

OUTFLOWS FROM BLACK HOLE ACCRETION

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THE BOTTOM LINE:

**THE EXISTENCE OF OUTFLOWS
FROM ACCRETING BLACK HOLES
IS GENERIC**

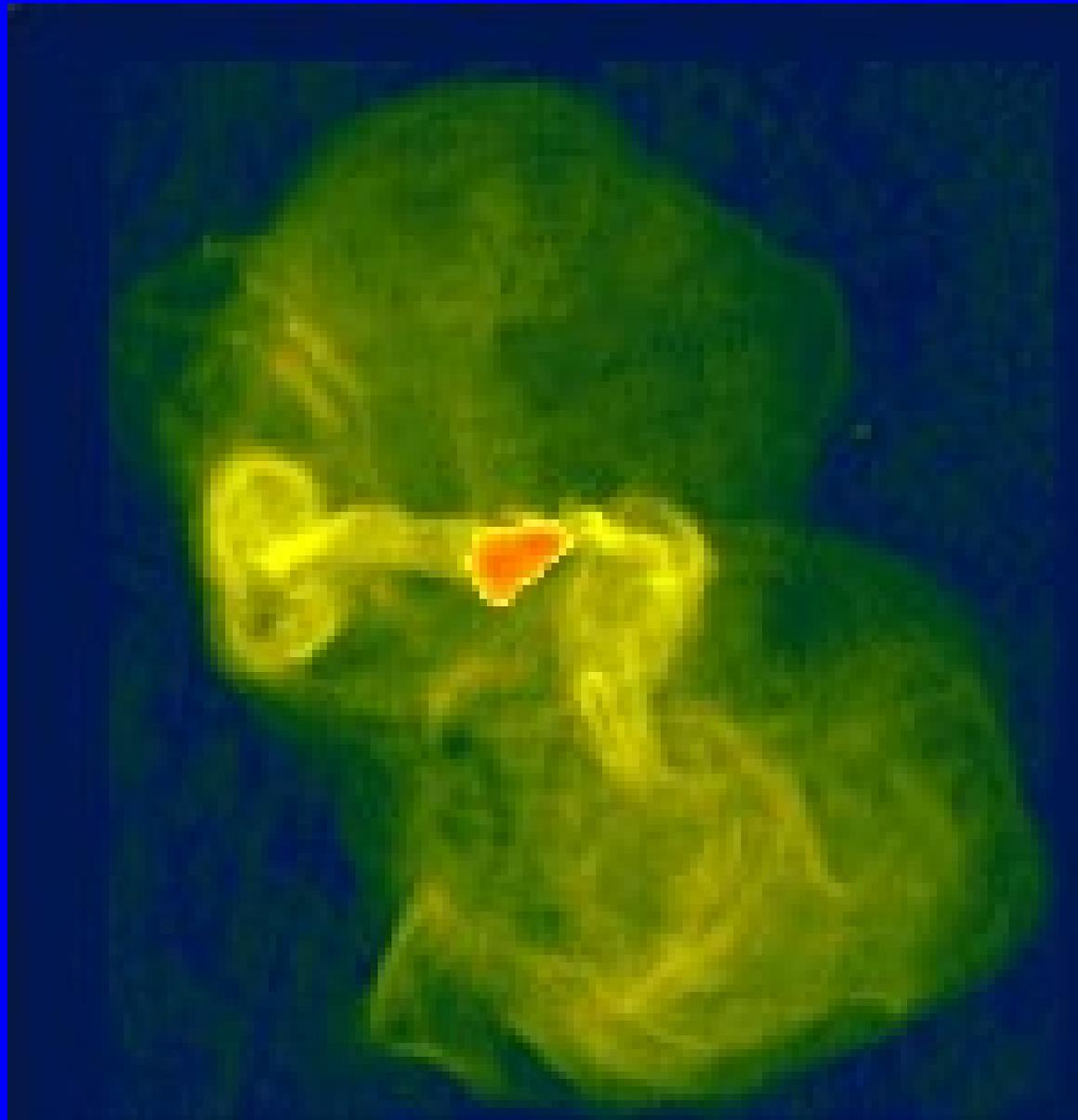
**BUT THE DETAILS ARE DIVERSE
AND SPECIFIC**

RADIO JETS & HALOES: M87

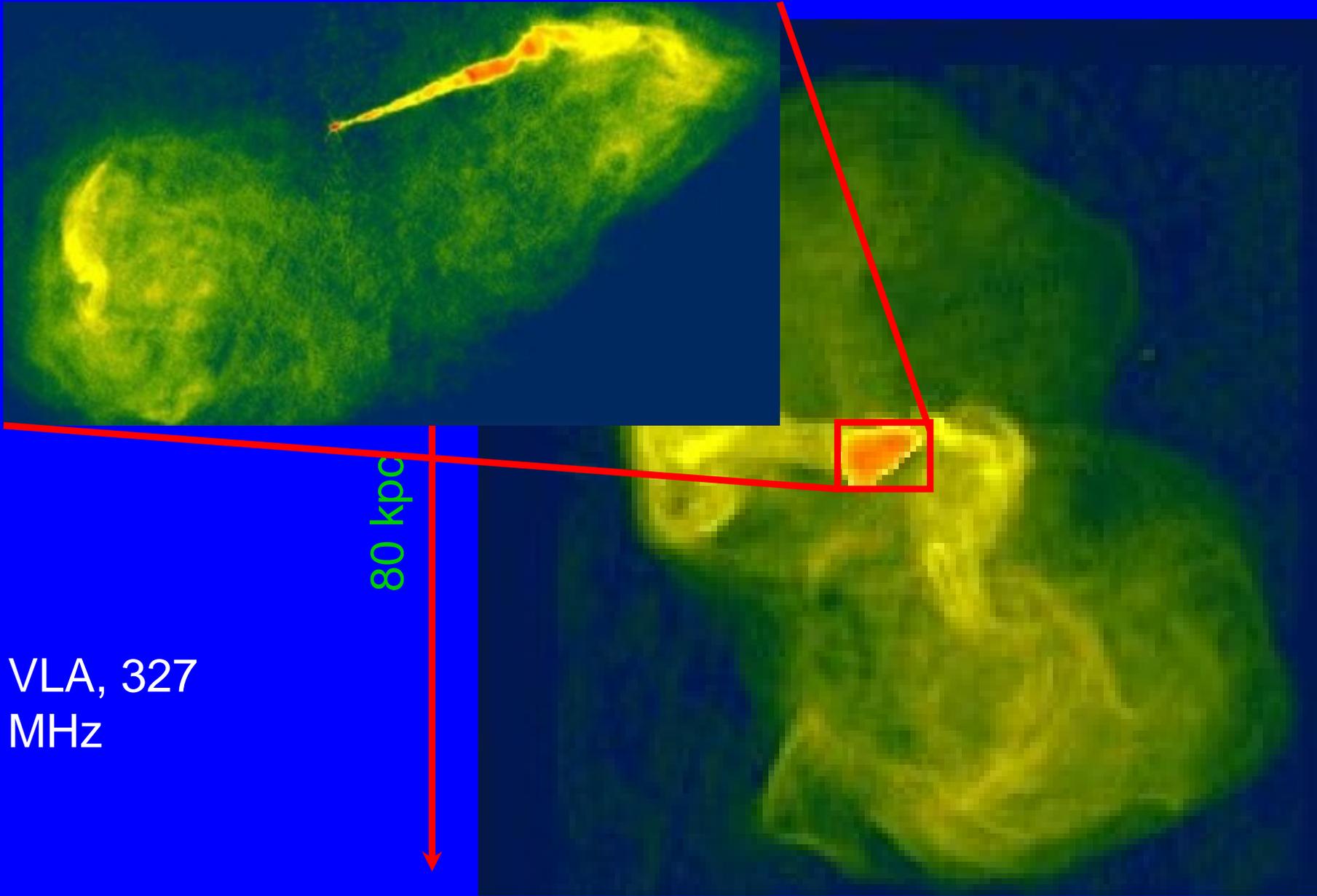
Owen, Eilek & Kassim 1999

80 kpc

VLA, 327
MHz



www.nrao.edu/~fowen/M87.html



VLA, 327
MHz

80 kpc

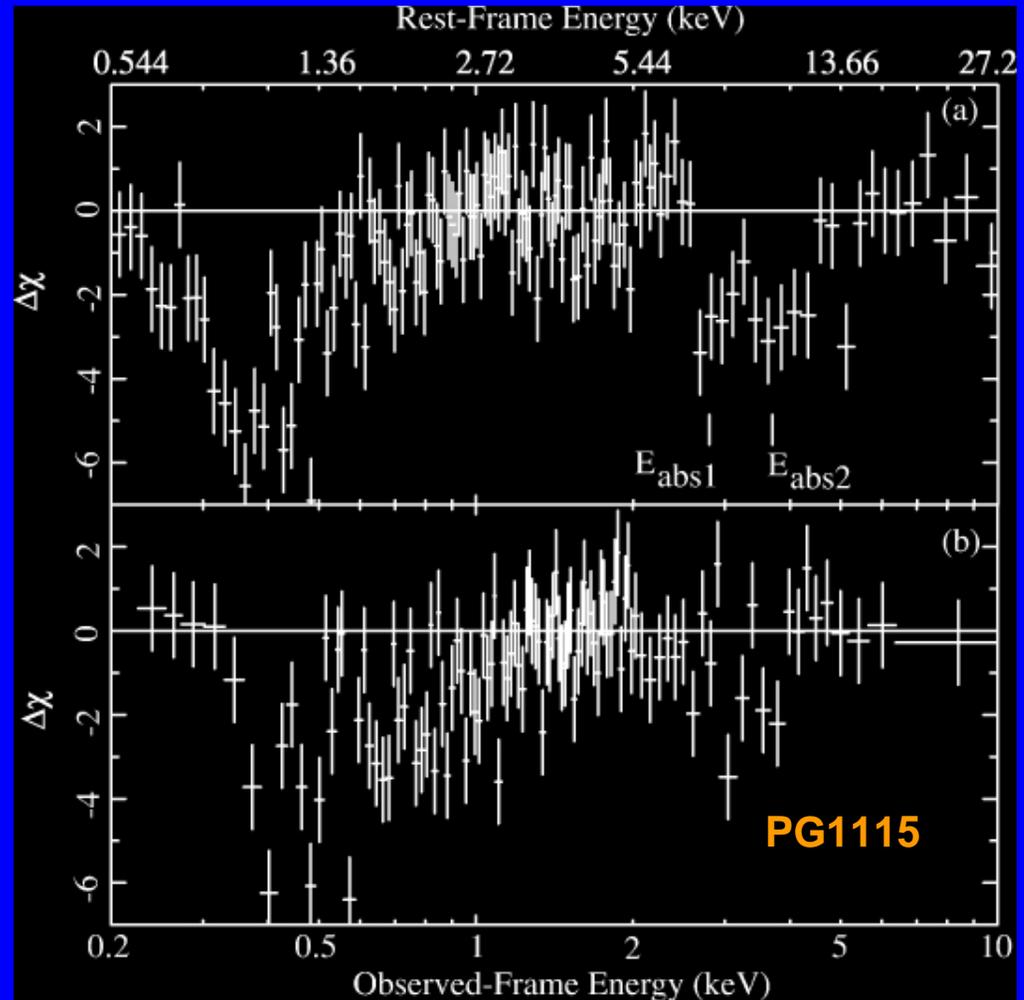
BAL Outflows in X-ray Absorption

$$N_H \sim 10^{23} \text{ cm}^{-2}$$

$$v \sim 0.2 c$$

$$\Omega \sim 4\pi/10$$

$$\frac{L_{KE}}{L_{Bol}} \sim 0.1 \left(\frac{r_{abs}}{10^{17} \text{ cm}} \right)$$



