
$B^- \rightarrow \omega K^- / \pi^-$ and Time-dependent CPV in $B^0 \rightarrow \eta' K_S$ at Belle

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1 Introduction

Charmless hadronic B decays are interesting for testing our understanding of heavy quark physics and search for CP violation. We report the measurements of charmless $B^- \rightarrow \omega \pi^-$ and $B^- \rightarrow \omega K^-$ decays. The $B^- \rightarrow \omega K^-$ mode was first reported by CLEO collaboration in 1998 [1], but superseded by nonobservation with a larger data set [2]. This result is supported by BaBar in 2001 [3]. The $B^- \rightarrow \omega \pi^-$ channel has also been reported by CLEO and BaBar collaborations. However, we find a significant signal in the $B^- \rightarrow \omega K^-$ mode. The data used in this analysis were based on an integrated luminosity of 29.4 fb^{-1} data sample on the $\Upsilon(4S)$ resonance taking by the Belle detector at the KEK B-factory.

Kobayashi and Maskawa (KM) showed CP violation could be accommodated in the Standard Model [4] (SM) in 1973. In the model, the CP violation arises from a complex phase in the weak interaction quark mixing matrix [5]. Belle and BaBar collaborations have measured the CP -violating parameter $\sin 2\phi_1$ from $B^0 \rightarrow c\bar{c}K^0$ decays. The $B^0 \rightarrow \eta' K_S$ decay is also sensitive to the parameter $\sin 2\phi_1$ and provides important tests of the KM model. Within the framework of SM, the charmless decay $B^0 \rightarrow \eta' K_S$ are dominated by $b \rightarrow s$ penguin diagrams and a small color-suppressed $b \rightarrow u$ tree contribution. However, since the branching fractions of $B \rightarrow \eta' K$ decays are larger than current theoretical prediction, there might be an additional new phase beyond SM.

The time-dependent CP asymmetry is given by

$$\begin{aligned} A_{CP}(\Delta t) &= \frac{\Gamma(\overline{B^0} \rightarrow \eta' K_S) - \Gamma(B^0 \rightarrow \eta' K_S)}{\Gamma(\overline{B^0} \rightarrow \eta' K_S) + \Gamma(B^0 \rightarrow \eta' K_S)} \\ &= A_{\eta' K_S} \cos(\Delta m \Delta t) + S_{\eta' K_S} \sin(\Delta m \Delta t), \end{aligned}$$

where the CP -violating parameters $A_{\eta' K_S}$ and $S_{\eta' K_S}$ are expressed as

$$A_{\eta'K_S} = \frac{|\lambda|^2 - 1}{|\lambda|^2 + 1}, \quad S_{\eta'K_S} = \frac{2\text{Im}\lambda}{|\lambda|^2 + 1}, \quad (1)$$

where λ is a complex parameter that depends on both $B^0 - \bar{B}^0$ mixing and on the decay amplitude for $B^0(\bar{B}^0) \rightarrow \eta'K_S$. $|\lambda|$ is close to 1 in the SM with a good approximation. The measurement of CP -violating parameters in the decay $B \rightarrow \eta'K$ is based on a 41.8 fb^{-1} data sample.

The KEKB asymmetric e^+e^- (3.5 GeV e^+ and 8 GeV e^-) collider operates on the $\Upsilon(4S)$ resonance. The $\Upsilon(4S)$ is produced with a Lorentz boost of $\beta\gamma = 0.425$ along the beam direction (z).

The Belle detector [6] is a general-purpose detector operating under 1.5 T magnetic field. It consists of a three layer silicon vertex detector (SVD), a central drift chamber (CDC) with 50 layers. Charged hadrons are distinguished by combining the information from an array of aerogel threshold Cerenkov counters (ACC), a time-of-flight scintillation counters (TOF), and the dE/dx measurements from CDC. Photons are detected in an array of 8736 CsI(Tl) crystals (ECL) located inside the magnet. The flux-return outside of the coil is instrumented for detection of muon and K_L (KLM).

