Frequency Scanned Interferometer for ILC Tracker Alignment

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8th International Linear Collider Workshop
SLAC, March 18-22, 2005
ILC - Silicon Detector

- Barrel – 5 layers, inner/outer radii – 20/125 cm, Silicon drift detector or microstrips
  \( \sigma_{r\phi} = 10 \, \mu m, \sigma_{rz} = 20 \, \mu m \)
- Forward – 5 disks, double-sided silicon microstrips
  \( \sigma_{r\phi} = 7 \, \mu m, \sigma_{rz} = 7 \, \mu m \)
- Coverage - \(|\cos(\theta)|=0.99\)
- Boundary between barrel and forward disks - \(|\cos(\theta)|=0.80\)
- Wafer size – 10cm x 10cm
- Wafer thickness – 150 \( \mu m \)

Ref: SLAC-R-570 (2001) hep-ex/0106058
A Possible SiD Tracker Alignment

752 point-to-point distance measurements
Physics Goals and Background

➔ To Carry out R&D toward a direct, quasi real time and remote way of measuring positions of critical tracker detector elements during operation.

➔ The 1-Dimension accuracy of absolute distance is on the order of 1 micron.

• Basic idea: To measure hundreds of absolute point-to-point distances of tracker elements in 3 dimensions by using an array of optical beams split from a central laser. Absolute distances are determined by scanning the laser frequency and counting interference fringes.

• Assumption: Thermal drifts in tracker detector on time scales too short to collect adequate data samples to make precise alignment.

Background – some optical alignment systems

• RASNIK system: used in L3, CHORUS and CDF, will be used in ATLAS and CMS

• Frequency Scanned Interferometer (FSI): will be used in ATLAS SCT [A.F. Fox-Murphy et al., NIM A383, 229(1996)]

• Focusing here on FSI system for ILC tracker detector
Principle of Distance Measurement

- The measured distance can be expressed by

\[ R = \frac{c \Delta N}{2n_g \Delta \nu} + \text{constant end corrections} \]

- Assuming the error of refractive index is small, the measured precision is given by:

\[ (\sigma_R / R)^2 = (\sigma_{\Delta N} / \Delta N)^2 + (\sigma_{\Delta \nu} / \Delta \nu)^2 \]

Example: \( R = 1.0 \text{ m}, \ \Delta \nu = 6.6 \text{ THz}, \ \Delta N \sim 2R \Delta \nu / c = 44000 \)

To obtain \( \sigma_R \approx 1.0 \mu m \), Requirements: \( \sigma_{\Delta N} \sim 0.02, \ \sigma_{\Delta \nu} \sim 3 \text{ MHz} \)
FSI Demonstration System (I)

- **Tunable Laser**: New Focus Velocity 6308, 3-4 mW, 665.1-675.2 nm.
- **Retroreflector**: Edmund, D=1”, angle tolerance: ±3 arc seconds.
- **Photodiode**: Thorlabs PDA55, DC-10MHz, Amplified Si Detector, 5 Gain Settings.
- **Thorlabs Fabry-Perot Interferometer SA200**, high finesse(>200) to determine the relative frequency precisely, Free Spectral Range (FSR) is 1.5 GHz, with peak FWHM of 7.5 MHz.
- **Thermistors and hygrometer** are used to monitor temperature and humidity respectively.
- **PCI Card**: NI-PCI-6110, 5 MS/s/ch, 12-bit simultaneous sampling DAQ.
- **PCI-GPIB Card**: NI-488.2, served as remote controller of laser.
- **Computers**: 1 for DAQ and laser control, 3 for analysis.
FSI Demonstration System (I)

- Photodetector
- Mirror
- Beamsplitters
- Laser
- Fabry-Perot Interferometer
- Retroreflector
Temperature Measurements

Outside of Box

Inside of Box
FSI with Optical Fibers (II)

Diagram showing the setup involving a Tunable Laser, Isolator, Fiber, BS, Fiber Coupler, Stage, BS, Retroreflector, Return Fiber, Femtowatt Photodetector, and Fabry Perot Interferometer.
A key issue for the optical fiber FSI is that the intensity of the return beams received by the optical fiber is very weak. 

