

Search for the Electric Dipole Moment of the τ Lepton

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We have searched for a T/CP violation signature arising from an electric dipole form factor (d_τ) of the τ lepton in the $e^+e^- \rightarrow \tau^+\tau^-$ reaction. Using an optimal observable method for 29.5 fb^{-1} of data collected with the Belle detector at the KEKB collider at $\sqrt{s} = 10.58 \text{ GeV}$, we obtained the preliminary result $Re(d_\tau) = (1.15 \pm 1.70) \times 10^{-17} \text{ ecm}$ and $Im(d_\tau) = (-0.83 \pm 0.86) \times 10^{-17} \text{ ecm}$ for the real and imaginary parts of d_τ , respectively, and set the 95% confidence level limits $-2.2 < Re(d_\tau) < 4.5$ (10^{-17} ecm) and $-2.5 < Im(d_\tau) < 0.8$ (10^{-17} ecm).

1. Introduction

While large CP violating asymmetry in the quark sector has recently been confirmed in B-meson decay at B factories [1,2], the Standard Model (SM) does not predict any appreciable CP violation (CPV) in the lepton sector. However, physics beyond the SM, such as multi-Higgs, SUSY, and lepto-quark models [3,4], could result in CPV in leptonic processes; since new bosons and Higgs would strongly couple with heavy particles through quantum loop effects, CP violating phenomena could be significantly enhanced, particularly in τ decay, due to its very large mass compared to other leptons.

The contribution from CP violating interactions in the τ -pair production process can be parametrized in a model-independent way using the electric dipole moment (or form factor) for the τ lepton. Up to now, searches for the electric dipole form factor d_τ have been performed at LEP [5,6] in the reaction $e^+e^- \rightarrow \tau^+\tau^-\gamma$ and by ARGUS [7] in the $e^+e^- \rightarrow \tau^+\tau^-$ reaction and the following upper limits at 95% confidence level were obtained: $-3.1 < d_\tau < 3.1$ (10^{-16} ecm) from L3 [5], $-3.8 < d_\tau < 3.6$ (10^{-16} ecm) from OPAL [6] and $|Re(d_\tau)| < 4.6 \times 10^{-16} \text{ ecm}$, $|Im(d_\tau)| < 1.8 \times 10^{-16} \text{ ecm}$ from ARGUS [7].

We have searched for a form factor d_τ at the $\gamma\tau\tau$ vertex in the process $e^+e^- \rightarrow \gamma^* \rightarrow \tau^+\tau^-$ using triple momentum and spin correlation observables. The effective Lagrangian can be expressed

as

$$\mathcal{L}_{CP} = -id_\tau(s)\bar{\tau}\sigma^{\mu\nu}\gamma_5\tau\partial_\mu A_\nu, \quad (1)$$

where the electric dipole form factor d_τ depends in general on s , the squared energy of the τ -pair system. In common with other analyses we ignore this possible s -dependence, assuming $d_\tau(s) \equiv d_\tau$, constant. The squared spin density matrix ($\mathcal{M}_{\text{prod}}^2$) for the reaction $e^+(\mathbf{p})e^-(-\mathbf{p}) \rightarrow \tau^+(\mathbf{k}, \mathbf{S}_+)\tau^-(-\mathbf{k}, \mathbf{S}_-)$ is given by [8]

$$\mathcal{M}_{\text{prod}}^2 = \mathcal{M}_{\text{SM}}^2 + Re(d_\tau)\mathcal{M}_{Re}^2 + Im(d_\tau)\mathcal{M}_{Im}^2, \quad (2)$$

$$\begin{aligned} \mathcal{M}_{\text{SM}}^2 = & \frac{e^4}{k_0^2} [k_0^2 + m_\tau^2 + |\mathbf{k}|^2(\hat{\mathbf{k}}\hat{\mathbf{p}})^2 - \mathbf{S}_+\mathbf{S}_-|\mathbf{k}|^2(1 - (\hat{\mathbf{k}}\hat{\mathbf{p}})^2)] \\ & + 2(\hat{\mathbf{k}}\mathbf{S}_+)(\hat{\mathbf{k}}\mathbf{S}_-)(|\mathbf{k}|^2 + (k_0 - m_\tau)^2(\hat{\mathbf{k}}\hat{\mathbf{p}})^2) \\ & - 2k_0(k_0 - m_\tau)(\hat{\mathbf{k}}\hat{\mathbf{p}})((\hat{\mathbf{k}}\mathbf{S}_+)(\hat{\mathbf{p}}\mathbf{S}_-) + (\hat{\mathbf{k}}\mathbf{S}_-)(\hat{\mathbf{p}}\mathbf{S}_+)) \\ & + 2k_0^2(\hat{\mathbf{p}}\mathbf{S}_+)(\hat{\mathbf{p}}\mathbf{S}_-), \end{aligned} \quad (3)$$