2 Observation of $B^\pm \rightarrow \omega K^\pm$ Decay

Charged tracks are required to be inconsistent with electrons or muons, and are identified as kaons or pions according to a K/π likelihood ratio (KID), $\mathcal{L}_K/(\mathcal{L}_K + \mathcal{L}_\pi)$, where the $\mathcal{L}_{K(\pi)}$ is given by

$$\mathcal{L}_{K(\pi)} = \mathcal{L}_{K(\pi)}^{\text{ACC}} \times \mathcal{L}_{K(\pi)}^{\text{TOF}} \times \mathcal{L}_{K(\pi)}^{dE/dx}, \quad (2)$$

where $\mathcal{L}_{K(\pi)}^{\text{ACC}}$, $\mathcal{L}_{K(\pi)}^{\text{TOF}}$, and $\mathcal{L}_{K(\pi)}^{dE/dx}$ are the likelihood for ACC, TOF, and dE/dx from CDC measurement respectively. Neutral pions are reconstructed from pairs of photons, each consisting of energy greater than 50 MeV, with the $m_{\gamma\gamma}$ within a $\pm 3\sigma$ ($\sigma = 5.4 \text{ MeV}/c^2$) mass window around the π^0 mass. A mass constrained fit is applied on the π^0 candidates. ω meson candidates are obtained by combining $\pi^+\pi^-\pi^0$ with invariant mass within $\pm 30 \text{ MeV}/c^2$ of the nominal ω mass. The momentum of candidate π^0 's under ω is required to be greater than $350 \text{ MeV}/c$. This cut gives 84% efficiency, but removes 60% combinatorial background.

A charged kaon or a pion track is selected to combine with an ω candidate to form a B^\pm candidate. B mesons are identified by two kinematic variables, the beam-constrained mass $M_{bc} = \sqrt{(E_{\text{beam}}^{\text{CM}})^2 - (p_B^{\text{CM}})^2}$ and the energy difference $\Delta E = E_B^{\text{CM}} - E_{\text{beam}}^{\text{CM}}$, where the $E_{\text{beam}}^{\text{CM}} = \sqrt{s}/2 = 5.29 \text{ GeV}$ is the CM beam energy. The resolutions are $3 \text{ MeV}/c^2$ for M_{bc} and 23 MeV for ΔE from Monte Carlo simulations.

Continuum $e^+e^- \rightarrow q\bar{q}$ process is the dominant background in this analysis. The $q\bar{q}$ consists of light quark ($u, d, s, \text{ or } c$) pairs. The continuum events which are jet-like

have to be suppressed compared to the spherical $B\bar{B}$ events. A Fisher discriminant [7] is defined containing modified Fox-Wolfram moments [8]. We use a likelihood ratio method combining the Fisher discriminant, B -meson direction, and the helicity angle which is defined as the angle between the B flight direction and the vector perpendicular to the ω decay plane in the ω rest frame. The probability density functions (PDFs) are determined by Monte Carlo simulation. A signal (background) likelihood \mathcal{L}_s (\mathcal{L}_b) is calculated by the PDFs for each event. The likelihood ratio is formed by $LR = \mathcal{L}_s/(\mathcal{L}_s + \mathcal{L}_b)$, and a cut $LR > 0.85$ is made to reject 95% continuum background but keeping 50% signal events.

Background from charmed and charmless B decays such as $B \rightarrow \omega K^*$, $B \rightarrow \omega \rho$, and non-resonant $B^- \rightarrow K^- \pi^+ \pi^- \pi^0$ are studied by Monte Carlo simulation up to 20 times larger than current data amount. These backgrounds are negligible in the M_{bc} - ΔE signal region ($M_{bc} > 5.27 \text{ GeV}/c^2$ and $|\Delta E| < 0.1 \text{ GeV}$).

