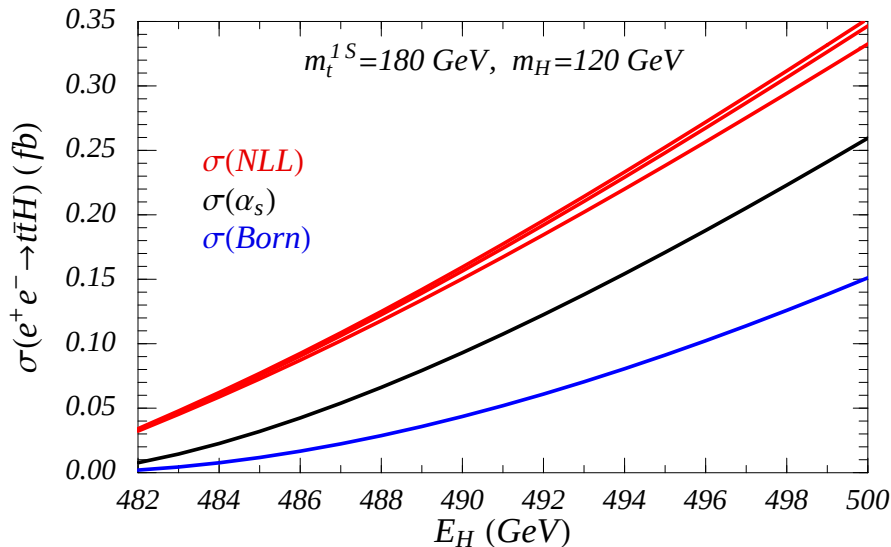
- large E_H endpoint regions increases for smaller $\sqrt{s}$ / larger m_H



Other Applications

$$e^+e^- \rightarrow t\bar{t}H$$

→ total cross section



- significant enhancement from summation of $(\alpha_s/v)^n$, $(\alpha_s \ln v)^n$ singularities
- essential for realistic studies for ILC (phase I)