$$\begin{aligned} \mathcal{M}_{Re}^2 = & 4\frac{e^3}{k_0} |\mathbf{k}| [-(m_\tau + (k_0 - m_\tau)(\hat{\mathbf{k}}\hat{\mathbf{p}})^2)(\mathbf{S}_+ \times \mathbf{S}_-) \hat{\mathbf{k}} \\ & + k_0(\hat{\mathbf{k}}\hat{\mathbf{p}})(\mathbf{S}_+ \times \mathbf{S}_-) \hat{\mathbf{p}}], \end{aligned} \quad (4)$$

$$\begin{aligned} \mathcal{M}_{Im}^2 = & 4\frac{e^3}{k_0} |\mathbf{k}| [-(m_\tau + (k_0 - m_\tau)(\hat{\mathbf{k}}\hat{\mathbf{p}})^2)(\mathbf{S}_+ - \mathbf{S}_-) \hat{\mathbf{k}} \\ & + k_0(\hat{\mathbf{k}}\hat{\mathbf{p}})(\mathbf{S}_+ - \mathbf{S}_-) \hat{\mathbf{p}}], \end{aligned} \quad (5)$$

where k_0 is the energy of the τ , m_τ is the τ mass, \mathbf{p} is the momentum vector of e^+ , \mathbf{k} is the momentum vector of τ^+ in the center-of-mass frame, \mathbf{S}_\pm are the spin vectors for τ^\pm , and the hat denotes a unit momentum. We disregard the higher order terms proportional to $|d_\tau|^2$ assuming $d_\tau(s)$ to be small. $\mathcal{M}_{\text{SM}}^2$ corresponds to the SM term. \mathcal{M}_{Re}^2 and \mathcal{M}_{Im}^2 are the interference terms (related to

the real and imaginary parts of d_τ , respectively) between the SM and CPV amplitudes. \mathcal{M}_{Re}^2 is CP odd and T odd, while \mathcal{M}_{Im}^2 is CP odd, but T even. In the above equations, e^+ and e^- are assumed to be unpolarized and massless particles.

We adapt the so-called optimal observable method [9], which maximizes the sensitivity to d_τ by optimizing the relevant quantities, \mathcal{O}_{Re} and \mathcal{O}_{Im} , for measured kinematic variables, as described below. Here the optimal observables are defined as

$$\mathcal{O}_{Re} = \frac{\mathcal{M}_{Re}^2}{\mathcal{M}_{SM}^2}, \quad \mathcal{O}_{Im} = \frac{\mathcal{M}_{Im}^2}{\mathcal{M}_{SM}^2}. \quad (6)$$

The mean value of the observables is expressed as

$$\begin{aligned} \langle \mathcal{O}_{Re} \rangle &\propto \int \mathcal{O}_{Re} d\sigma \propto \int \mathcal{O}_{Re} \mathcal{M}_{\text{prod}}^2 d\phi \\ &= \int \mathcal{M}_{Re}^2 d\phi + \text{Re}(d_\tau) \int \frac{(\mathcal{M}_{Re}^2)^2}{\mathcal{M}_{SM}^2} d\phi, \end{aligned} \quad (7)$$

where the integration is over the phase space (ϕ) sustained by the relevant kinematic variables. The cross-term containing the integral of the product of \mathcal{M}_{Re}^2 and \mathcal{M}_{Im}^2 drops out because of their different symmetry properties. The expression for the imaginary part is similar. The means of the observables $\langle \mathcal{O}_{Re} \rangle$ and $\langle \mathcal{O}_{Im} \rangle$ are thus expressed as linear functions of d_τ

$$\begin{aligned} \langle \mathcal{O}_{Re} \rangle &= a_{Re} \cdot \text{Re}(d_\tau) + b_{Re}, \\ \langle \mathcal{O}_{Im} \rangle &= a_{Im} \cdot \text{Im}(d_\tau) + b_{Im}. \end{aligned} \quad (8)$$