The signal yields are extracted by an unbinned maximum likelihood fit on M_{bc} and ΔE together. The signal PDF for M_{bc} is a Gaussian, and for ΔE [9], parameters are determined by MC simulation and calibrated by the $B^- \rightarrow D^0 \pi^-$ decay, where D^0 decays to $K^- \pi^+ \pi^0$. Background PDFs are threshold functions [10] determined from MC simulation. The yields are summarized in Table 1. The M_{bc} and ΔE projections of candidate events are shown in Figure 1.

The systematic uncertainties associated with the fit are determined by varying the parameters in the fitting functions by $\pm 1\sigma$ of their nominal values. The systematic error of ω detection efficiency is determined from detailed studies of charged particle tracking, π^0 detection, and particle identification. A 5% systematic error is assigned to the continuum suppression cut from $B^- \rightarrow D^{*0} \pi^-$ data and MC study. The combined uncertainty of the efficiency is 10.1%.

We make the consistency check of our analysis without KID requirement. Figure 2 shows the ΔE distribution and scatter plot of KID versus ΔE . The ωK^- signals are peaking at -51 MeV , where the $\omega \pi^-$ signals are placed at zero. The ΔE distributions confirm the ωK^- signals are larger than $\omega \pi^-$. We also examine the ω properties to confirm the ωK^- signal candidates. The ω invariant mass spectrum and the cosine of the helicity angle of the B candidate events are also shown in Figure 2.

We also measure the asymmetry in $B^\pm \rightarrow \omega K^\pm$ decays to search for direct CP violation. The asymmetry is defined as

$$A_{CP} = \frac{N(\omega K^-) - N(\omega K^+)}{N(\omega K^-) + N(\omega K^+)}. \quad (3)$$

The asymmetry is measured to be $-0.21 \pm 0.28 \pm 0.03$. The 90% confidence level interval is given by $-0.70 < A_{CP} < 0.28$.

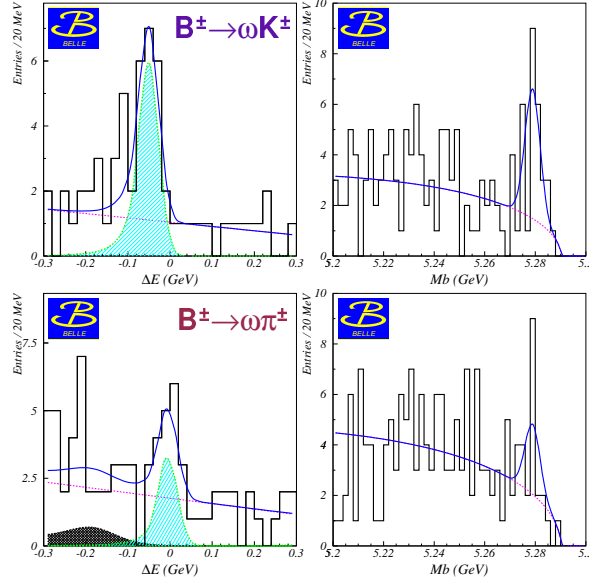


Figure 1: The M_{bc} and ΔE projection plots of $B^- \rightarrow \omega K^-$ and $B^- \rightarrow \omega \pi^-$ modes.

3 Search of Time-dependent CP Asymmetry in $B \rightarrow \eta' K$ Decay

The $B \rightarrow \eta' K$ event selection is described in the previous published branching ratio measurement [11]. The CP fit analysis is done nearly the same as the $B \rightarrow \pi^+ \pi^-$ case.

The $K_S \rightarrow \pi^+ \pi^-$ candidates are reconstructed with a pair of charged tracks with pion mass assumption. A ± 16 MeV/ c^2 to the nominal K_S mass window cut is applied. Two decay channels are used for η' reconstruction: $\eta' \rightarrow \eta \pi^+ \pi^-$ with the η decays to two photons and $\eta' \rightarrow \rho^0 \gamma$. Candidate photons from η (η') are required to have a energy greater than 50 MeV (100 MeV). The B meson candidates are formed by combining η' and a K_S (K^\pm). The kinematic variables, ΔE and M_{bc} are used to identify B signal candidates.