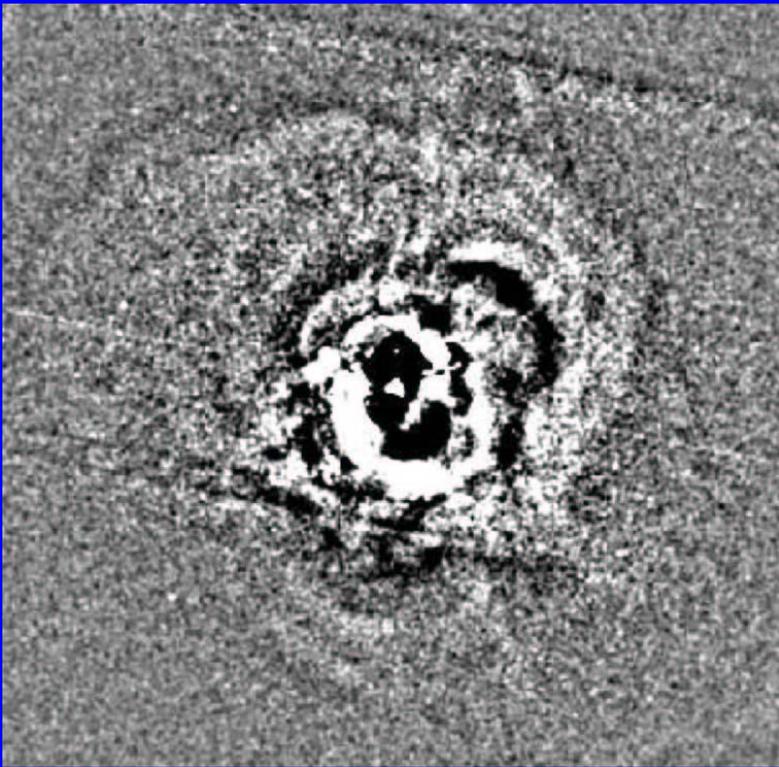
Chartas et al. 2003

Sound Waves in the Intracluster Medium

Perseus cluster

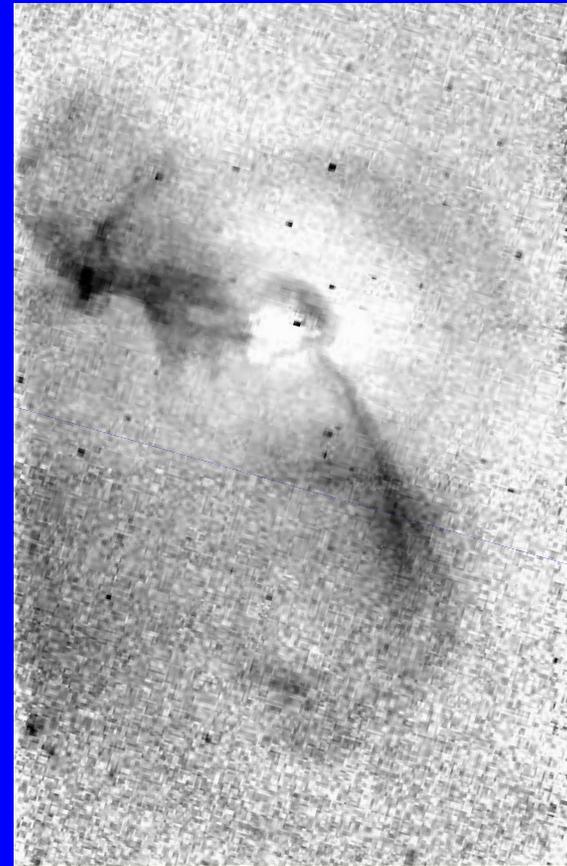
Unsharp masked Chandra image

← 131 kpc →



Fabian et al. 2003

M87 + Virgo cluster



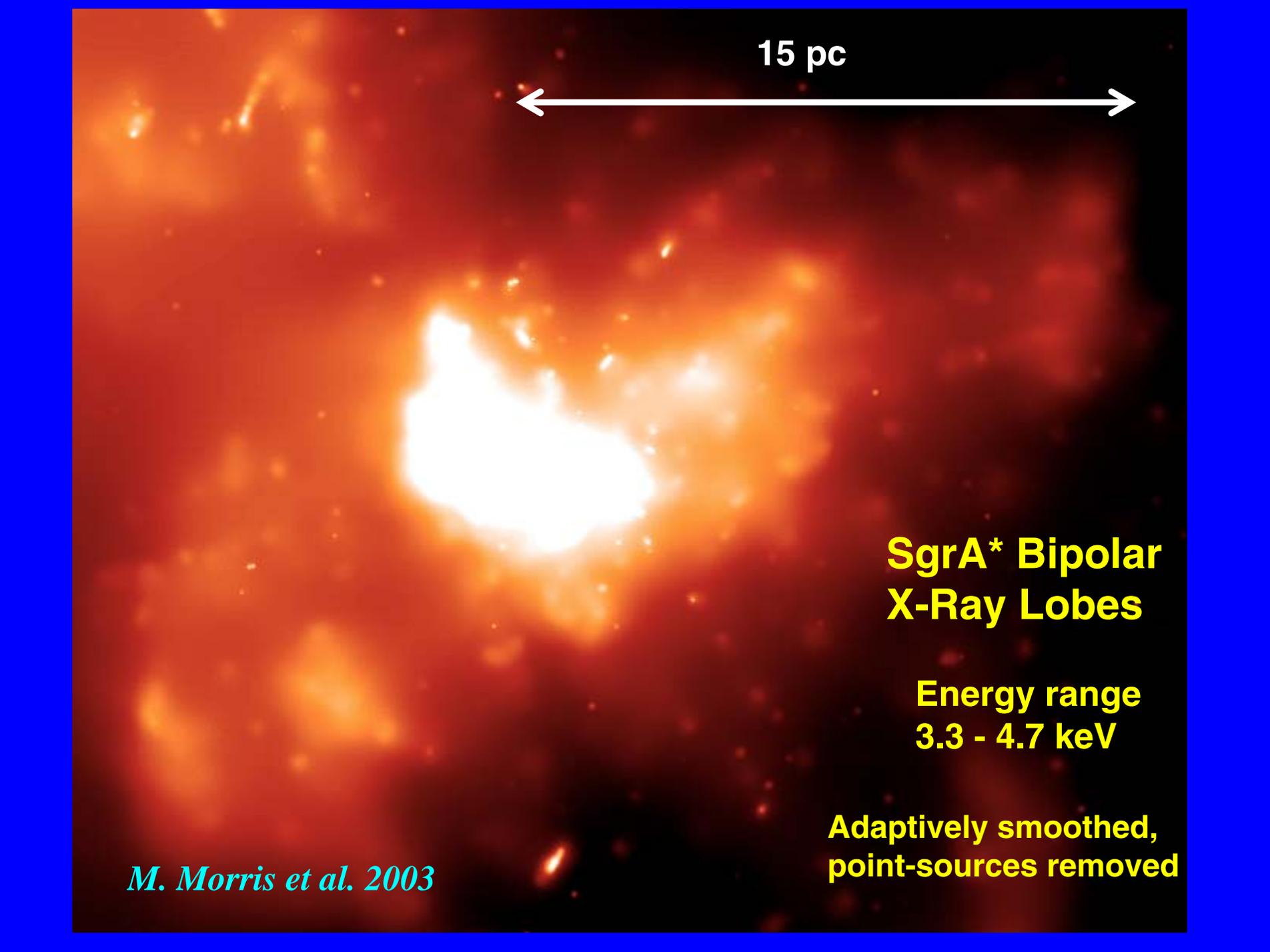
Forman et al. 2004

THE CLUSTER ENTROPY PROBLEM

- Hierarchical clustering $\longrightarrow L_X \propto T^2$
- Observed: $L_X \propto T^3$ \longrightarrow entropy “floor”
- Need ~ 1 keV/baryon EXTRA during cluster assembly
 - mixing/conduction not enough
- Supernova heating inadequate (?)



AGN HEATING?

The image shows a central bright white and yellow region representing the SgrA* source, surrounded by two diffuse, bipolar lobes extending outwards. The lobes are primarily orange and red, with some brighter spots. A white double-headed arrow at the top indicates a scale of 15 parsecs. The background is dark with scattered point sources.