Eight different final states in the decay of τ -pairs, $(e\nu\bar{\nu})(\mu\nu\bar{\nu})$, $(e\nu\bar{\nu})(\pi\nu)$, $(\mu\nu\bar{\nu})(\pi\nu)$, $(e\nu\bar{\nu})(\rho\nu)$, $(\mu\nu\bar{\nu})(\rho\nu)$, $(\pi\nu)(\rho\nu)$, $(\rho\nu)(\rho\bar{\nu})$, and $(\pi\nu)(\pi\bar{\nu})$, are analyzed, where all particles except ν and $\bar{\nu}$ are positively or negatively charged. Because of the undetectable particles, we can not fully reconstruct the quantities \mathbf{k} and \mathbf{S}_\pm . Therefore, for each event we calculate possible kinematic configurations and obtain the mean value of \mathcal{M}_{SM}^2 , \mathcal{M}_{Re}^2 and \mathcal{M}_{Im}^2 by averaging over the calculated configurations. In the case when both τ leptons decay hadronically ($\tau \rightarrow \pi\nu$ or $\rho\nu$), the τ flight direction is calculated with a two-fold ambiguity and we take the average of \mathcal{M}_{SM}^2 , \mathcal{M}_{Re}^2 and \mathcal{M}_{Im}^2 over the two solutions. In the case when either one or both τ leptons decay leptonically ($\tau \rightarrow e\nu\bar{\nu}$ or $\mu\nu\bar{\nu}$), a Monte Carlo (MC)

treatment is adopted to take into account the additional ambiguity in the effective mass of the $\nu\bar{\nu}$ system ($m_{\nu\bar{\nu}}$). For each event we generate 100 MC configurations using a hit-and-miss approach by varying $m_{\nu\bar{\nu}}$, and compute the averaged \mathcal{M}_{SM}^2 , \mathcal{M}_{Re}^2 and \mathcal{M}_{Im}^2 over successful tries in which the τ direction can be constructed kinematically. In the calculation, we ignore the effect of undetected photons coming from initial state radiation, radiative τ decays and bremsstrahlung. This effect has been examined and found to be negligible compared to the total errors.

2. Data and event selection

In this analysis, we used 26.8 million τ -pairs (29.5 fb^{-1}) accumulated with the Belle detector [10] at the KEKB accelerator [11]. KEKB is an asymmetric energy e^+e^- collider with a beam crossing angle of 22 mrad. Its center-of-mass energy is 10.58 GeV, corresponding to the $\Upsilon(4S)$ resonance, with beam energies of 8 and 3.5 GeV for electrons and positrons, respectively. Belle is a general-purpose detector with an asymmetric structure along the beam direction. Among the detector elements, the central drift chamber (CDC) and the silicon vertex detector (SVD) are essential to obtain the momentum vectors of charged particles. The combined information from the silica Aerogel Cherenkov counters (ACC), the time-of-flight counters (TOF), the CsI electromagnetic calorimeter (ECL), and the μ/K_L detector (KLM) is used for particle identification.

The MC event generators KORALB / TAUOLA [12] are used for τ -pair production and decays. The detector simulation is performed by a GEANT-based program, GSIM. Actual data and MC generated events are reconstructed by the same program written by the Belle collaboration.

We chose the eight final state τ -pair modes mentioned above. All of the final state particles were reconstructed with the following conditions. Each charged track is required to have a transverse momentum $p_t > 0.1 \text{ GeV}/c$. Photon candidates should deposit an energy of $E > 0.1 \text{ GeV}$

Table 1

Yield, purity and background rate obtained by the event selection described in the text, where the purity was evaluated by MC simulation and its error comes from the MC statistics.

	Yield	Purity (%)	Background mode (%)
$e\mu$	250,948	96.6 ± 0.1	$2\gamma \rightarrow \mu\mu(1.9), \tau\tau \rightarrow e\pi(1.1)$.
$e\pi$	132,574	82.5 ± 0.1	$\tau\tau \rightarrow e\rho(6.0), eK(5.4), e\mu(3.1), eK^*(1.3)$.
$\mu\pi$	123,520	80.6 ± 0.1	$\tau\tau \rightarrow \mu\rho(5.7), \mu K(5.3), \mu\mu(2.9), 2\gamma \rightarrow \mu\mu(2.0)$.
$e\rho$	240,501	92.4 ± 0.1	$\tau\tau \rightarrow e\pi\pi^0\pi^0(4.4), eK^*(1.7)$.
$\mu\rho$	217,156	91.6 ± 0.1	$\tau\tau \rightarrow \mu\pi\pi^0\pi^0(4.2), \mu K^*(1.6), \pi\rho(1.0)$.
$\pi\rho$	110,414	77.7 ± 0.1	$\tau\tau \rightarrow \rho\rho(5.1), K\rho(4.9), \pi\pi\pi^0\pi^0(3.8), \mu\rho(2.7)$.
$\rho\rho$	93,016	86.2 ± 0.1	$\tau\tau \rightarrow \rho\pi\pi^0\pi^0(8.0), \rho K^*(3.1)$.
$\pi\pi$	28,348	70.0 ± 0.2	$\tau\tau \rightarrow \pi\rho(9.2), \pi K(9.2), \pi\mu(4.7), \pi K^*(2.0)$.

in the ECL. A signal event is required to have two charged tracks with zero net-charge and no photon apart from $\rho^\pm \rightarrow \pi^\pm \pi^0$, $\pi^0 \rightarrow \gamma\gamma$.