\[\text{e.g. the core of the single mode optical fiber has diameter of } ~5 \text{ } \mu\text{m}.\]  

\[\text{Geometrical Efficiency: } \sim 6.25 \times 10^{-10} \text{ for a distance of 0.5 m}\]

A novelty in our design is the use of a gradient index lens (GRIN lens – 0.25 pitch lens with D=1mm, L=2.58mm) to collimate the output beam from the optical fiber. The density of the outgoing beam is increased by a factor of \(\sim 1000\) by using the GRIN lens. This makes it possible to split the laser beam into many beams to serve a set of interferometers simultaneously.
Multiple-Measurement Techniques

- If drift error ($\varepsilon$) occurs during the laser scanning, it will be magnified by a factor of $\Omega (\Omega \equiv v/\Delta v \sim 67$ for full scan of our tunable laser),

$$\text{OPD}_{\text{measured}} = \text{OPD}_{\text{true}} + \Omega \varepsilon$$

$\Rightarrow$ Plastic box and PVC pipes are constructed to reduce thermal drift.

- Assuming a vibration with one frequency:

$$x_{\text{vib}}(t) = a_{\text{vib}} \times \cos(2\pi f_{\text{vib}} t + \phi_{\text{vib}})$$

- Fringe phase at time $t$:

$$\Phi(t) = 2\pi \times [\text{OPD}_{\text{true}} + 2x_{\text{vib}}(t)]/\lambda(t)$$

$$\Delta N = [\Phi(t)−\Phi(t0)]/2\pi = \text{OPD}_{\text{true}} \times \Delta v/c + [2x_{\text{vib}}(t)/\lambda(t)- 2x_{\text{vib}}(t0)/\lambda(t0)]$$

- If we assume $\lambda(t) \sim \lambda(t0) = \lambda$, measured OPD can be written as,

$$\text{OPD}_{\text{meas}} = \text{OPD}_{\text{true}} + \Omega \times [2x_{\text{vib}}(t)− 2x_{\text{vib}}(t0)] \quad (1)$$

$$\text{OPD}_{\text{meas}} = \text{OPD}_{\text{true}} − \Omega \times 4a_{\text{vib}} \sin[\pi f_{\text{vib}}(t-t0)] \times \sin[\pi f_{\text{vib}}(t+t0)+\phi_{\text{vib}}] \quad (2)$$

$\Rightarrow$ Two new multiple-distance measurement techniques are presented to extract vibration and to improve the distance measurement precision based on Eq.1 and Eq.2, respectively.
Two Multiple-Measurement Techniques

- Fix the measurement window size \((t-t_0)\) and shift the window one F-P peak forward each time to make a set of distance measurements. The average value of all measurements is taken to be the final measured distance of the scan.

- If \(t_0\) is fixed, the measurement window size is enlarged one F-P peak for each shift. An oscillation of a set of measured OPD reflects the amplitude and frequency of vibration.
Vibration Measurement

- A PZT transducer was employed to produce controlled vibration of the retroreflector, \( f_{\text{vib}} = 1.01 \pm 0.01 \text{ Hz}, \) \( \text{amp}_{\text{vib}} = 0.14 \pm 0.02 \text{ \textmu m} \)

- Magnification factor \( \Omega \) for each distance measurement depends on the scanned frequency of the laser beam in the measurement window with smaller \( \Omega \) for larger window - plot(a). Since the vibration is magnified by \( \Omega \) for FSI during the scan, the expected reconstructed vibration amplitude is \( \sim 10.0 \text{ \textmu m} \) assuming \( \Omega \sim 70 \) – plot(b).

\[ \Rightarrow \] The extracted vibration – plot(c)

\( f_{\text{vib}} = 1.007 \pm 0.0001 \text{ Hz}, \)
\( \text{amp}_{\text{vib}} = 0.138 \pm 0.0003 \text{ \textmu m} \)
Vibration Measurement

* Controlled vibration source with very low amplitude

\[ f_{\text{vib}} = 1.01 \pm 0.01 \text{ Hz}, \quad \text{amp}_{\text{vib}} = 9.5 \pm 1.5 \text{ nanometers} \]