15 pc

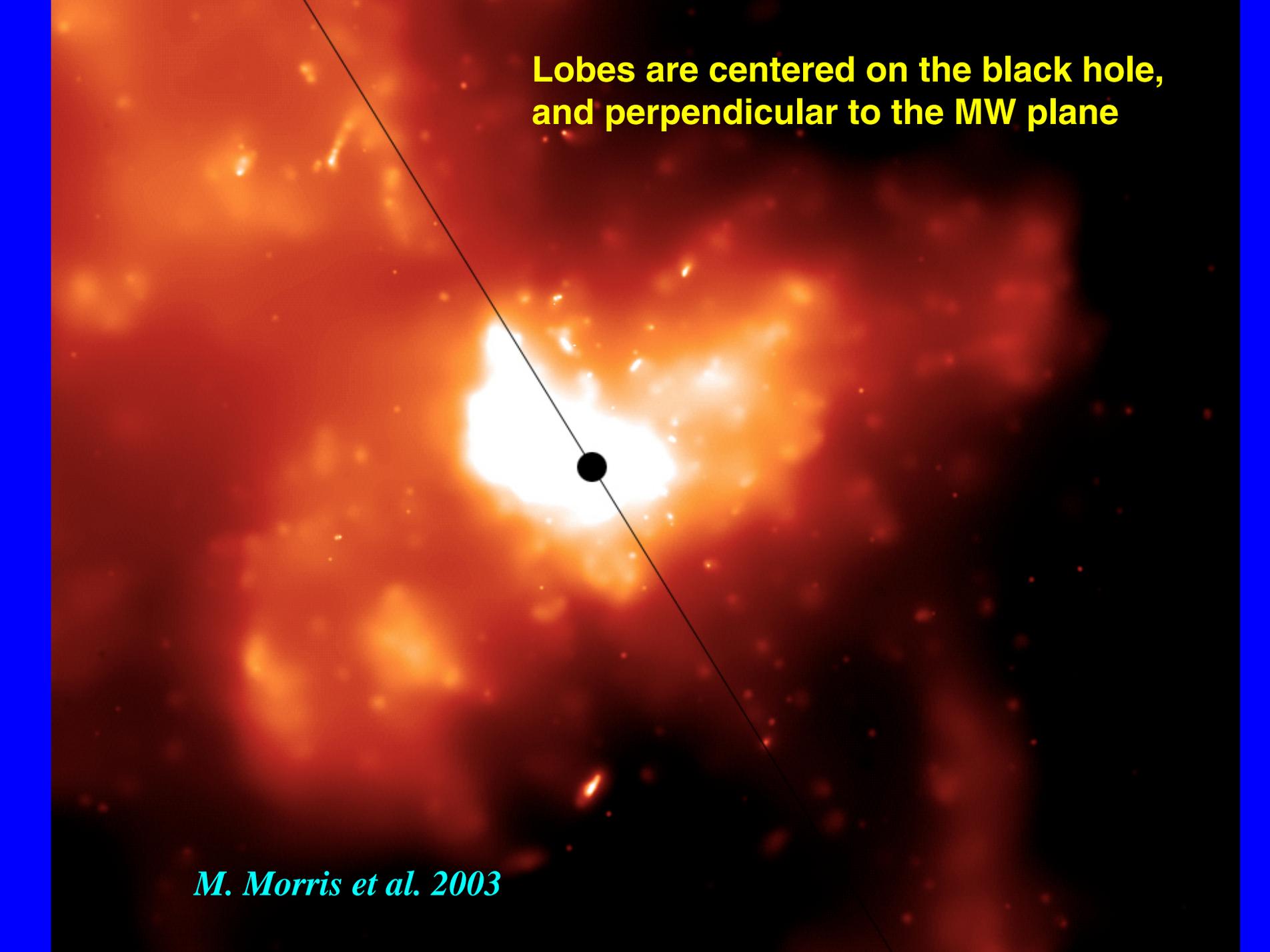
**SgrA* Bipolar
X-Ray Lobes**

**Energy range
3.3 - 4.7 keV**

**Adaptively smoothed,
point-sources removed**

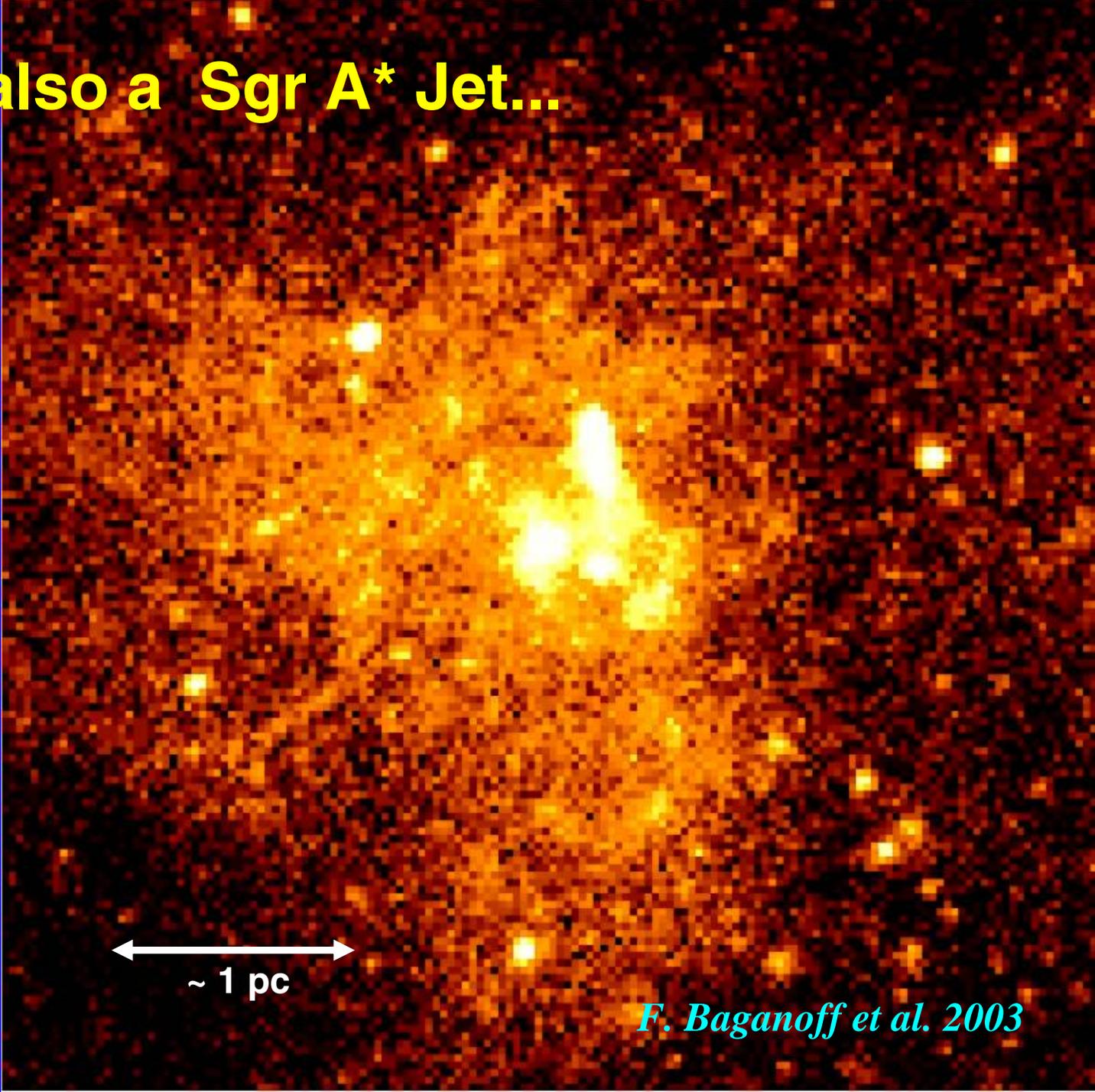
M. Morris et al. 2003

**Lobes are centered on the black hole,
and perpendicular to the MW plane**



M. Morris et al. 2003

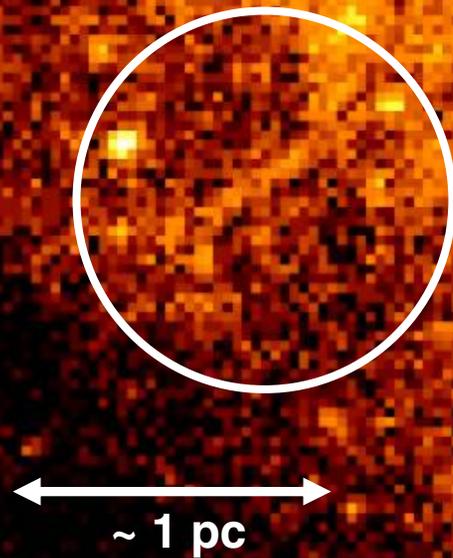
There is also a Sgr A* Jet...

The image shows a dense field of stars in the Sgr A* region. A prominent feature is a bright, elongated structure oriented vertically, which is the jet. The jet is composed of several bright, compact sources. The surrounding field of stars is dense and extends across the entire frame. A scale bar at the bottom left indicates a distance of approximately 1 parsec.

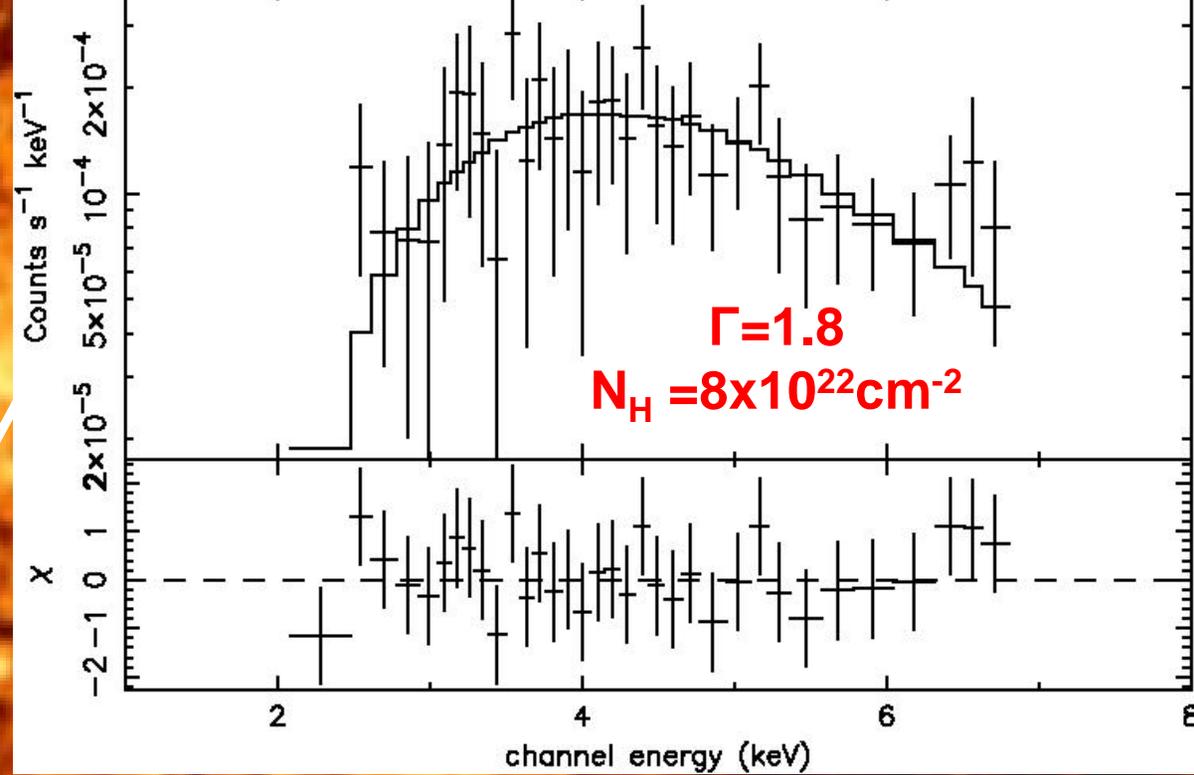
~ 1 pc

F. Baganoff et al. 2003

There is also a Sgr A* Jet...



F. Baganoff et al. 2003



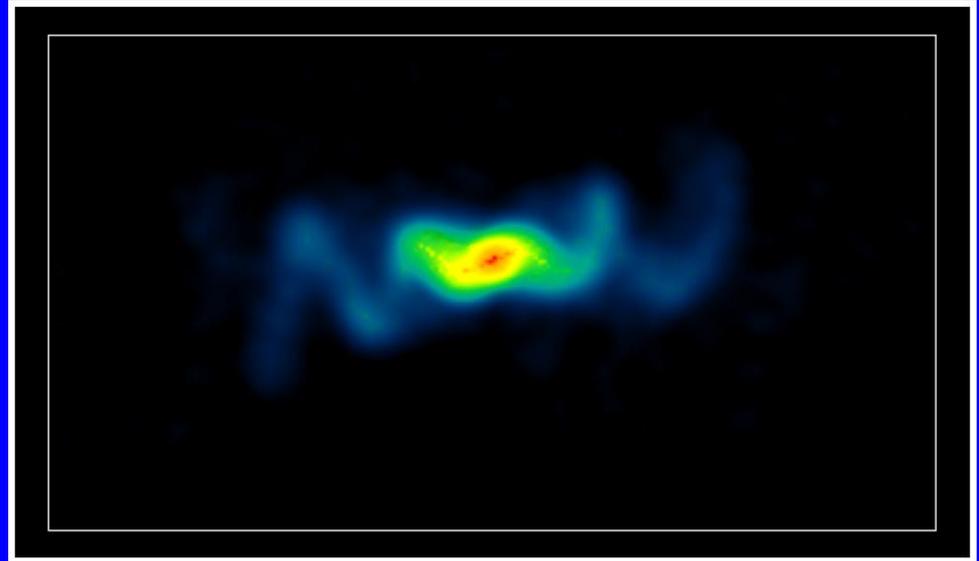
*F. Baganoff et al. 2003,
and in prep. 2005*

...with a hard nonthermal spectrum

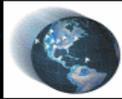
SS433: Super-Eddington accretion + precessing jets

VLA, 5GHz

Blundell & Bowler 2004



SS433
VLBA



Amy Mioduszewski
Michael Rupen
Craig Walker
Greg Taylor

Mioduszewski et al. 2004

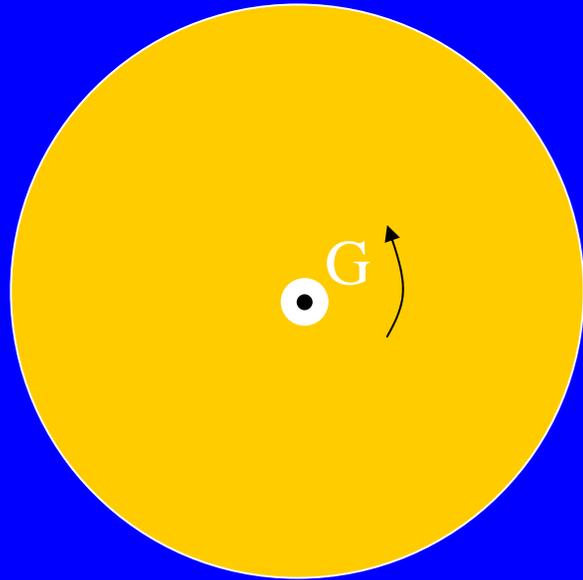
**WHY ARE BLACK
HOLES**

SUCH FUSSY EATERS?

