

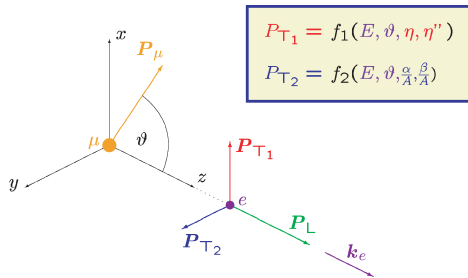
# Measurement of the Polarization Vector of the Positrons from the Decay of Polarized Muons

## as a Test of Time Reversal Invariance and a Quest for Physics Beyond the SM

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### Observables



$$P_{T1} = f_1(E, \theta, \eta, \eta'')$$

$$P_{T2} = f_2(E, \theta, \lambda, \lambda')$$

The standard model predicts:

$$\langle P_{T1} \rangle_E = 0.003$$

$$P_{T2} \equiv 0$$

A nonzero  $P_{T2}$  would signal violation of time reversal invariance. This is the only purely leptonic reaction for which TRI has been tested up to now.

### Matrix Element

$$\mathcal{M} = \frac{4G_F}{\sqrt{2}} \sum_{\gamma=S,V,T} \sum_{\epsilon,\mu=R,L} g_{\epsilon\mu}^{\gamma} \langle \bar{e}_\epsilon | \Gamma^\gamma | \nu_e \rangle \langle \bar{\nu}_\mu | \Gamma_\gamma | \mu \rangle$$

The index  $\gamma$  labels the type of interaction:

$$\Gamma^S = 4\text{-scalar}$$

$$\Gamma^V = 4\text{-vector}$$

$$\Gamma^T = 4\text{-tensor}$$

The indices  $\epsilon, \mu$  indicate the chiralities of the spinors of the observed (charged) leptons. The chiralities  $n, m$  of the neutrinos are uniquely determined for given  $\gamma, \epsilon$  and  $\mu$ .

The transverse polarization component  $P_{T1}$  yields the low energy parameter  $\eta$  without the suppression factor  $m_e/m_\mu$  of  $\eta$  in the energy spectrum of the decay positron. These interference terms allow for sizeable effects.

$$\eta = \frac{1}{2} \text{Re} \{ g_{LL}^V g_{RR}^{S*} + g_{RR}^V g_{LL}^{S*} + g_{LR}^V (g_{RL}^{S*} + g_{RL}^{T*}) + g_{RL}^V (g_{LR}^{S*} + g_{LR}^{T*}) \}$$

In the standard model:  $g_{LL}^V = 1$

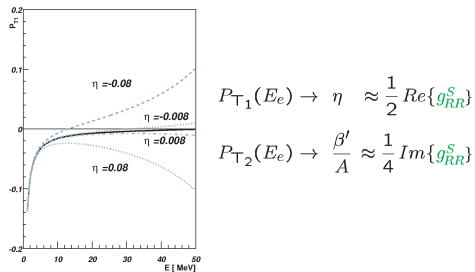
$$g_{\epsilon\mu}^T = 0$$

(all other interactions)

Assuming 1 additional interaction and knowing that

$$g_{LL}^V \approx 1,$$

one deduces:



$$P_{T1}(E_e) \rightarrow \eta \approx \frac{1}{2} \text{Re} \{ g_{RR}^S \}$$

$$P_{T2}(E_e) \rightarrow \frac{\beta'}{A} \approx \frac{1}{4} \text{Im} \{ g_{RR}^S \}$$

Main scientific interests:

$P_{T1}$ : Precise determination of Fermi coupling constant  $G_F$

$P_{T2}$ : Test of time reversal invariance

### Fermi Coupling Constant

Should be independent of masses and radiative corrections:

Universal coupling constant

$$G_F^2 = 192\pi^3 \cdot \frac{\hbar}{\tau_\mu} \cdot \frac{1}{m_\mu^5} \cdot \left\{ 1 + \frac{\alpha}{2\pi} \left( \pi^2 - \frac{25}{4} \right) \right\} \cdot \left\{ 1 - \frac{3}{5} \left( \frac{m_\mu}{m_W} \right)^2 \right\} \cdot \left\{ 1 - 4\eta \cdot \frac{m_e}{m_\mu} - 4\lambda \cdot \frac{m\nu_\mu}{m_\mu} + 8 \left( \frac{m_e}{m_\mu} \right)^2 + 8 \left( \frac{m\nu_\mu}{m_\mu} \right)^2 \right\}$$

New in  $G_F^2$ :  $\lambda \approx \frac{1}{2} \text{Re} \{ g_{LL}^V \cdot g_{LR}^{V*} \}$