* Measured vibration

\[ f_{\text{vib}} = 1.025 \pm 0.002 \text{ Hz}, \]
\[ \text{amp}_{\text{vib}} = 9.3 \pm 0.3 \text{ nanometers} \]

→ Measurable range

\[ f_{\text{vib}} = 0.1 \sim 100 \text{ Hz}, \]
\[ \text{amp}_{\text{vib}} = \text{few nm} \sim 0.4 \mu\text{m} \]
Absolute Distance Measurements

- The scanning rate was 0.5 nm/s and the sampling rate was 125 KS/s.
- The measurement residual versus the No. of measurements/scan shown in Fig.,
  (a) for one typical scan,
  (b) for 10 sequential scans.

It can be seen that the distance errors decrease with increasing $N_{\text{meas}}$.

- $N_{\text{meas}}=1$, precision=1.1 µm (RMS)
- $N_{\text{meas}}=1200$, precision=41 nm (RMS)

Multiple-distance measurement technique is well suited for reducing vibration effects and uncertainties from fringe & frequency determination, BUT not good for drift errors such as thermal drift.
Absolute Distance Measurements

Each precision listed is for standard deviation (RMS) of 10 scans.

<table>
<thead>
<tr>
<th>Distance (cm)</th>
<th>Precision ($\mu m$)</th>
<th>Scanning Rate (nm/s)</th>
<th>FSI System (Optical Fiber or Air)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>open box</td>
<td>closed box</td>
<td></td>
</tr>
<tr>
<td>10.385107</td>
<td>1.1</td>
<td>0.019</td>
<td>2.0</td>
</tr>
<tr>
<td>10.385105</td>
<td>1.0</td>
<td>0.035</td>
<td>0.5</td>
</tr>
<tr>
<td>20.555075</td>
<td>-</td>
<td>0.036, 0.032</td>
<td>0.8</td>
</tr>
<tr>
<td>20.555071</td>
<td>-</td>
<td>0.045, 0.028</td>
<td>0.4</td>
</tr>
<tr>
<td>41.025870</td>
<td>4.4</td>
<td>0.056, 0.053</td>
<td>0.4</td>
</tr>
<tr>
<td>44.982897</td>
<td>-</td>
<td>0.041</td>
<td>0.5</td>
</tr>
<tr>
<td>61.405952</td>
<td>-</td>
<td>0.051</td>
<td>0.25</td>
</tr>
<tr>
<td>65.557072</td>
<td>3.9, 4.7</td>
<td>-</td>
<td>0.5</td>
</tr>
<tr>
<td>70.645160</td>
<td>-</td>
<td>0.030, 0.034, 0.047</td>
<td>0.5</td>
</tr>
</tbody>
</table>

* Distance measurement precisions for various setups using the multiple-distance-measurement technique.
Dispersion Effect

- Dispersive elements, beamsplitter, corner cube prism etc. can create significant offset in measured distance for FSI system since the small OPD change caused by dispersion is magnified by a factor of $\Omega$.

- Sellmeier formula for dispersion in crown glass (BK7)
  \[ n^2(\lambda^2)=1+B1*\lambda^2 / (\lambda^2 -C1)+B2*\lambda^2 / (\lambda^2 -C2)+B3*\lambda^2 / (\lambda^2 -C3) \]

  - $B1=1.03961212$, $B2=0.231792344$, $B3=1.01046945$
  - $C1=0.00600069867$, $C2=0.0200179144$, $C3=103.560653$

- Numerical simulation results (thickness of the corner cube prism = 1.86 cm)
  \[ R_1 – R_{true} = 373.876 \text{ um}, R_{2000} – R_{true} = 367.707 \text{ um} \]

  - $R_1 – R_{2000} = 6.2 \pm 0.2 \text{ um}$

- Real data - fitted result
  \[ R_1 – R_{2000} = 6.14 \pm 0.1 \text{ um} \]

Dispersion effects can be avoided by using hollow retroreflector and put the beamsplitter’s anti-reflecting surface facing the optical fiber.
Error Estimations

