Top Quark Properties in Little Higgs Models

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We study the shifts in the gauge couplings of the top quark induced in the Littlest Higgs model with and without T parity. We find that the ILC will be able to observe the shifts throughout the natural range of model parameters.

1. INTRODUCTION

Identifying the mechanism which breaks electroweak symmetry and generates fermion masses is one of the main physics goals for both the LHC and the ILC. Studies of the top quark have the potential to illuminate this issue; since it is the heaviest of the Standard Model (SM) fermions, the top is expected to couple strongly to the symmetry-breaking sector. Consequently, the structure of that sector can have significant, potentially observable effects on the properties of the top. For example, it is well known that the vector and axial tTZ form factors receive large corrections (of order 5-10%) in certain models of dynamical electroweak symmetry breaking [1]. At future colliders such as the LHC and the ILC, we will be able to pursue a program of precision top physics, similar to the program studying the Z at LEP and SLC. In this manuscript, we study the corrections to the top quark properties in “Little Higgs” models of electroweak symmetry breaking [2], and compare the expected deviations from the SM predictions with expected sensitivities of experiments at the LHC and the ILC.

In the Little Higgs models, electroweak symmetry is driven by the radiative effects from the top sector, including the SM-like top and its heavy counterpart, a TeV-scale “heavy top” T. Probing this structure experimentally is quite difficult. While the LHC should be able to discover the T quark, its potential for studying its couplings is limited [3, 4]. Direct production of the T will likely be beyond the kinematic reach of the ILC. However, we will show below that the corrections to the gauge couplings of the SM top, induced by its mixing with the T, will be observable at the ILC throughout the parameter range consistent with naturalness. Measuring these corrections will provide a unique window on the top sector of the Little Higgs.

Many Little Higgs models have been proposed in the literature. We will consider two examples in this study, the “Littlest Higgs” model [5], and its variation incorporating T parity [6].

2. THE LITTLEST HIGGGS

As our first example, consider the SU(5)/SO(5) Littlest Higgs (LH) model [5]. Since the original model turned out to be severely constrained by precision electroweak data [7], we focus on the version with a reduced gauge group, SU(2) × SU(2) × U(1), which is significantly less constrained. We follow the conventions and notation of Ref. [3]. The new TeV-scale states are the gauge bosons W_0^T, W_0^H, a vector-like weak-singlet quark T (the “heavy top”), and a weak-triplet scalar field φ. The model is parametrized by the symmetry breaking scale f (assumed to be of order 1 TeV), the SU(2) mixing angle ψ, the hT coupling constant λ_T, and the triplet vacuum expectation value (vev) v'.

ALCPG0428
It can be shown that $v' \sim v^2/f \ll v$, where $v = 246$ GeV is the SM Higgs vev. In this analysis, we will set $v' = 0$, since the effects of a non-vanishing $v'$ on the observables considered here are numerically small. Instead of $f$, we will use the more physical quantity, the heavy top mass $m_T$, in our plots; these are related by $m_T/f = (\lambda_T^2 + \lambda_T^2)/\lambda_T$, where $\lambda_T \approx 1$ is the SM top Yukawa. Naturalness arguments put an upper bound on this parameter, $m_T \lesssim 2$ TeV.

Corrections to the gauge couplings of the top quark in the LH model arise from two sources: the mixing of the (left-handed) top with the heavy top $T$, and the mixing of the SM gauge bosons $W^\pm, Z^0$ with their heavy counterparts, $W^\pm_H$ and $Z_H$. Using the superscripts “t” and “g” to denote the contributions from these two sources, the corrections to the $t\bar{t}Z$ coupling can be written as

$$
\begin{align*}
\delta g_{L/R}^{Z_t} &= 0, \\
\delta g_{L/R}^{Z_g} &= \frac{v^2}{4f^2} \frac{c_{\psi}^2 s_{\psi}^2}{c_{W}^2 - s_{W}^2} g_{R}, \\
\delta g_{L/R}^{Z_t} &= \frac{\lambda_T^2 v^2 g_A^2}{m_T^2}, \\
\delta g_{L/R}^{Z_g} &= \frac{v^2}{4f^2} \left[ 2g_A^2 c_{\psi} + g_R^2 \frac{c_{\psi}^2 s_{\psi}^2}{c_{W}^2 - s_{W}^2} \right].
\end{align*}
$$