A track was identified as an electron using a likelihood ratio combining dE/dx in the CDC, the ratio of energy deposited in the ECL and momentum measured in the CDC, the shower shape of the ECL and the hit pattern from the ACC. The identification efficiency was estimated to be 92% with a π^\pm fake rate of 0.3% [13]. A muon was evaluated from its range and hit pattern at the KLM detector; its efficiency and fake rate were obtained by MC as 91% and 2%, respectively. A pion track was found by requiring that a track be identified as a hadron by the KLM information, and not identified as an electron. The efficiency was estimated to be 81%. The purity for the obtained samples is about 89%. A ρ^\pm was reconstructed from a charged track and a π^0 where the track should be neither an electron nor a muon, and for $\pi^0 \rightarrow \gamma\gamma$, the reconstructed π^0 should have an invariant mass between 110 and 150 MeV/ c^2 and a momentum in the laboratory frame larger than 0.2 GeV/ c . In order to suppress the background and to enhance the particle identification ability, the detection of leptonic particles was restricted to within the barrel region, $-0.60 < \cos\theta^{\text{lab}} < 0.83$, and that of single pions to within the KLM barrel region, $-0.50 < \cos\theta^{\text{lab}} < 0.62$, where θ^{lab} is the polar angle relative to the e^- beam direction. For the same reason, the particle momentum was required to be greater than 0.5 GeV/ c for an electron, 1.2 GeV/ c for both a muon and pion, and

1.0 GeV/ c for ρ^\pm in the laboratory frame.

The dominant backgrounds are due to two-photon as well as Bhabha and $\mu\mu$ processes. In order to remove two-photon events, we required the missing momentum not to be directed towards the beam-pipe region (imposing a selection $-0.950 < \cos\theta^{\text{lab}} < 0.985$), and to reject the latter processes we required that the sum of the charged track momenta be less than 9 GeV/ c in the center-of-mass frame. Additional selections were imposed particularly on the $e\pi$ mode where a large number of Bhabha events could contribute through misidentification. For the $e\pi$ mode, we remove events which satisfy the following criteria: the opening angle of the two tracks in the plane perpendicular to the beam axis is greater than 175° , and their momentum sum is greater than 6 GeV/ c in the τ -pair rest frame. Finally, we removed events in which the τ flight direction could not be kinematically reconstructed, which mostly arise from τ -pairs having hard initial-state radiation and misidentified τ -pair backgrounds.

The yield of events passing this selection is given in Table 1 for each of the eight selected modes. The mean energy of the τ -pair system in the obtained sample is $\sqrt{s} = 10.38$ GeV; this sets the scale at which $d_\tau(s)$ is measured. Because of events with soft radiated photons, the energy scale is slightly lower than the beam energy. The dominant background sources are also listed in the table. Hadronic τ decays with two or more π^0 's make a large contribution of a few percent. For the modes including π^\pm , the misidentification of kaons and muons as pions yields other back-

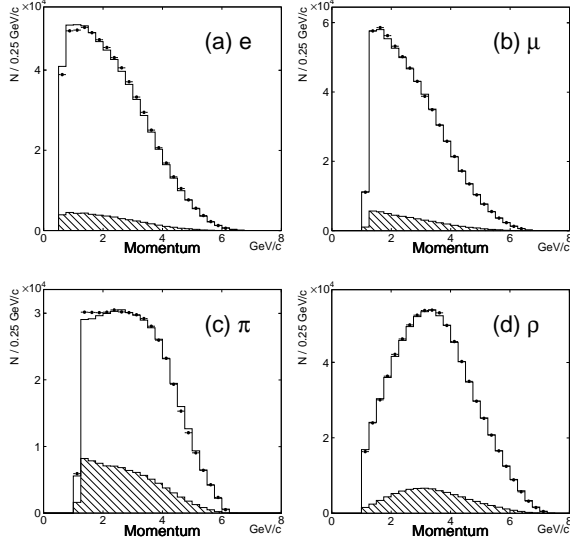


Figure 1. Momentum distributions of (a) e^\pm , (b) μ^\pm , (c) π^\pm , and (d) ρ^\pm in the laboratory frame. The points with error bars are the data and the histogram is the MC expectation. The latter is scaled to the total number of entries. The hatched histogram is the background distribution evaluated by MC.

grounds; for example, 5.3% of μK final states is included in the $\mu\pi$ mode. The other backgrounds are estimated by MC to be a few percent from two-photon processes, and less than 1% from Bhabha, $\mu\mu$, and multihadronic processes.