The dominant background to $B \rightarrow \eta' K$ modes are $e^+ e^- \rightarrow q \bar{q}$ continuum events, just like the $B \rightarrow \omega h$ modes. A loose cut of $|\cos \theta_T| < 0.9$ [11] is used rejecting 50% of the

	Signal Yield	Σ	$\epsilon(\%)$	$BF (\times 10^{-6})$	$UL (\times 10^{-6})$
ωK^-	$19.7^{+5.8+0.7}_{-4.7-0.5}$	6.4σ	6.3	$9.9^{+2.7}_{-2.4} \pm 1.0$	-
$\omega \pi^-$	$10.6^{+4.3+0.4}_{-4.5-0.6}$	3.3σ	7.7	$4.3^{+2.0}_{-1.8} \pm 0.5$	< 8.2

Table 1: The results of $B^- \rightarrow \omega K^-$ and $B^- \rightarrow \omega \pi^-$ decays.

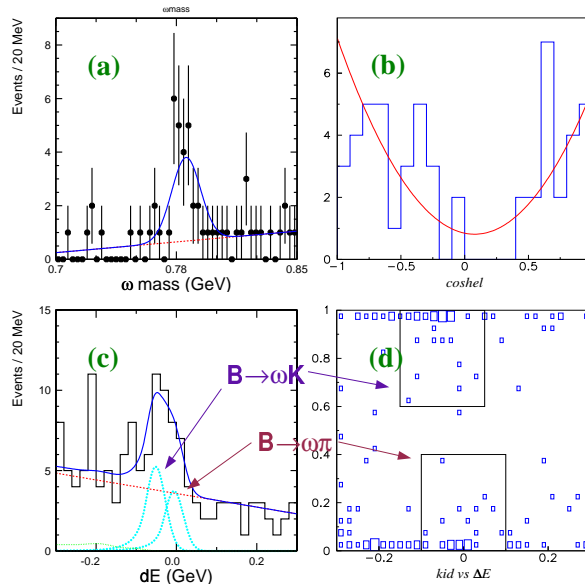


Figure 2: (a) The $\pi^+\pi^-\pi^0$ invariant mass spectrum. (b) The cosine of the helicity angle of $B\rightarrow\omega h$ modes. (c) The ΔE distribution without KID information. (d) The Belle KID versus ΔE spectrum.

background but retaining 90% of signal events.

The $\eta'\rightarrow\eta\pi^+\pi^-$ mode is clean enough after the $\cos\theta_T$ selection, but the $\eta'\rightarrow\rho^0\gamma$ is not as clean. We introduce a likelihood ratio method which is very close to the $B\rightarrow\omega h$ analysis. The only difference is the helicity angle of $\eta'\rightarrow\rho^0\gamma$ mode is defined with ρ^0 in η' rest frame. Finally, a $LR > 0.5$ is applied to reject 81% of the background and keep 80% of the signal. Other generic $B\bar{B}$ backgrounds mainly come from the decay of B to charm. They contribute 8.1% for the $\eta'(\rightarrow\rho^0\gamma)K_S$ mode but negligible in the $\eta'\rightarrow\eta\pi^+\pi^-$ modes. The projection plots for ΔE and M_{bc} of $\eta'K^\pm$ and $\eta'K_S$ decays are shown in Figure 3.

Leptons, charged pions, Λ 's and kaons are not associated with the reconstructed $B^0\rightarrow\eta'K_S$ decay are used to identify the flavor of the accompanying B meson. The algorithm is the same as the $\sin 2\phi_1$ measurement. The vertex positions for $\eta'K_S$ and B_{tag} decays are reconstructed using tracks with at least two SVD hits in z axis and tracks from K_S candidates are removed.