[CLUE: it's not just accreting black holes.... (cf. WDs, NSs, protostars)]

The culprit: Excess angular momentum, which must be transported outward... by internal torques, or winds

TORQUE TRANSPORTS ENERGY



Angular Momentum Flux:

Torque $G \sim \dot{M}\ell$ outward

Energy Flux:

$G\Omega$ outward

IN A THIN ACCRETION DISK:

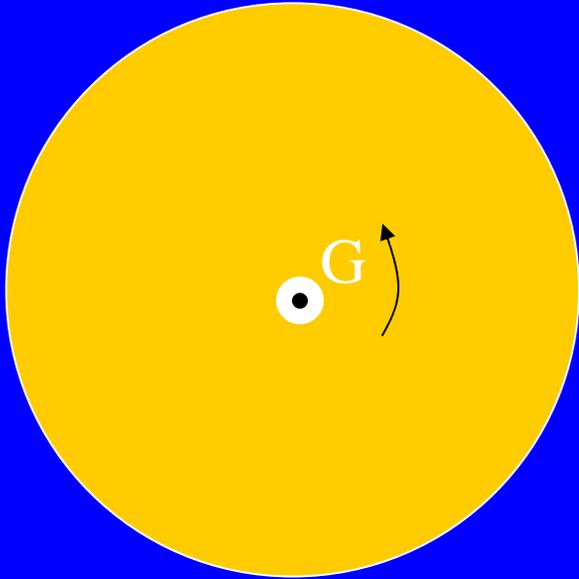
Local rate of energy release:

$$\dot{M} \frac{GM}{2R}$$

Local rate of dissipation:

$$3 \times \left(\dot{M} \frac{GM}{2R} \right)$$

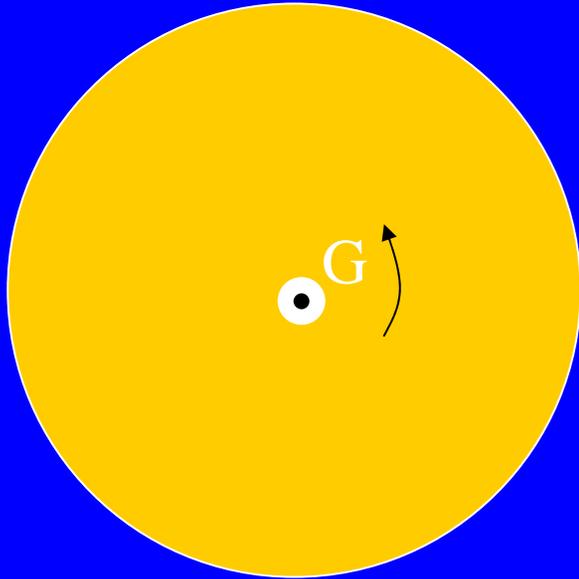
2/3 of energy dissipated at R transported from $<R$ by viscous torque



IN A RADIATIVELINEFFICIENT ACCRETION FLOW (RIAF):

Energy Transport:

$$G\Omega \sim \dot{M} B$$



$$B = \frac{v^2}{2} + \Phi + h > 0$$



Bernoulli Function

Energy transport from small R by torque unbinds gas at large R unless radiative efficiency $> 2/3$

RIAFs are EXPLOSIVE

1 g of gas accreting at $r \sim m$
can liberate 1 kg of gas at $r \sim 1000 m$

- Torque a “conveyor belt” for liberated energy
- Flow **must** lose energy OR limit accretion
 - Mass loss or circulation
 - Small fraction of supplied mass reaches BH

ADIOS = ADIABATIC INFLOW-OUTFLOW
SOLUTION (Blandford & Begelman 99)

Adiabatic accretion occurs in 3 limits:

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 - gas highly opaque, radiates but photons can't escape

$$\dot{M} > \dot{M}_{\text{Eddington}}$$

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- **Low accretion rate**

- If:**
- Dissipated energy goes mainly into protons
 - Protons-electron thermal coupling weak

- gas is tenuous, falls into BH before radiating

$$\dot{M} < \alpha^2 \dot{M}_{\text{Eddington}}; \quad \alpha < 1$$

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- **Everything in-between?**

- Coronae are common

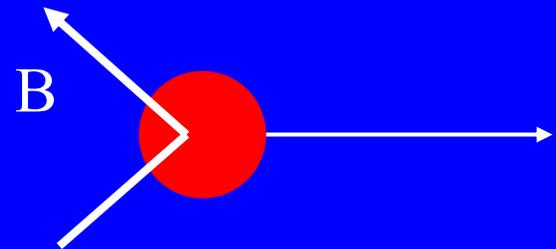
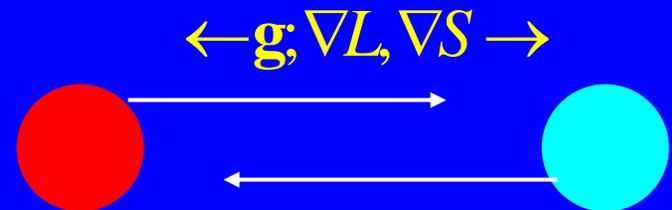
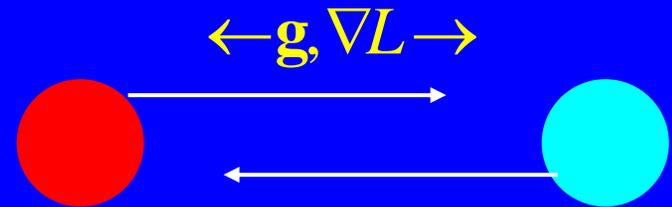
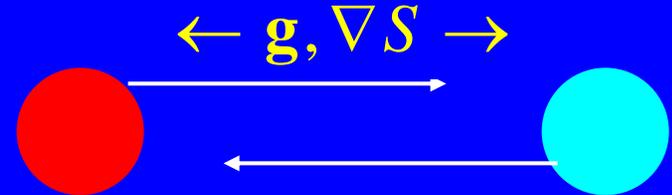
Energy argument is *generic*...

**PHYSICALLY, WHAT DRIVES THE
OUTFLOW or CIRCULATION?**

- **Candidate mechanisms:**
 - magnetic torques, flares
 - radiation pressure
 - radiative heating
- **Convection = “minimal” mechanism**
 - all else being equal, adiabatic flows **must** be convectively unstable

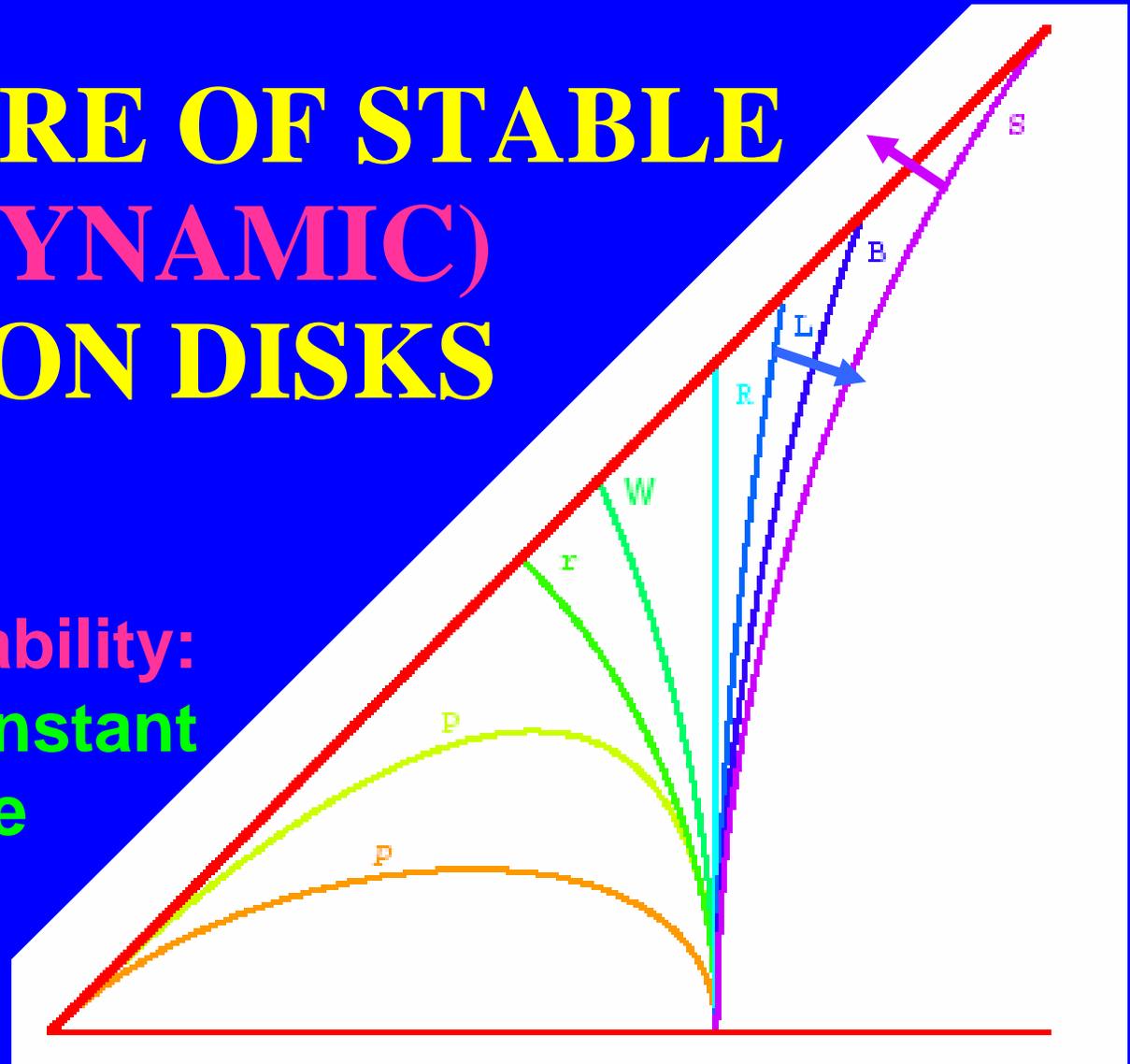
CONVECTION IN ACCRETION DISKS

- **Entropy Gradient**
 - Schwarzschild Criterion
- **Ang. Mom. Gradient**
 - Rayleigh Criterion
- **Both Gradients**
 - Høiland Criterion
- **MHD**
 - MRI, magnetic buoyancy