Figs 1 and 2 show the resulting momentum and $\cos\theta^{lab}$ distributions, respectively, for charged particles in the laboratory frame. Very good agreement with MC is found, except for low-momentum electrons (Fig. 1(a)) and pions (Fig. 1(c)). The dip in the $\cos\theta$ distribution of the muon (Fig. 2(b)) is due to an efficiency drop at the overlapping region between the barrel and endcap KLM elements.

The resultant \mathcal{O}_{Re} and \mathcal{O}_{Im} distributions are shown in Fig. 3 along with those obtained from MC simulation with $d_\tau = 0$. Good agreement is found between the experimental data and the MC samples. The ratio of the data to MC shows flat

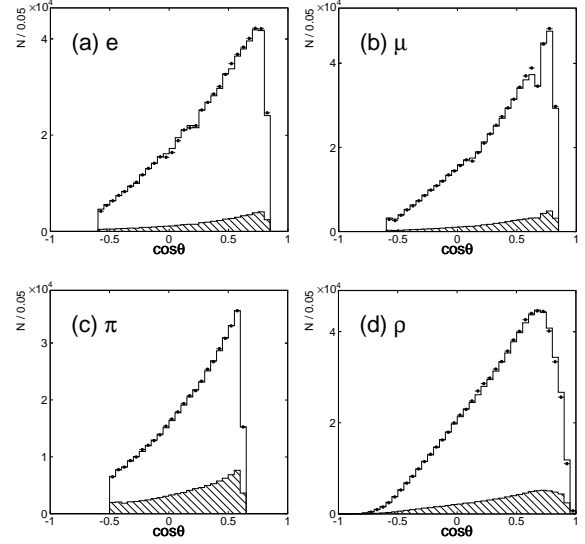


Figure 2. The $\cos\theta$ distributions of (a) e^\pm , (b) μ^\pm , (c) π^\pm , and (d) ρ^\pm in the laboratory frame. The meanings of the points and histograms are the same as in Fig. 1.

distributions around 1.0, although they are not shown here.

3. Extraction of d_τ

In order to extract the d_τ value from the observable using Eq. (8), we have to know the coefficient a and the offset b . In the ARGUS analysis [7] which also used the optimal observable method, the first term of Eq. (7) was assumed to be negligible because of the property of the CP violating term under the ideal detection hypothesis $\int \mathcal{M}_{Re}^2 d\phi = 0$. The value of d_τ was thus obtained as the ratio of the observable's mean to the second term, $Re(d_\tau) = \langle \mathcal{O}_{Re} \rangle / \langle \mathcal{O}_{Re}^2 \rangle$. Experimentally, the detector acceptance η affects the means of the observables as

$$\langle \mathcal{O}_{Re} \rangle \propto \int \eta(\phi) \mathcal{O}_{Re} \mathcal{M}_{prod}^2 d\phi. \quad (9)$$

Similar expressions are obtained for the imaginary part. This means that the first term of Eq. (7) is not necessarily zero and the coefficient

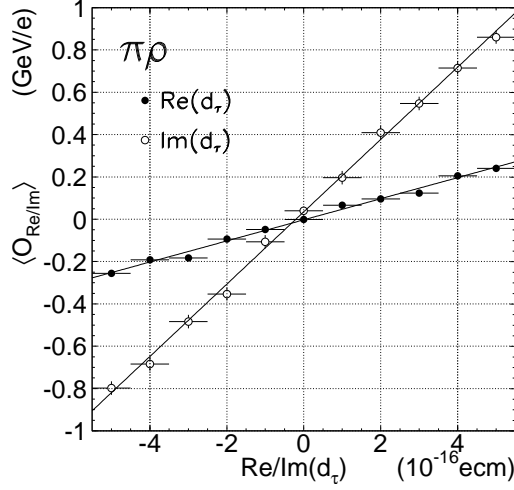


Figure 4. $Re(d_\tau)(Im(d_\tau))$ dependence of the mean of the observable $\langle \mathcal{O}_{Re} \rangle (\langle \mathcal{O}_{Im} \rangle)$ for the $\pi\rho$ mode. The closed circles show the dependence for $Re(d_\tau)$ and the open circles show the dependence for $Im(d_\tau)$. The lines show the fitted linear functions.

may differ from $\langle \mathcal{O}_{Re}^2 \rangle$ when the detector acceptance is taken into account. In their study, the acceptance effect produced the largest systematic uncertainty, on the order of 10^{-16} ecm .