The proper-time interval resolution for the signal, $R_{\text{sig}}(\Delta t)$, is obtained by convolving a sum of two Gaussians. A main component due to the SVD vertex resolution and charmed meson lifetimes, with a tail component caused by poorly reconstructed tracks. The mean and width of the Gaussians are calculated event-by-event from the vertex fit error and the χ^2 values. The background resolutions, $R_{q\bar{q}}(\Delta t)$ and $R_{B\bar{B}}(\Delta t)$, have the same form but the parameters are determined from sideband

events and generic B Monte Carlo. We obtain the lifetimes of charged and neutral B mesons with the same procedure, and the results agree with the world average values.

The CP violation parameters, $A_{\eta'K_S}$ and $S_{\eta'K_S}$, are determined by performing an unbinned maximum-likelihood fit to the measured Δt distributions. The expected PDF for signal,

$$P_{\text{sig}}(\Delta t) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}} \{1 + q \cdot (1 - 2w_l) \cdot A_{CP}(\Delta t)\}, \quad (4)$$

is used in the fit. The τ_{B^0} and Δm_d are fixed to their world average values [12]. The PDF used for the background distribution is

$$P_{\text{bkg}}(\Delta t) = \{f_\tau \cdot \frac{e^{-|\Delta t|/\tau_{\text{bkg}}}}{2\tau_{\text{bkg}}} + (1 - f_\tau) \cdot \delta(\Delta t)\}/2, \quad (5)$$

where f_τ is the background fraction with an effective lifetime τ_{bkg} . Parameters for continuum background are determined from sideband data. The PDF shape for $B\bar{B}$ background is from Monte Carlo simulation.

The likelihood value for each event are defined as:

$$P_i = \int f_{\text{sig}}^l P_{\text{sig}}(\Delta t') R_{\text{sig}}(\Delta t_i - \Delta t') + f_{q\bar{q}}^l P_{q\bar{q}}(\Delta t') R_{q\bar{q}}(\Delta t_i - \Delta t') \\ + f_{B\bar{B}}^l P_{B\bar{B}}(\Delta t') R_{B\bar{B}}(\Delta t_i - \Delta t') d\Delta t', \quad (6)$$

where the f_k^l (subscript $k = \text{sig}, q\bar{q}, \text{ or } B\bar{B}$, and $l = 6$) are the weighted probability functions determined on an event-by-event basis as a function of ΔE and M_{bc} .

In the fit, $A_{\eta'K_S}$ and $S_{\eta'K_S}$ are free parameters that are determined by maximizing the likelihood function $\mathcal{L} = \prod P_i$ over all $B^0 \rightarrow \eta'K_S$ candidates. The CP asymmetry parameters obtained from a total 73 $\eta'K_S$ signal events after subtracting the background are

$$S_{\eta'K_S} = 0.28 \pm 0.55 \text{ (stat)}_{-0.08}^{+0.07} \text{ (syst)}, \\ A_{\eta'K_S} = 0.13 \pm 0.32 \text{ (stat)}_{-0.06}^{+0.09} \text{ (syst)}.$$

Figure 4 shows the Δt distributions for B^0 (\bar{B}^0) tagged events, and raw asymmetry of the $\eta'K_S$ sample.

We determine the systematic error due to fitting parameters (physics parameters, wrong tag fractions, signal and background PDFs) by repeating the fit after varying the parameters by $\pm 1\sigma$ determined from the data or MC.

A number of checks are also performed. The $\eta'K^\pm$ sample is analyzed in the same way as $\eta'K_S$. A fit to 230 candidate events give $S = 0.11 \pm 0.29$ and $A = -0.27 \pm 0.17$ consistent with no asymmetry. The other check by floating the B^\pm (B^0) lifetime

in the B^\pm (B^0) sample is performed. The B^\pm (B^0) lifetime determined with $\eta'K^\pm$ ($\eta'K_S$) samples is $1.54^{+0.14}_{-0.13}$ ($1.58^{+0.31}_{-0.26}$) ps, consistent with world average value [12].