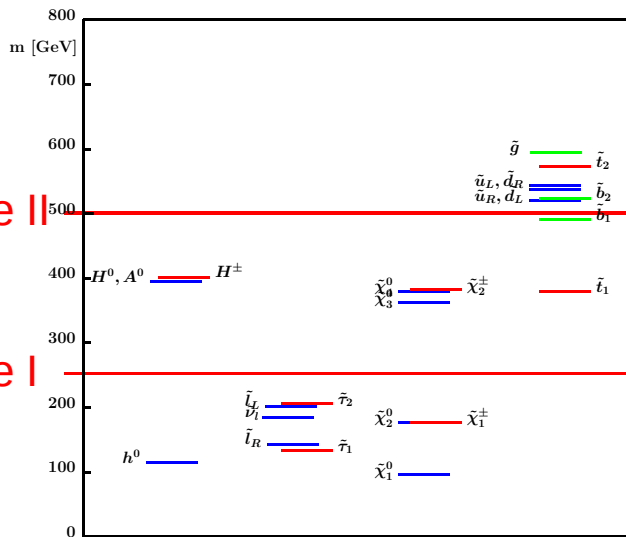


Other Applications

$$e^+e^- \rightarrow \tilde{q}\tilde{q}^*$$

- in many models for SUSY breaking squark pair production is possible at the ILC

SPS 1a

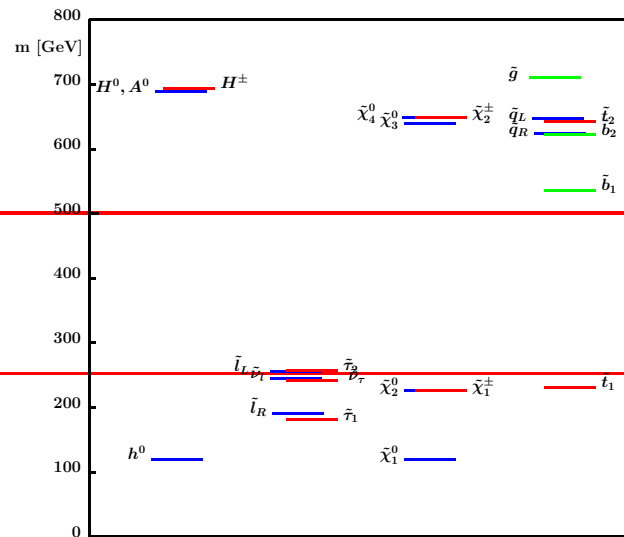


$$m_{\tilde{t}_1} = 396 \text{ GeV}$$

$$\tilde{t}_1 \rightarrow b\chi_1^+, \dots$$

$$\Gamma_{\tilde{t}_1} = 1.92 \text{ GeV}$$

SPS 5

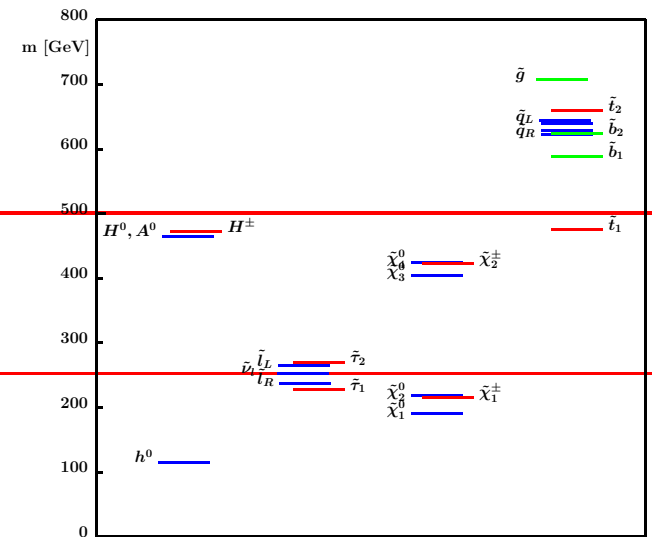


$$m_{\tilde{t}_1} = 240 \text{ GeV}$$

$$\tilde{t}_1 \rightarrow b\chi_1^+, c\chi_1^0$$

$$\Gamma_{\tilde{t}_1} = 0.04 \text{ GeV}$$

SPS 6



$$m_{\tilde{t}_1} \simeq 490 \text{ GeV}$$

$$\tilde{t}_1 \rightarrow b\chi_1^+, \dots$$

$$\Gamma_{\tilde{t}_1} \simeq 3.2 \text{ GeV}$$

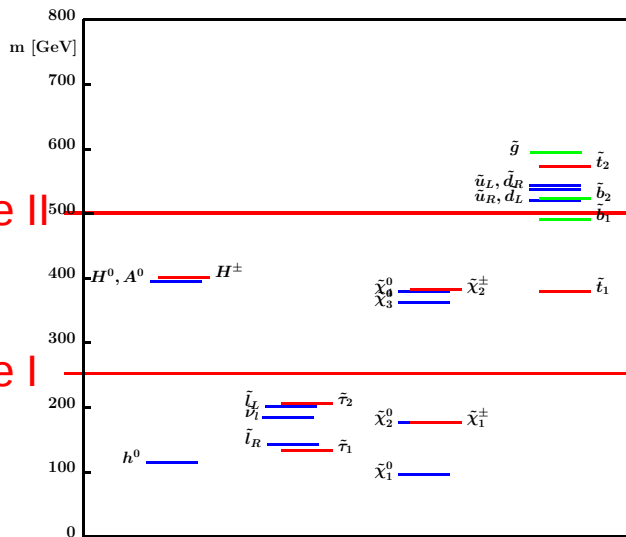


Other Applications

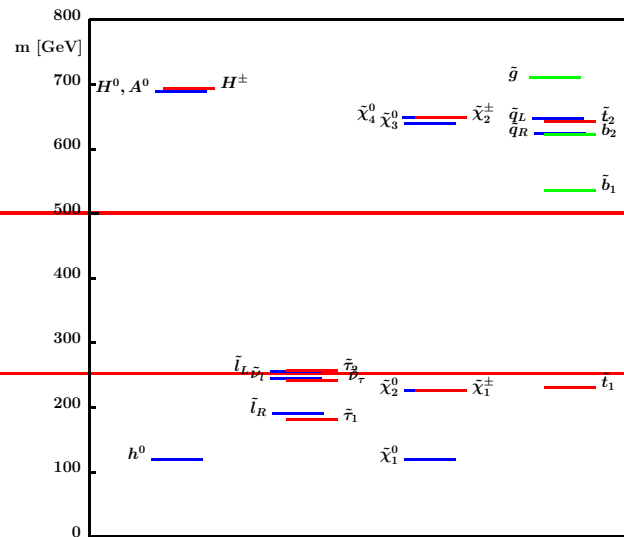
$$e^+e^- \rightarrow \tilde{q}\tilde{q}^*$$