In order to reduce this systematic effect, we use both parameters a and b extracted from the correlation between $\langle \mathcal{O}_{Re} \rangle (\langle \mathcal{O}_{Im} \rangle)$ and $Re(d_\tau)(Im(d_\tau))$ obtained by a full MC including acceptance effects. An example of the correlation between $\langle \mathcal{O}_{Re} \rangle (\langle \mathcal{O}_{Im} \rangle)$ and $Re(d_\tau)(Im(d_\tau))$ is shown in Fig. 4. Each point is obtained from MC with detector simulation and event selection. By fitting the correlation plot with Eq. (8), the parameters a and b are obtained. The misidentified background from τ -pair events shows some dependence on d_τ , because the spin direction is correlated with the momenta of the final state particles. Therefore, both the coefficient a and offset b are corrected with the parameters obtained from the MC misidentified background. The resultant coefficients and offsets are shown in Fig. 5.

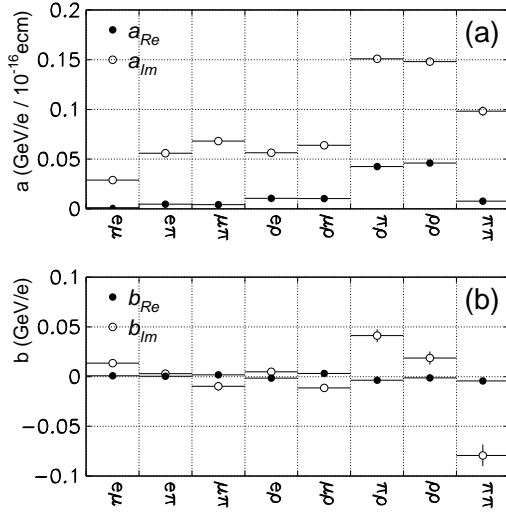


Figure 5. Sensitivity a (a) and offset b (b) for each mode from MC. The closed circles show the parameters for $Re(d_\tau)$ and the open circles show the parameters for $Im(d_\tau)$. The errors are due to MC statistics.

As can be seen from the values of the coefficient a showing the sensitivity to d_τ , the $\pi\rho$ and $\rho\rho$ modes have the highest sensitivities, while the $\pi\pi$ mode yields a somewhat lower value. This is because both \mathbf{S} and \mathbf{k} in the $\pi\pi$ mode are inferred only from the π momentum vector, so that taking the mean value of two possible solutions with equal weights significantly reduces the achievable sensitivity. On the other hand, the momentum sum and angular distribution of $\rho \rightarrow \pi\pi^0$ decay in the $\pi\rho$ and $\rho\rho$ modes provide different weights on the possible solutions, and therefore there is less reduction in sensitivity. Other modes, which include leptons, have low sensitivities because of a lack of information about the τ^\pm directions and spin orientations due to the additional missing neutrinos.

Non-zero offsets b_{Im} are seen for the imaginary part, due to the forward/backward asymmetry in the acceptance of the detector.

Table 2
Systematic errors for $Re(d_\tau)$ and $Im(d_\tau)$ in units of 10^{-16} ecm .

$Re(d_\tau)$	$e\mu$	$e\pi$	$\mu\pi$	$e\rho$	$\mu\rho$	$\pi\rho$	$\rho\rho$	$\pi\pi$
Mismatch of distribution	0.80	0.58	0.70	0.11	0.15	0.21	0.16	0.06
Charge asymmetry	0.00	0.01	0.01	0.01	0.01	0.01	-	-
Background variation	0.43	0.12	0.07	0.07	0.08	0.03	0.04	0.05
Momentum reconstruction	0.16	0.09	0.24	0.04	0.06	0.06	0.04	0.45
Detector alignment	0.02	0.02	0.01	0.00	0.01	0.01	0.02	0.03
Total	0.92	0.60	0.74	0.14	0.18	0.22	0.17	0.46
$Im(d_\tau)$	$e\mu$	$e\pi$	$\mu\pi$	$e\rho$	$\mu\rho$	$\pi\rho$	$\rho\rho$	$\pi\pi$
Mismatch of distribution	0.43	0.02	0.05	0.12	0.01	0.05	0.10	0.41
Charge asymmetry	0.13	0.44	0.43	0.02	0.09	0.15	-	-
Background variation	0.08	0.07	0.02	0.01	0.03	0.02	0.03	0.06
Momentum reconstruction	0.03	0.03	0.06	0.00	0.02	0.02	0.04	0.04
Detector alignment	0.01	0.03	0.02	0.01	0.01	0.02	0.01	0.05
Total	0.46	0.45	0.44	0.13	0.10	0.16	0.11	0.42