In summary, we measure the time-dependent CP asymmetry parameters in $B^0(\overline{B}^0) \rightarrow \eta'K_S$ decay based on 41.8 fb^{-1} data sample collected with the Belle detector. The results, $S_{\eta'K_S} = 0.28 \pm 0.55$ (stat) $^{+0.07}_{-0.08}$ (syst) and $A_{\eta'K_S} = 0.13 \pm 0.32$ (stat) $^{+0.09}_{-0.06}$ (syst), are the first measurements of CP asymmetry parameters related to ϕ_1 with a charmless B^0 decay. With more data, the uncertainty in $S_{\eta'K_S}$ can be reduced and impose tighter constraints on the phases from physics beyond the Standard Model.

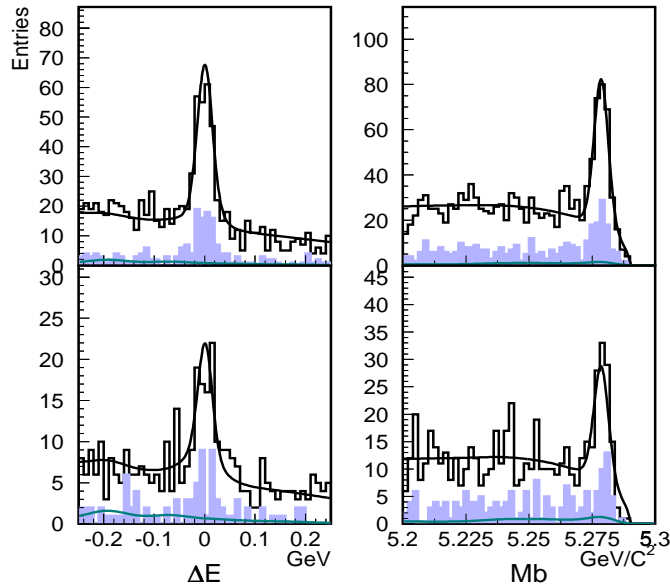


Figure 3: The M_{bc} and ΔE projection plots of $B^\pm \rightarrow \eta'K^\pm$ and $B^0 \rightarrow \eta'K_S$ modes.

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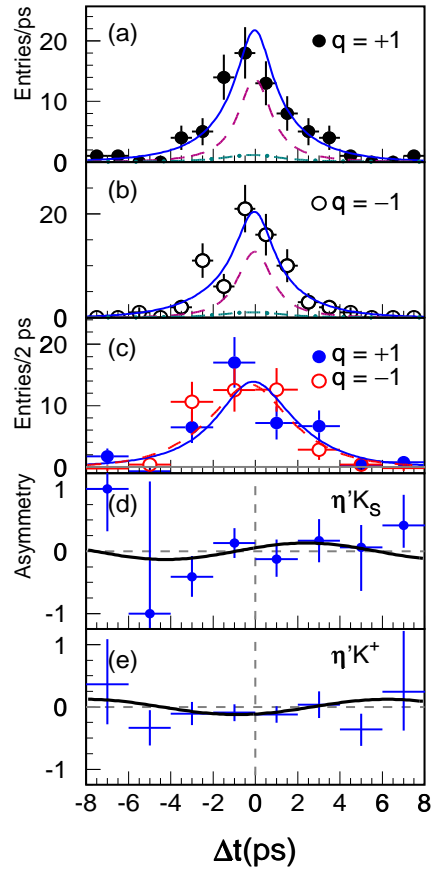


Figure 4: (a) The Δt distribution for B^0 tagged events. (b) The Δt distribution for \overline{B}^0 tagged events. (c) B^0 and \overline{B}^0 tagged events overlaid after background subtraction. (d) The asymmetries versus Δt plot. The solid curve is obtained from CP fit. (e) The asymmetries versus Δt plot for $\eta'K^\pm$.

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