STRUCTURE OF STABLE (HYDRODYNAMIC) ACCRETION DISKS

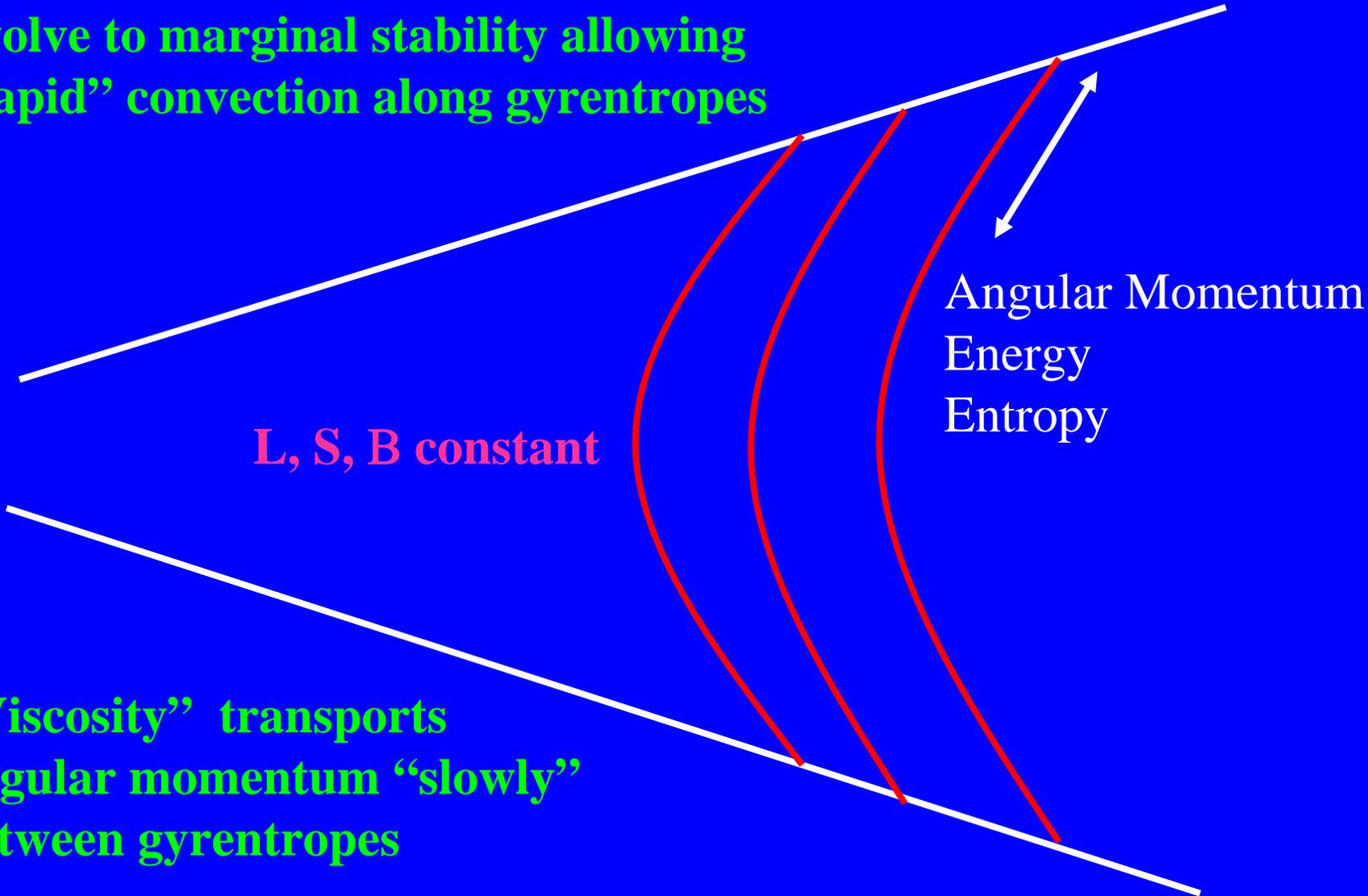
At marginal stability:
Surfaces of constant
 S , L , B coincide



Flow is “GYRENTROPIC”

“GYRENTROPIC HYPOTHESIS”

Evolve to marginal stability allowing
“rapid” convection along gyrentropes



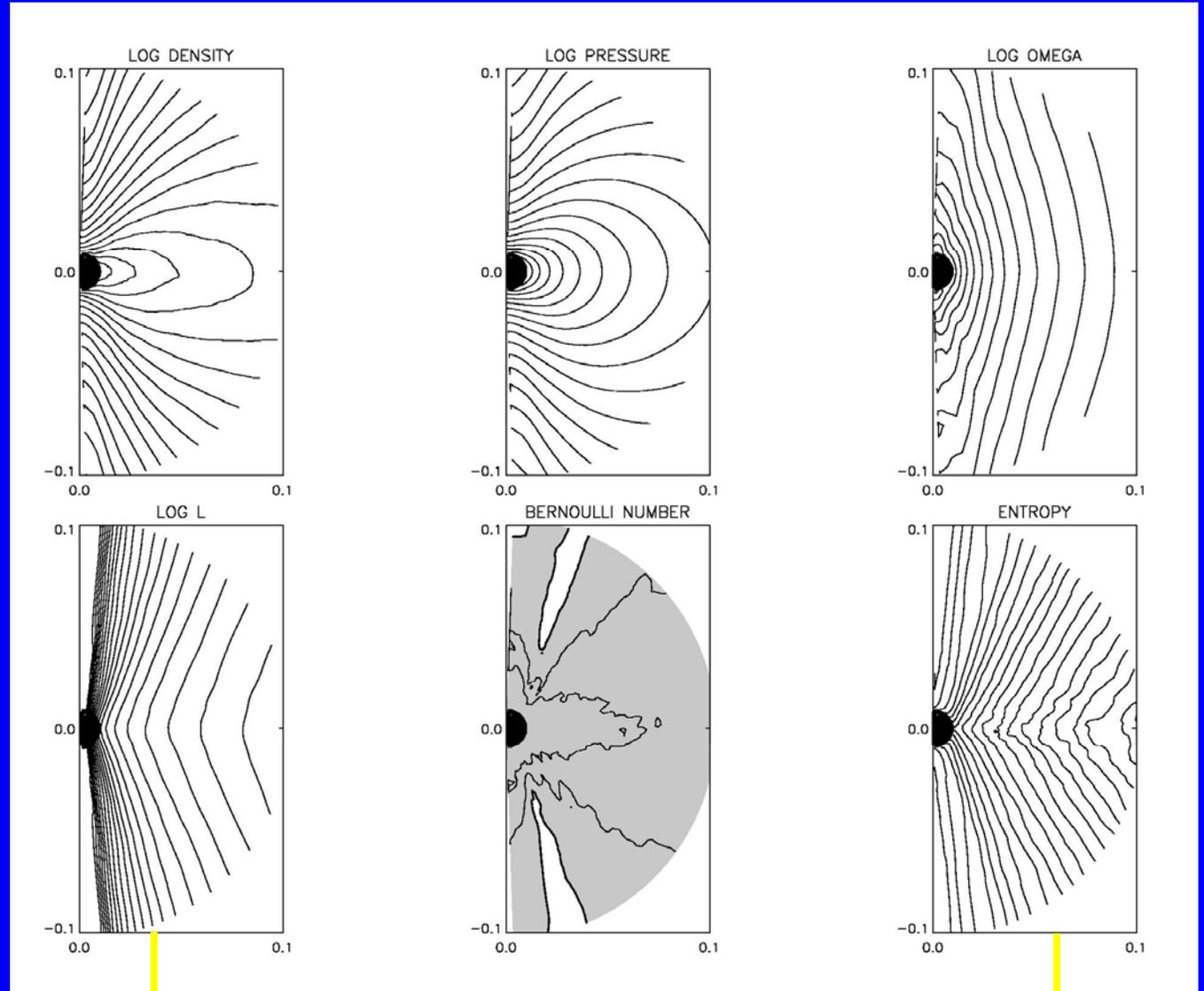
L, S, B constant

Angular Momentum
Energy
Entropy

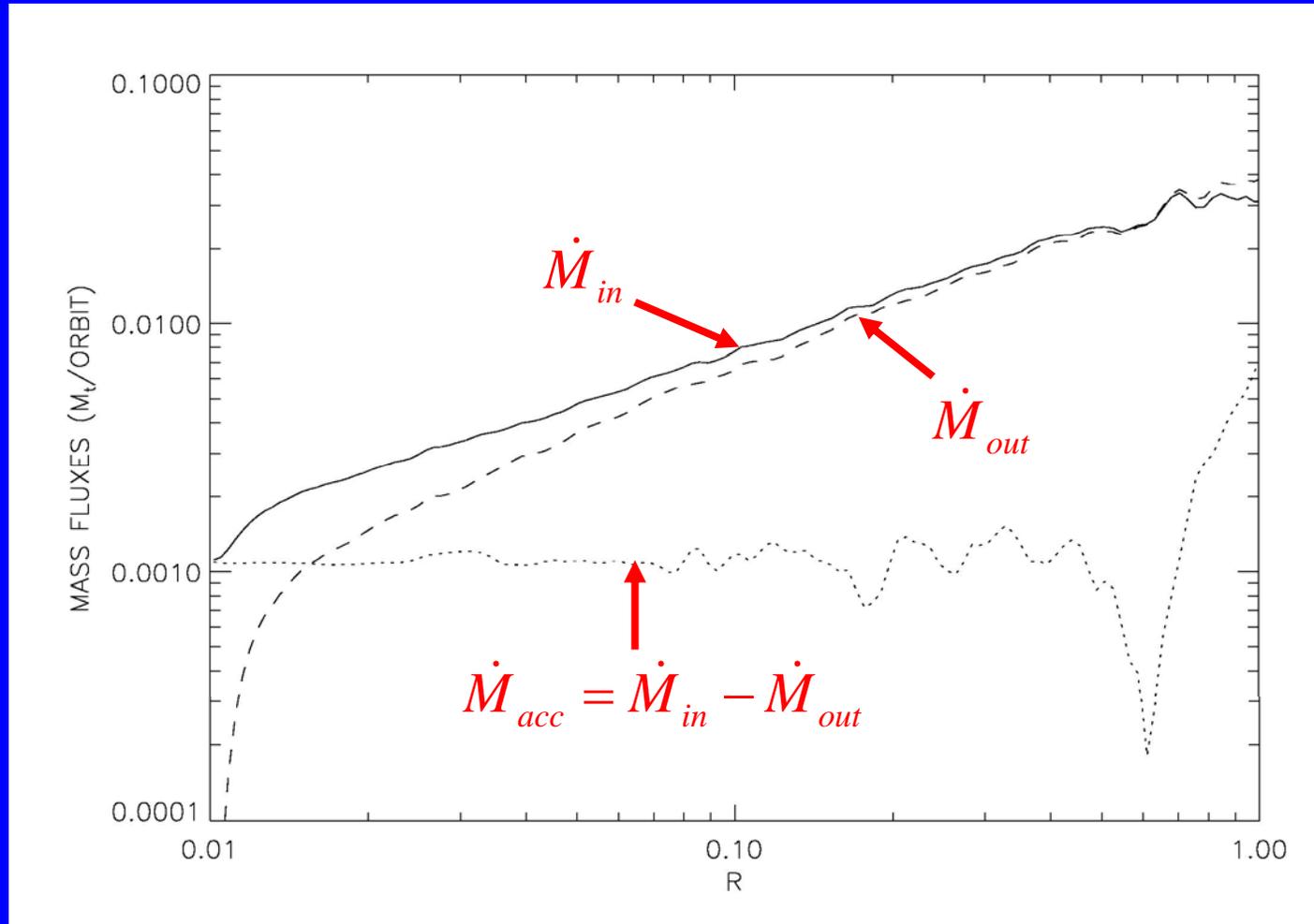
“Viscosity” transports
angular momentum “slowly”
between gyrentropes

2-D, adiabatic α -model

(time average)