4. Systematic uncertainties

Although in general the MC simulation reproduces the observed kinematic distributions well, the disagreement seen in Figs. 1 and 2 causes an error dominating the systematic uncertainty. Its uncertainty with respect to d_τ was studied by reweighting the MC distributions by the ratio of data to MC. The second significant uncertainty originates from possible charge asymmetry in the detection efficiency. The ratio of yields, $N(\alpha^+\beta^-)/N(\alpha^-\beta^+)$, for data and MC agrees within 1% accuracy, where α and β are the relevant charged particles from the τ decays. This systematic uncertainty was evaluated by varying the detection efficiency by $\pm 1\%$. The effect is of the same size as the statistical error for $Im(d_\tau)$, while it is negligible for $Re(d_\tau)$. The backgrounds provide additional systematic errors for d_τ because the parameters a and b are corrected for the background contributions, which are assessed by varying the background rate. The influence of the momentum reconstruction of charged particles and photons was also checked by applying scaling factors corresponding to the difference between the data and MC. In order to examine a possible asymmetry arising from the alignment of the tracking devices, we measured the differences of the polar angles, $\Delta\theta = \theta_{\text{CM}}^+ - \theta_{\text{CM}}^-$, and the azimuthal angles, $\Delta\phi = \phi_{\text{CM}}^+ - \phi_{\text{CM}}^-$, between

two tracks in $e^+e^- \rightarrow \mu^+\mu^-$ events, and found a small deviation from a back-to-back topological alignment in each direction as $\Delta\theta = 1.48 \text{ mrad}$ and $\Delta\phi = 0.36 \text{ mrad}$. By applying an artificial angular deviation of this magnitude to one of the charged tracks, the residual value of the observables is calculated to be negligible compared with the other errors. The systematic errors are listed in Table 2.

5. Result

The values of d_τ extracted using Eq. (8) are listed in Table 3 along with the statistical and systematic errors. All results are consistent with zero within the errors.

Finally, we obtain mean values for $Re(d_\tau)$ and $Im(d_\tau)$ over the eight different $\tau^+\tau^-$ modes weighted by quadratically summed statistical and systematic errors. The resultant preliminary result for the electric dipole form factors are

$$Re(d_\tau) = (1.15 \pm 1.70) \times 10^{-17} \text{ ecm}, \quad (10)$$

$$Im(d_\tau) = (-0.83 \pm 0.86) \times 10^{-17} \text{ ecm}, \quad (11)$$

with corresponding 95% confidence limits

$$-2.2 < Re(d_\tau) < 4.5 \quad (10^{-17} \text{ ecm}), \quad (12)$$

$$-2.5 < Im(d_\tau) < 0.8 \quad (10^{-17} \text{ ecm}). \quad (13)$$

This investigation has improved the sensitivity to the τ lepton's electric dipole form factor by

Table 3

Results for the electric dipole form factor. The first error is statistical and the second is systematic.

	$Re(d_\tau) (10^{-16} ecm)$	$Im(d_\tau) (10^{-16} ecm)$
$e\mu$	$2.25 \pm 1.26 \pm 0.92$	$-0.41 \pm 0.22 \pm 0.46$
$e\pi$	$0.43 \pm 0.64 \pm 0.60$	$-0.22 \pm 0.19 \pm 0.45$
$\mu\pi$	$-0.41 \pm 0.87 \pm 0.74$	$0.15 \pm 0.19 \pm 0.44$
$e\rho$	$0.00 \pm 0.36 \pm 0.14$	$-0.01 \pm 0.14 \pm 0.13$
$\mu\rho$	$0.04 \pm 0.42 \pm 0.18$	$-0.02 \pm 0.14 \pm 0.10$
$\pi\rho$	$0.34 \pm 0.25 \pm 0.22$	$-0.22 \pm 0.13 \pm 0.16$
$\rho\rho$	$-0.08 \pm 0.25 \pm 0.17$	$-0.12 \pm 0.14 \pm 0.11$
$\pi\pi$	$0.42 \pm 1.17 \pm 0.46$	$0.24 \pm 0.34 \pm 0.42$
All	0.115 ± 0.170	-0.083 ± 0.086

a factor of more than 10 compared to previous measurements.