(1)

Here, $g_{L/R}^{Z_t}$ are the SM left- and right-handed $t\bar{t}Z$ couplings, $g_{L/R}^{Z_g} = (g_R^2 + g_Z^2)/2$ and $g_{L/R}^{Z_g} = (g_R^2 - g_Z^2)/2$ are their vector and axial combinations, $c_{W}, s_{W}$ are respectively the cosine and sine of the weak mixing angle, and $s_{\psi} \equiv \sin\psi$, $c_{\psi} \equiv \cos\psi$. The predicted shifts in the $t\bar{t}Z$ axial and vector couplings for $m_T = 0.5, 1.0, 2.0$ TeV, and $\lambda_T = 0.5, 1, 2$, are plotted in Fig. 1 (left panel), along with the experimental sensitivities expected at the LHC [8] and the ILC [1]. The mixing angle $\psi$ is varied between 0 and $\pi/2$. Note that the shifts have a definite sign. While only a rather small part of the parameter space is accessible at the LHC even with 3000 fb$^{-1}$ integrated luminosity, the ILC experiments will be able to easily observe the shifts in most of the parameter space preferred by naturalness considerations.

The corrections to the $t\bar{b}W$ coupling have the form

$$
\begin{align*}
\delta g_{L/R}^{W_t} &= 0, \\
\delta g_{L/R}^{W_g} &= 0, \\
\delta g_{L/R}^{W_t} &= \frac{1}{4} \frac{\lambda_T^2 v^2 g_W^2}{m_T^2}, \\
\delta g_{L/R}^{W_g} &= \frac{v^2}{4f^2} g_W^2 c_{\psi}^2 \left( c_{\psi}^2 - s_{\psi}^2 - \frac{c_{\psi}^2 c_W^2}{c_{W}^2 - s_{W}^2} \right),
\end{align*}
$$

(2)

where $g_W$ is the SM $t\bar{b}W$ coupling. These corrections induce a shift in the top width, $\delta \Gamma_t/\Gamma_t = 2\delta g_{L/R}^{W_t}/g_W$. The induced shift, as a function of the angle $\psi$, is plotted in the right panel of Fig. 1, where the parameter $\lambda_T$ is varied between 0.5 and 2 for $m_T = 0.5, 1.0$ and 2.0 TeV. The accuracy of the top width measurement expected at the ILC [1] will allow to observe this effect in most of the natural parameter space.
3. LITTLEST HIGGS WITH T PARITY

The LH model can be extended to include a discrete symmetry, T parity [6], which greatly reduces the contributions to precision electroweak observables [9]. The main new feature in this model is the absence of the gauge boson mixing, since light and heavy gauge bosons have opposite charges under T parity. The top-heavy top mixing is still present, however. The resulting corrections to the t\(\bar{t}Z\) and \(t\bar{t}W\) vertices are identical to the corresponding shifts in the model without T parity, \(\delta g_{\mu}^{A}\) and \(\delta g_{\mu}^{W}\), given in Eqs. (1) and (2). The shift in the axial \(t\bar{t}Z\) coupling is plotted in the left panel of Fig. 2. (The shift in the vector coupling is identical up to a sign.) The correction to the top width is shown in the right panel of Fig. 2. Again, both effects should be observable at the ILC.

Acknowledgments

The authors wish to thank Shrihari Gopalakrishna, JoAnne Hewett, Michael Peskin and Tom Rizzo for useful discussions. CFB is supported by the US Department of Energy under contract DE-AC02-76SF00515, MP is supported by the NSF grant PHY-0355005, and FP is supported by the University of Wisconsin Research Committee with funds granted by the Wisconsin Alumni Research Foundation.

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