2-D, adiabatic α -model: ADIOS



SELF-SIMILAR DISK-WIND MODEL

Disk: Viscous flow
with $B < 0$

Entropy increases at
disk-wind interface

Wind: Inviscid outflow
with $B < 0$

Jet: Evacuated cone

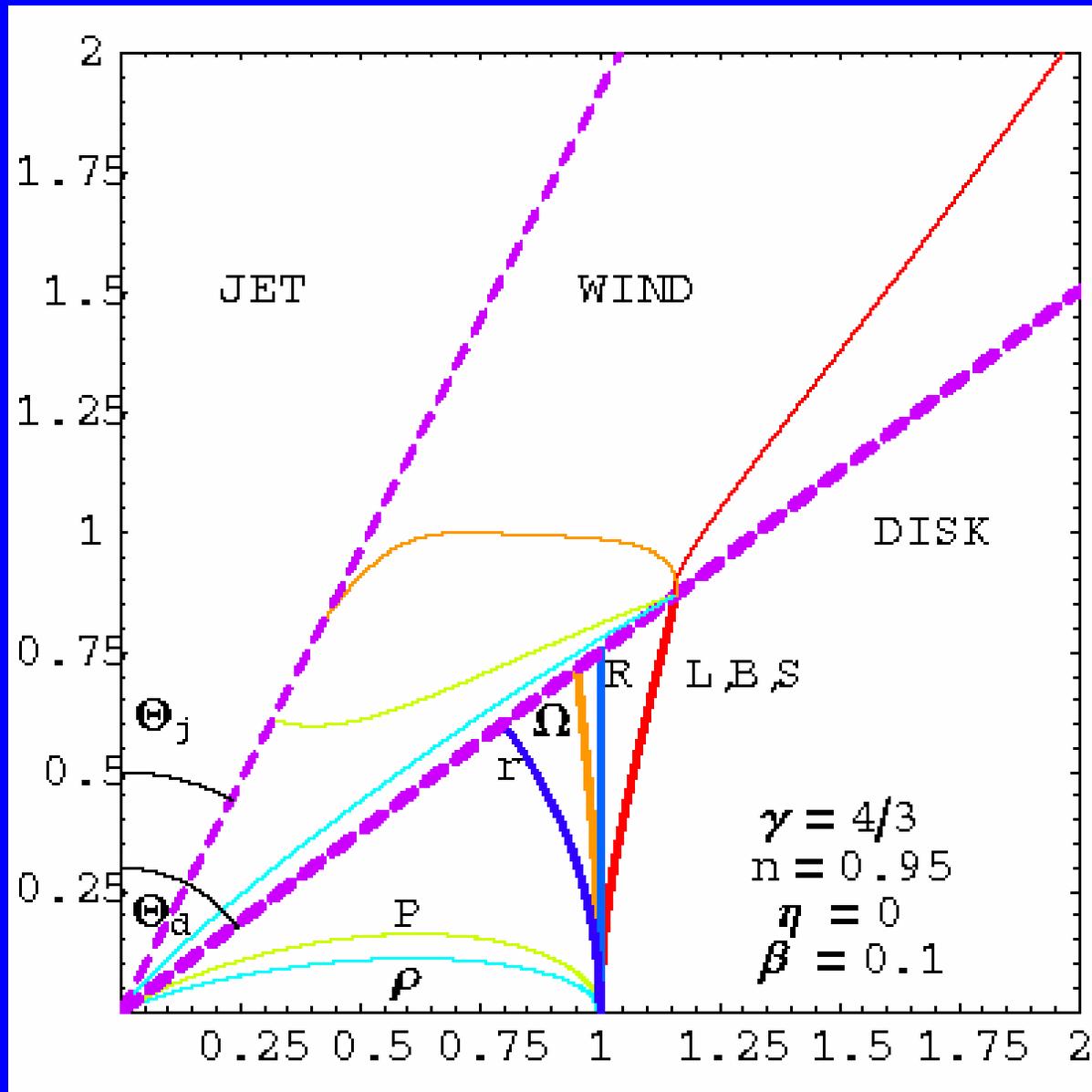
$$dM/dt \sim r^{0.95}$$

$$P \sim r^{-1.55}$$

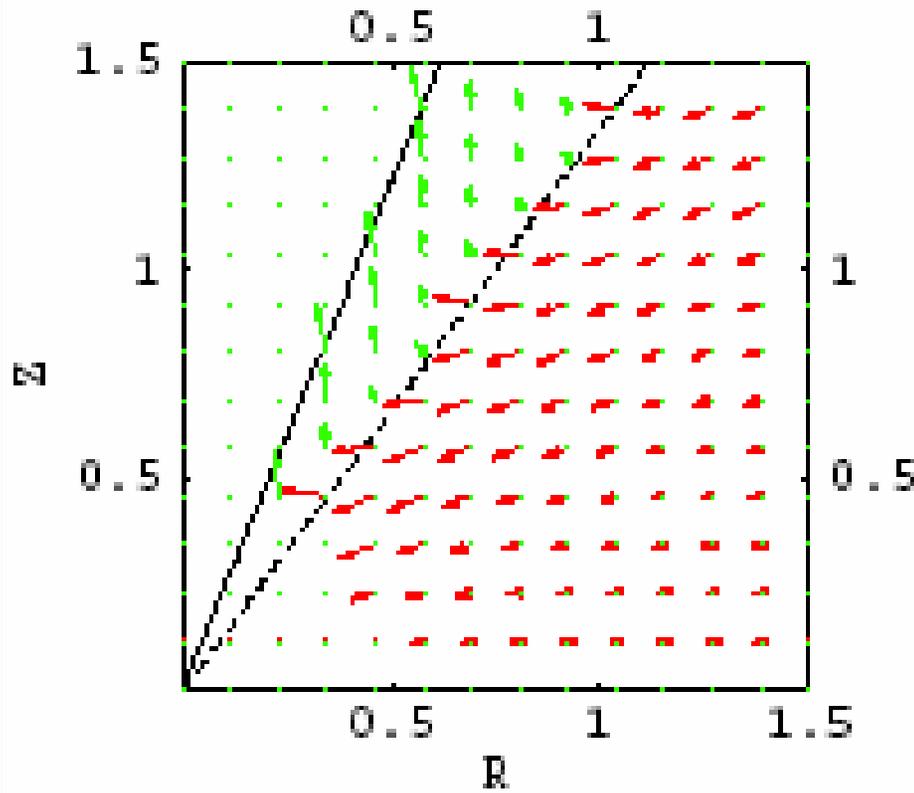
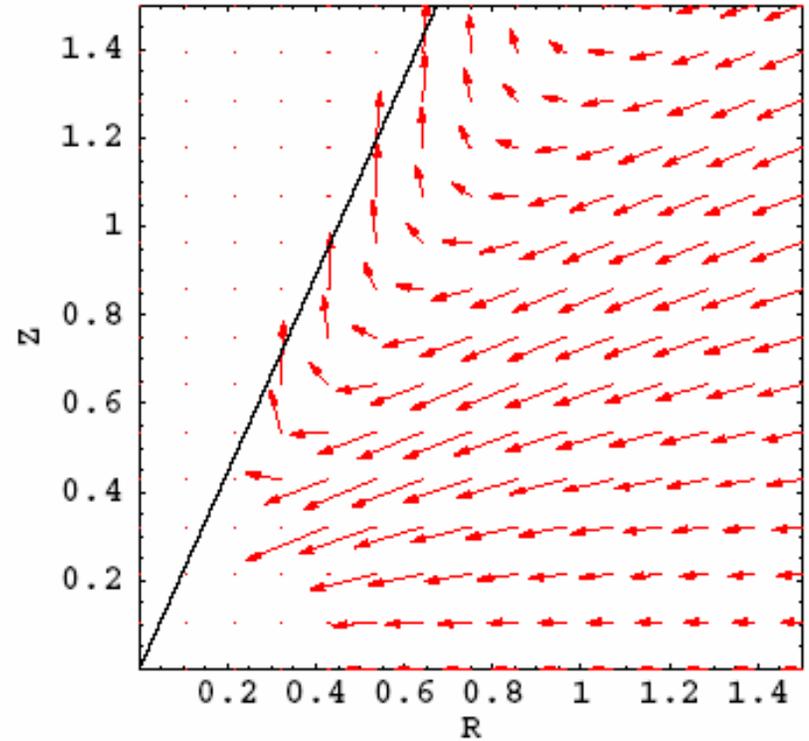
$$\rho \sim r^{-0.55}$$

$$V \sim r^{-0.5}$$

Blandford & Begelman 2004



Convection drives
meridional circulation...