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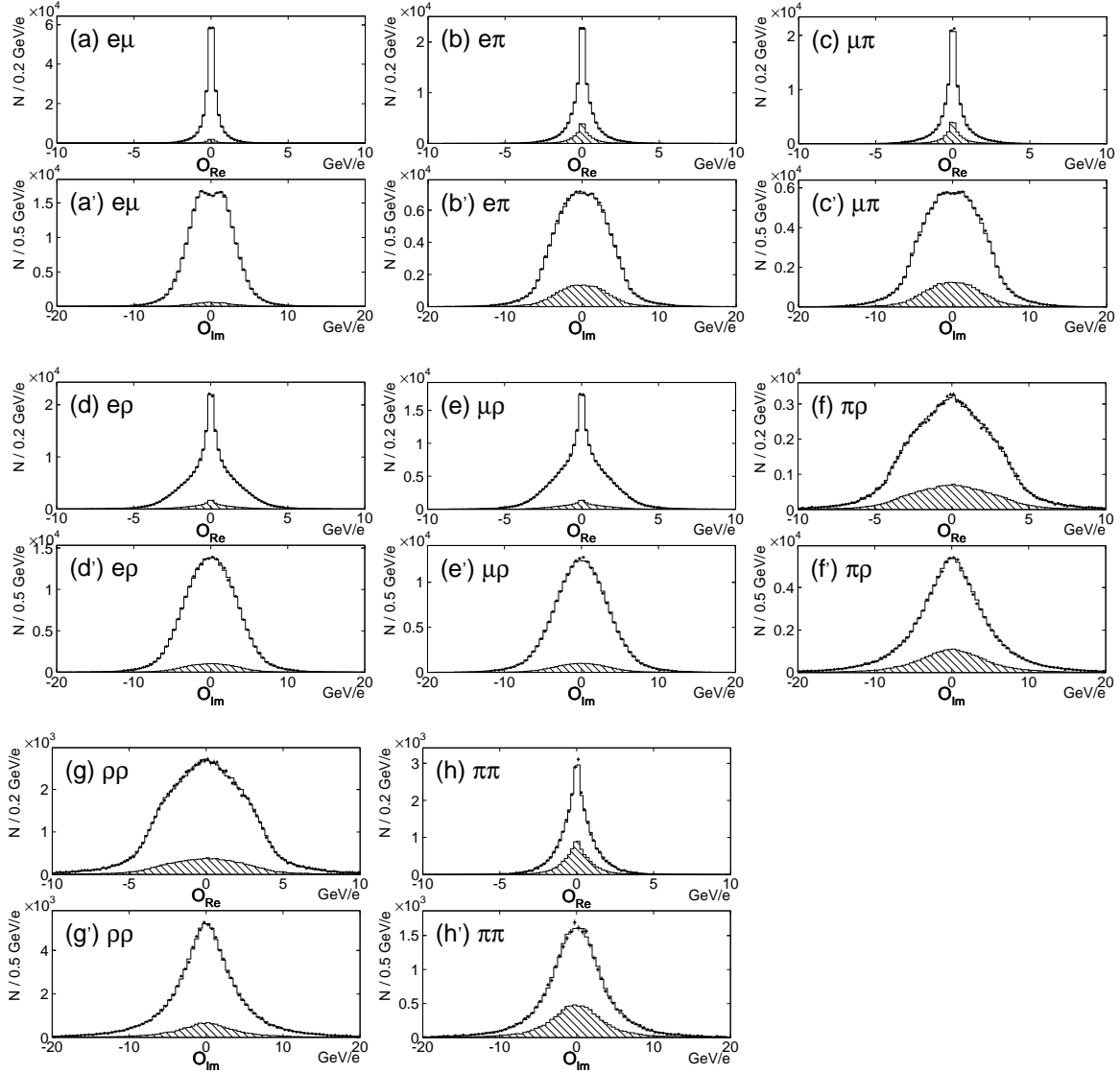


Figure 3. Distributions of the optimal observables \mathcal{O}_{Re} and \mathcal{O}_{Im} for each mode. The upper figure of each mode is for \mathcal{O}_{Re} and lower figure is for \mathcal{O}_{Im} . The closed circles are the experimental data and the histogram is the MC expectation with $d_\tau = 0$, normalized to the number of entries. The hatched histogram is the background distribution evaluated by MC.