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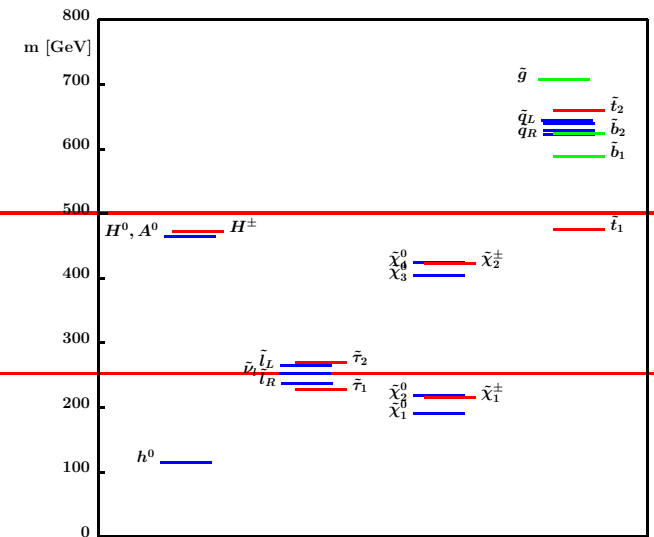
SPS 1a



SPS 5



SPS 6



- low-energy QCD dynamics very similar to $t\bar{t}$ threshold physics
- e.w. weak effects & phenomenology can differ significantly: $m_{\tilde{q}}$, Γ , ...
- no coherent theoretical analysis of QCD & electroweak effects exists

Ruiz-Femenia, Teubner, AHH



Other Applications

$$e^+e^- \rightarrow \tilde{q}\tilde{q} \rightarrow \text{vNRQCD}$$

quarks: $\Lambda = -i \frac{\mathbf{S} \cdot (\mathbf{p}' \times \mathbf{p})}{\mathbf{k}} , \quad T = \boldsymbol{\sigma}_1 \boldsymbol{\sigma}_2 - \frac{3(\mathbf{k}\boldsymbol{\sigma}_1)(\mathbf{k}\boldsymbol{\sigma}_2)}{\mathbf{k}^2}$

$$V = \left[\frac{\mathcal{V}_c(\nu)}{\mathbf{k}^2} + \frac{\mathcal{V}_k(\nu)\pi^2}{m|\mathbf{k}|} + \frac{\mathcal{V}_r(\nu)(\mathbf{p}^2 + \mathbf{p}'^2)}{2m^2\mathbf{k}^2} + \frac{\mathcal{V}_2(\nu)}{m^2} + \frac{\mathcal{V}_s(\nu)}{m^2} \mathbf{S}^2 + \frac{\mathcal{V}_\Lambda(\nu)}{m^2} \Lambda + \frac{\mathcal{V}_t(\nu)}{m^2} T \right]$$

$$\frac{d}{d \ln \nu} \ln C_{3S_1}(\nu) = -\frac{\mathcal{V}_c(\nu)}{16\pi^2} \left(\frac{\mathcal{V}_c(\nu)}{4} + \mathcal{V}_2(\nu) + \mathcal{V}_r(\nu) + \mathbf{S}^2 \mathcal{V}_s(\nu) \right) + \frac{\mathcal{V}_k(\nu)}{2}$$

$$C_{3S_1}(\nu = 1) = 1 - \frac{8}{3} \frac{\alpha_s(m_t)}{\pi}$$



Other Applications

$$e^+e^- \rightarrow \tilde{q}\tilde{q}^*$$

→ scalar vNRQCD

squarks:

- no spin-dependent interactions
- P-wave production
- full NLL QCD running completed

[P. Ruiz-Femenia, AHH]

$$V = \left[\frac{\mathcal{V}_c(\nu)}{\mathbf{k}^2} + \frac{\mathcal{V}_k(\nu)\pi^2}{m|\mathbf{k}|} + \frac{\mathcal{V}_r(\nu)(\mathbf{p}^2 + \mathbf{p}'^2)}{2m^2\mathbf{k}^2} + \frac{\mathcal{V}_2(\nu)}{m^2} \right]$$

$$\frac{d}{d\ln\nu} \ln C_{1P_1}(\nu) = -\frac{\mathcal{V}_c(\nu)}{48\pi^2} \left(\frac{\mathcal{V}_c(\nu)}{4} + \mathcal{V}_2(\nu) + \mathcal{V}_r(\nu) \right) + \frac{\mathcal{V}_k(\nu)}{6}$$

$$C_{1P_1}(\nu=1) = 1 - \frac{8}{3} \frac{\alpha_s(m_t)}{\pi}$$



Conclusion

- Top pair total rate predictions become more and more realistic
 - full set of electroweak corrections (w.i.p)
 - full NNLL QCD correction still unknown $\rightarrow d\sigma/\sigma \sim \pm 6\%$
- reaching goals like $\delta m_t \sim 100$ MeV is no free lunch
still considerable work has to be invested
but prospects are very good
(all remaining problems appear solvable)
- Many interesting applications of threshold physics exist.



Colors

This is blue

This is red

This is brown

This is magenta

This is Dark Green

This is Dark Blue

This is Green

This is Cyan

Test how this color looks

Test how this color looks

Test how this color looks

Test how this color looks

Test how this color looks