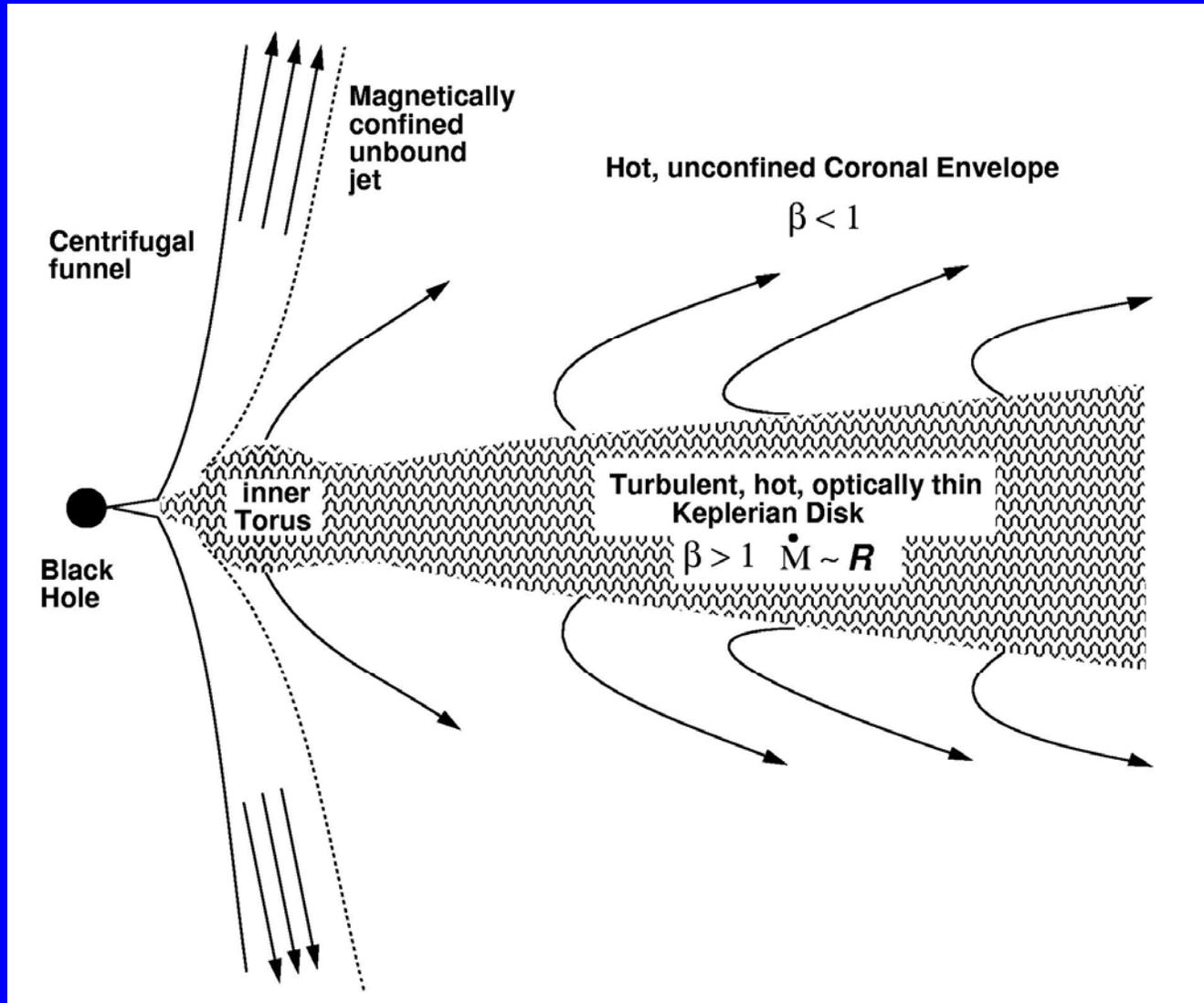


...which matches onto
thermal wind

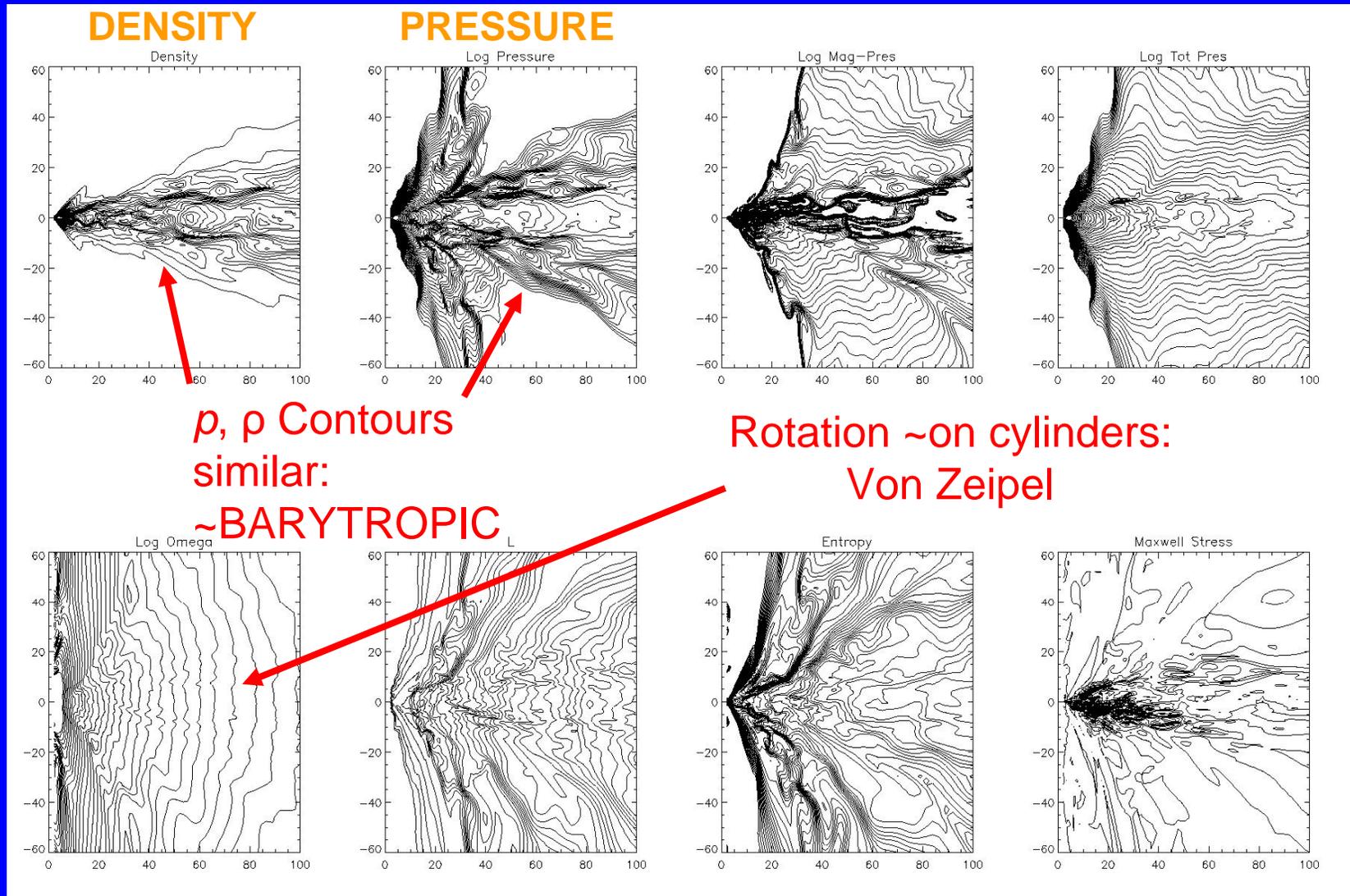


Blandford & Begelman 2004

ADIOS behavior also in MHD...



...but there are differences



ρ , ρ Contours
similar:
~BARYTROPIC

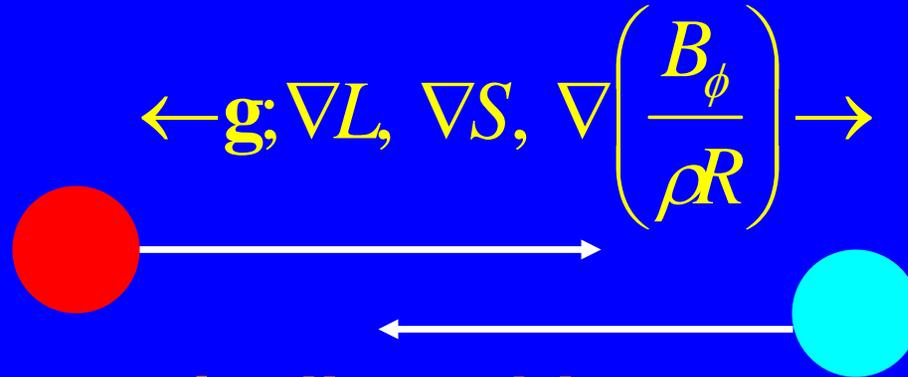
Rotation ~on cylinders:
Von Zeipel

What's Different About MHD?

- **No marginal stability:** always unstable to MRI
So why is there any systematic structure?
(“looks like” marginal stability)
- **Clue 1:** MRI most effective on small scales, convection works best on large scales.
- **Clue 2:** MRI “winds up” B_ϕ more than B_{pol} , buoyancy of B_ϕ can affect disk structure
- **Result?** MRI dominates on small scales, but magnetic buoyancy governs overall disk structure...maybe (similar to hydro case, but different details)

Magnetic Høiland Criterion

- 3 quantities to conserve

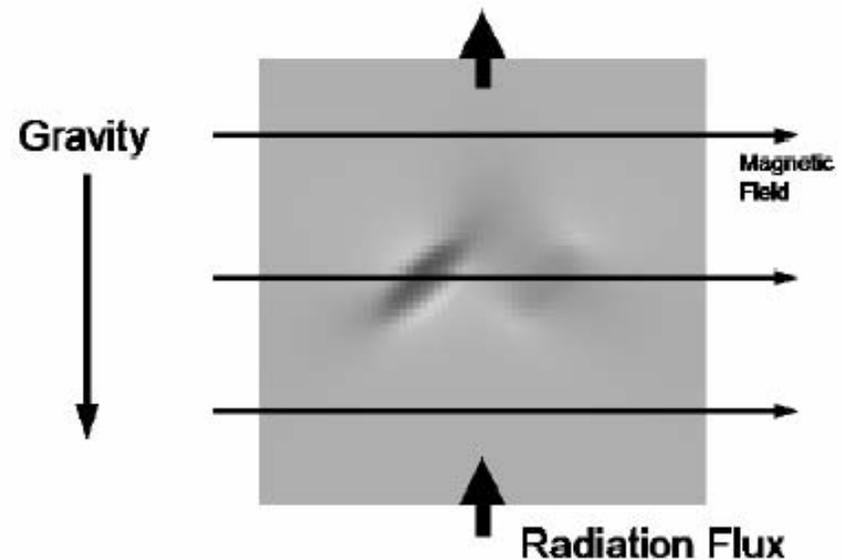
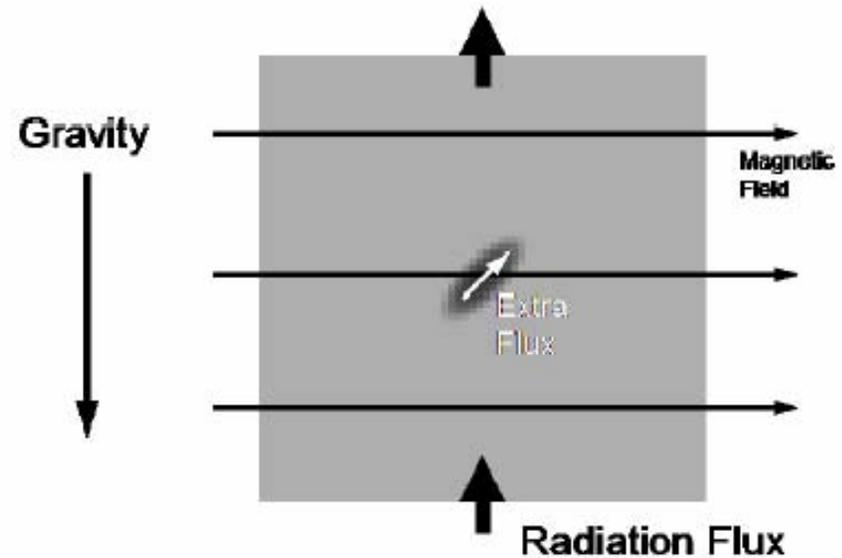


- No unique marginally stable state

– special case: "Campogyrentropic"

- Modest B_ϕ can make flow rotate on \sim cylinders
- Does magnetic convection power the outflows?

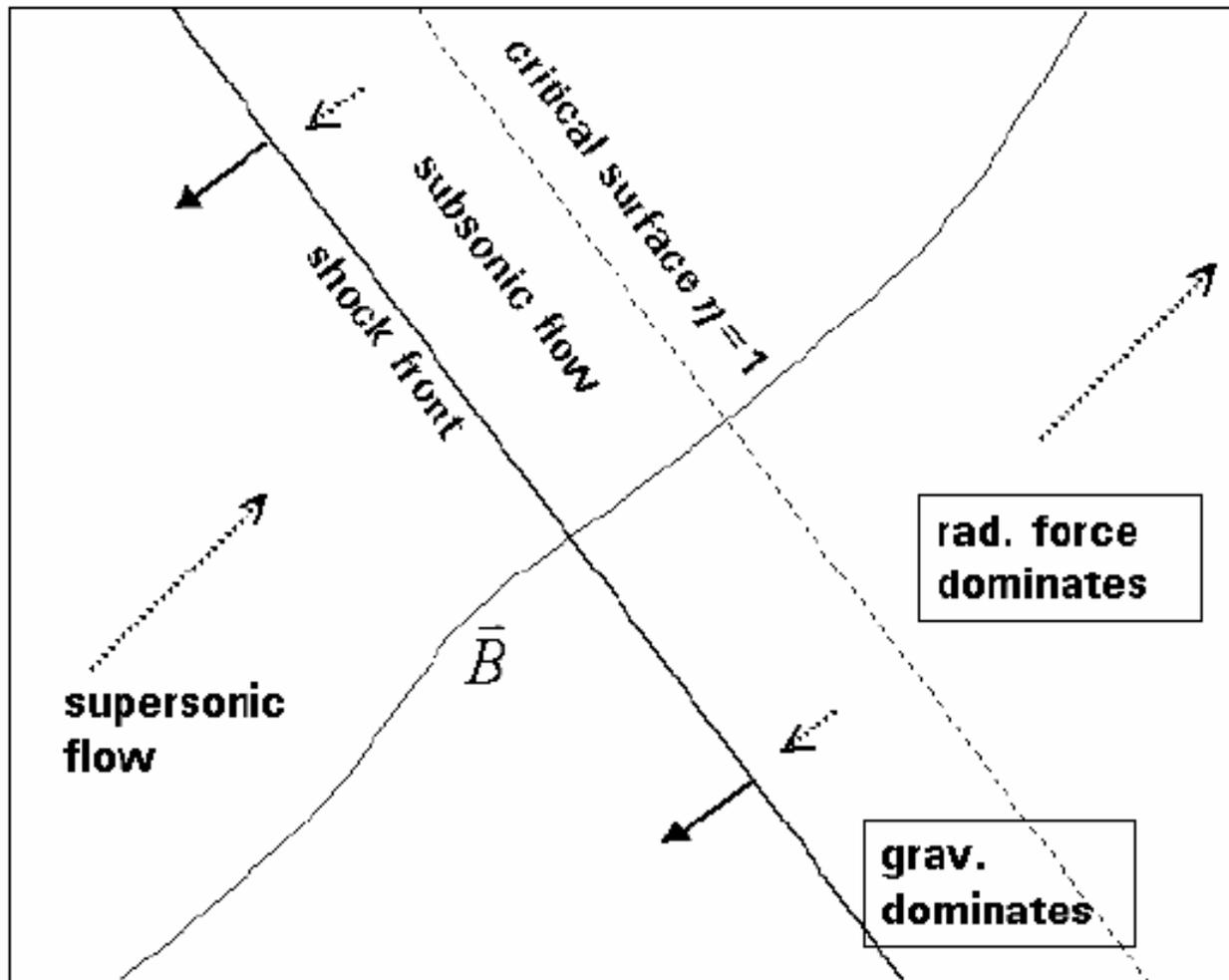
PHOTON BUBBLE INSTABILITY...