- Error from uncertainties of fringe and frequency determination, \( \frac{\Delta R}{R} \sim 1.9 \text{ ppm} \); if \( N_{\text{meas}} = 1200 \), \( \frac{\Delta R}{R} \sim 77 \text{ ppb} \)
- Error from vibration. \( \frac{\Delta R}{R} \sim 0.4 \text{ ppm} \); if \( N_{\text{meas}} = 1200 \), \( \frac{\Delta R}{R} \sim 10 \text{ ppb} \)
- Error from thermal drift. Temperature fluctuations are well controlled down to 0.5 mK(RMS) in Lab by plastic box on optical table and PVC pipes shielding the volume of air near the laser beam. An air temperature change of 1 °C will result in a 0.9 ppm change of refractive index at room temperature. The drift will be magnified during scanning. if \( N_{\text{meas}} = 1200 \), \( \frac{\Delta R}{R} \sim 0.9 \text{ ppm/K} \times 0.5\text{mK} \times \Omega(94) \sim 42 \text{ ppb} \).
- Error from air humidity and pressure, \( \frac{\Delta R}{R} \sim 10 \text{ ppb} \).

The total error from the above sources is \( \sim 89 \text{ ppb} \) which agrees well with the measured residual spread of \( \sim 90 \text{ ppb} \) over different days and times of measurement.
Systematic Error Estimations

* The major systematic bias comes from uncertainty of the Free Spectral Range (FSR) of the Fabry Perot interferometer used to determine scanned frequency range precisely, the relative error would be dR/R ~ 50 ppb if the FSR was calibrated by an wavemeter with a precision of 50 ppb. A wavemeter of this precision was not available for the measurement described here.

* The systematic bias from the multiple-distance-measurement technique was also estimated by changing the starting point of the measurement window, the window size and the number of measurements, the uncertainties typically range from 10-30 nanometers (< 50 ppb).

* The systematic bias from uncertainties of temperature, air humidity and barometric pressure scales should have negligible effect.

The total systematic error is ~ 70 ppb.
Comparison of FSI performances

- National Institute of Standards and Technology (NIST):
  Air transport FSI, Distance: 30 cm – 5 m,
  Precision: ~ 250 nm by averaging measurements of 80 independent scans.
  [J.A. Stone et.al, Applied Optics, V38. No. 28, 5981(1999)]

- University of Oxford – ATLAS Group
  Optical fiber FSI, Distance: 20 cm – 1.2 m,
  Precision: ~215 nm by using dual-laser technique to reduce drift errors

- University of Michigan – ILC Group
  Optical fiber FSI, Distance: 10 cm – 0.6 m (measurable distance limited
  by bandwidth of our femtowatt photodetector, 30-750 Hz)
  Precision: ~50 nm by using new multiple-distance measurement technique
  under well controlled laboratory conditions.
  Vibration: 0.1-100 Hz, > few nanometers, can be extracted precisely
  using new vibration extraction technique.
  [physics/0409110, Accepted for publication by Applied Optics, 2004]
A dual-laser FSI intended to reduce the drift errors is under study currently. Two lasers are operating simultaneously, but the laser beams are isolated by using two choppers.

**Laser #1:** $D_1 = D_{\text{true}} + \Omega_1 \varepsilon_1$

**Laser #2:** $D_2 = D_{\text{true}} + \Omega_2 \varepsilon_2$

Drift errors: $\varepsilon_1 \approx \varepsilon_2 = \varepsilon$

$D_{\text{true}} = \left( D_2 - \rho D_1 \right) / (1 - \rho)$,

Where $\rho = \Omega_2 / \Omega_1$
Fringes & F-P Peaks for Dual-Laser FSI

Chopper edge effects and low photodiode duty cycle per laser complicate measurement – requires study
Summary and Outlook

- Two FSI demonstration systems, with or without optical fibers, were constructed to make high-precision absolute distance measurements.
- Two new multi-distance-measurement analysis techniques were presented to improve absolute distance measurement and to extract the amplitude and frequency of vibration.
- A high precision of ~50 nm for distances up to 60 cm under laboratory conditions was achieved.
- Major error sources were estimated, and the expected error was in good agreement with spread in data.
- We are investigating dual-laser scanning technique used by Oxford ATLAS group currently.
- Michigan group has extended the frontier of FSI technology, but much work lies ahead.
RASNIK provides alignment monitoring with submicron precision, developed at NIKHEF.