N. Turner et al. 2004

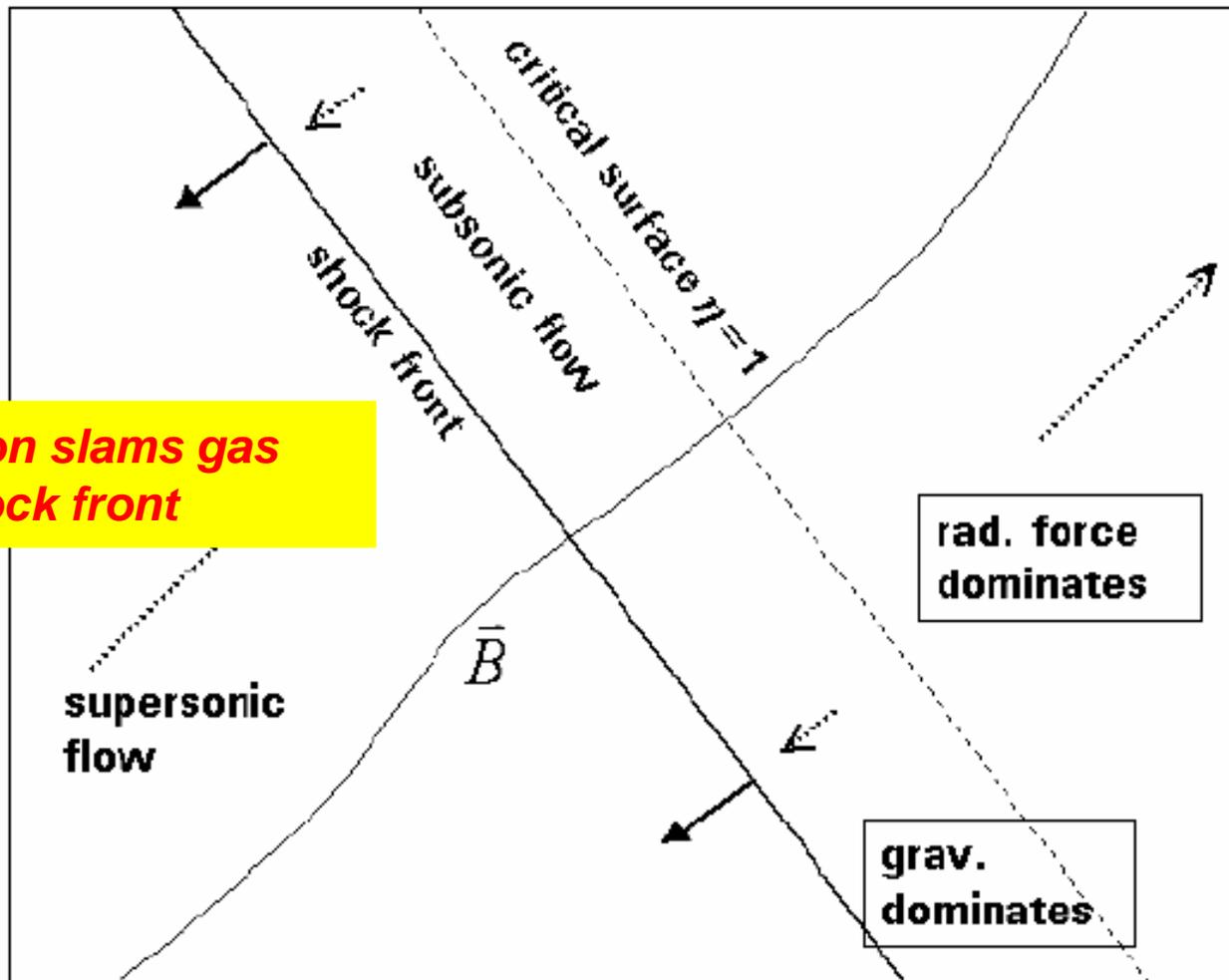
Nonlinear Evolution:

PHOTON BUBBLE SHOCK TRAINS



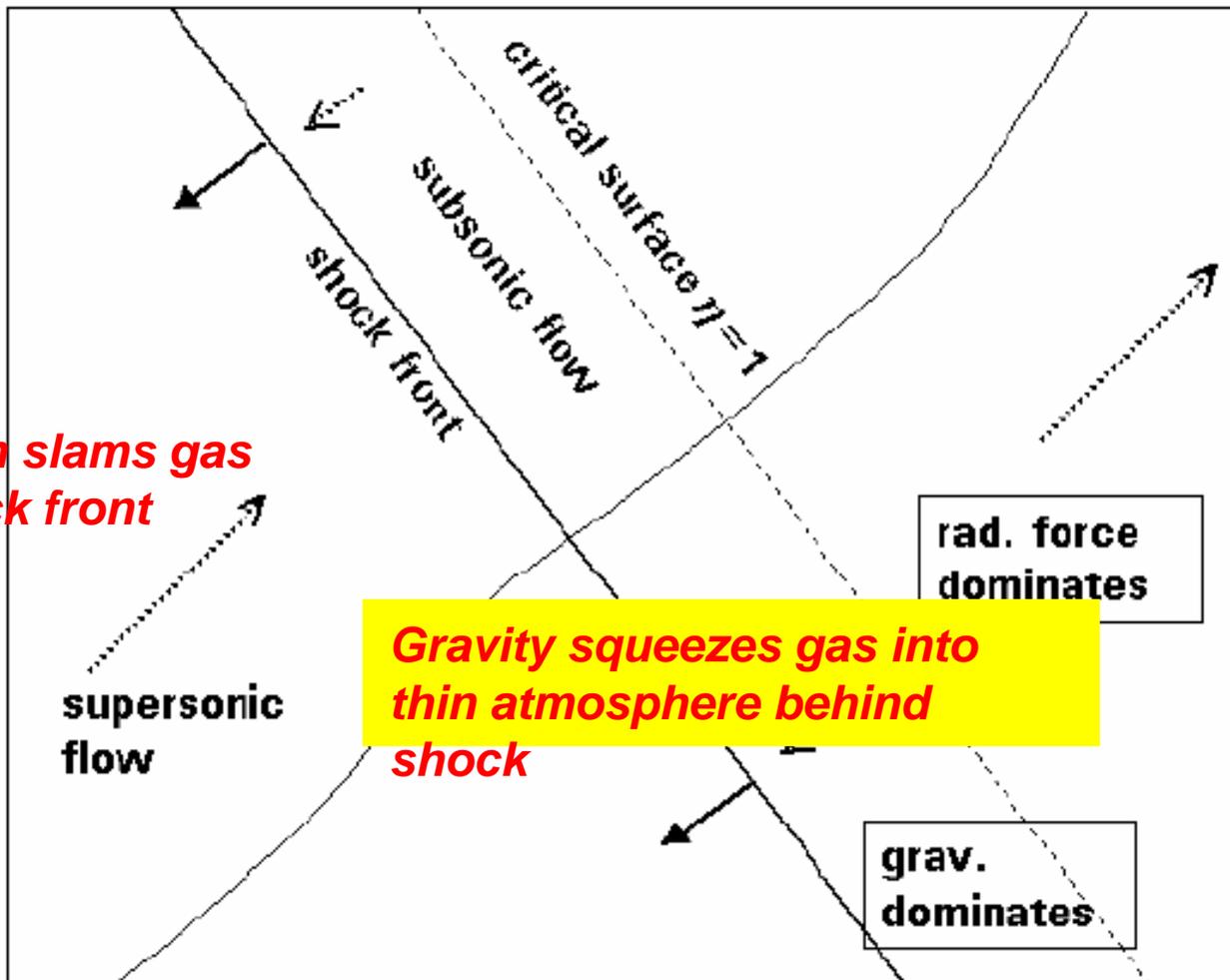
(Begelman 2001)

PHOTON BUBBLE SHOCK TRAINS



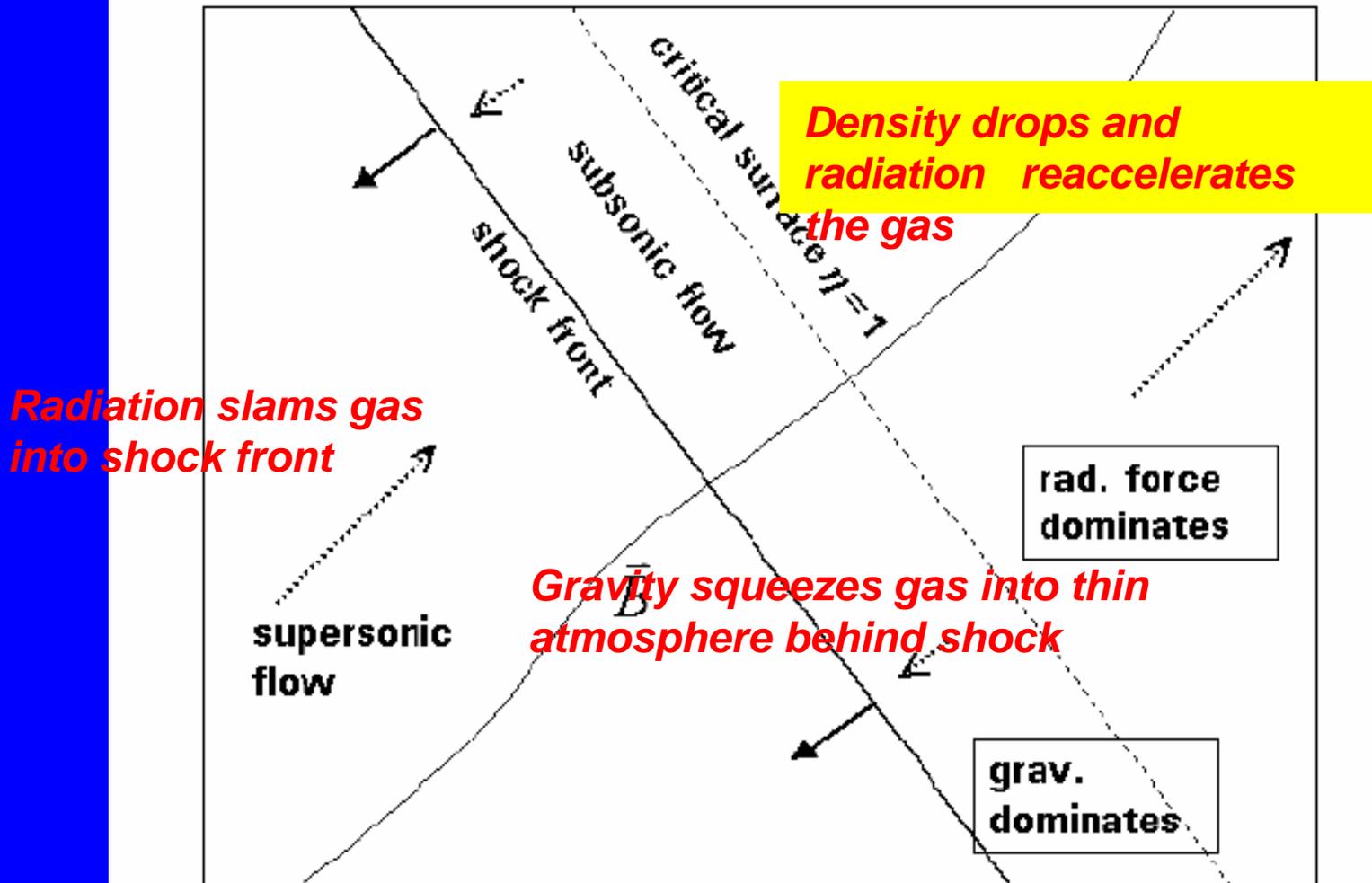
Radiation slams gas into shock front

PHOTON BUBBLE SHOCK TRAINS



PHOTON BUBBLE SHOCK TRAINS

← ...and the process repeats... →



SUPER-EDDINGTON ACCRETION DISKS?

Photon bubbles \Rightarrow porous disk \Rightarrow

$L > L_E$ possible without blowing disk apart

Max. luminosity:

$$\frac{L}{L_E} \sim 40 \left(\frac{\alpha}{0.01} \right) \left(\frac{m}{10} \right)^{1/5} \left(\frac{\xi}{0.1} \right)^{4/5}$$

Near the top of the atmosphere...

- Optical depth must fall
- Weakens density-dependence of flux
- Drives strong wind if super-Eddington

...but...

- Energetics marginal for dispersing the disk (and regulating L to $\sim L_E$)
 - Requires effective radiation trapping
 - magnetically dominated corona may retain and recycle escaping gas

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