In left-right symmetric models with mixing angle  $\zeta$ :

$$\lambda \approx \frac{1}{2} \zeta$$

Contribution from	$\frac{\Delta G_F}{G_F} [10^{-6}]$	
	$\mu^+ \rightarrow \bar{\nu}_\mu e^+ \nu_e$	$\tau^+ \rightarrow \bar{\nu}_\tau \mu^+ \nu_\mu$
$\Delta m_W$	0.0	1
$\Delta m_{\mu,\tau}$	0.2	421
$\Delta \tau$	9.1	2070
$\Delta(\lambda m_{\bar{\nu}})$	70.0	22500
$\Delta \eta$	193.0	6500
$\Delta \Gamma_{\tau \rightarrow \mu}$	-	100000

### Experimental Method

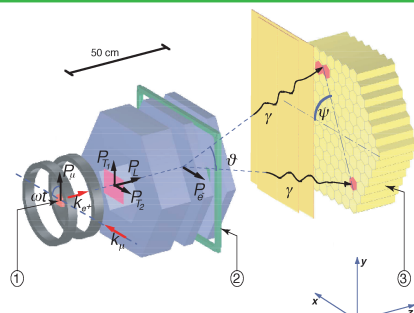
Measure the complete polarization vector of the decay positrons:

$$P_{e^+} = \begin{pmatrix} P_{T1} \\ P_{T2} \\ P_L \end{pmatrix} \equiv \begin{pmatrix} P_T \cdot \cos \varphi \\ P_T \cdot \sin \varphi \\ P_L \end{pmatrix}$$

with 3 simultaneous and independent measurements:

Observable	Method
$P_T$	Time dependence of annihilation
$\varphi$	Remnant $\mu$ SR effect
$P_L$	Spatial dependence of annihilation

Setup of the Experiment and Principle of Measurement:



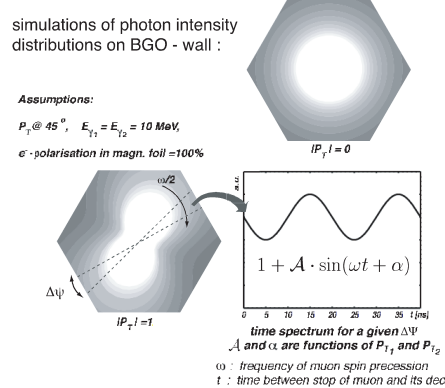
- 1: Beryllium stop target within spin precession magnet
- 2: magnetized Vacoflux foil within iron return yoke
- 3: calorimeter consisting of 127 BGO crystals

- Highly polarized  $\mu^+$  beam at  $\mu$ E1 area of PSI: (91%)
- Muon stop rate in Be target:  $(20 - 80) \times 10^6 \text{ s}^{-1}$
- Precession in homogeneous  $B$  field; precession frequency = cyclotron frequency (50.8 MHz)
- Burst width 3.9 ns (FWHM)  $\Rightarrow$  80% muon polarization in Be stop target
- Positron tracking with drift chambers
- Annihilation with polarized  $e^-$
- Detection of annihilation quanta with 127 BGO crystals

### Event Reconstruction and Data Analysis

- Energy calibration with cosmic rays
- Cluster recognition
- Background suppression
- Analysing power of annihilation events

•  $e^+$  annihilation-in-flight as analysing reaction for transverse polarization using the Dependence of the Annihilation Cross-Section on the Relative Orientation of Spins



Assumptions:

$$P_e @ 45^\circ, E_e = E_\mu = 10 \text{ MeV},$$

$e^-$  polarisation in magn. foil = 100%

$$IP_{e^+} = 0$$

$$IP_{e^-} = 1$$

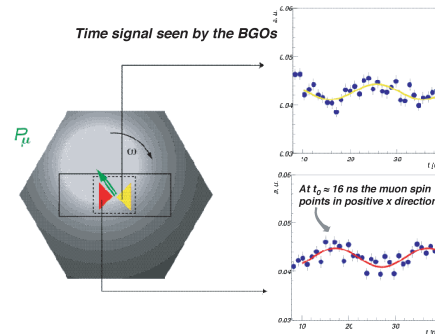
$$\Delta \Psi$$

$$1 + A \cdot \sin(\omega t + \alpha)$$

$$\omega$$
 : frequency of muon spin precession
   

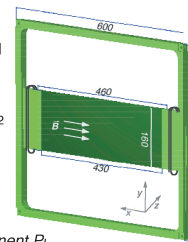
$$t$$
 : time between stop of muon and its decay

- Muon decay asymmetry ( $\mu$ SR) yields time zero  $\mu$ SR Effect is used to find the direction of the muon spin



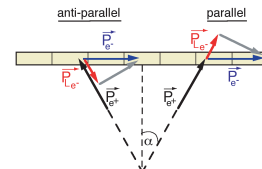
- Longitudinal polarization using information about position on magnetized Vacoflux foil (determined by tracks reconstructed from drift-chamber data) where annihilations take place

area on foil taken into account:  $140^2 \text{ mm}^2$  area divided into rectangular bins ( $ij$ ), 17 bins in  $x$ - and  $y$ -direction, respectively



Tracks that do not hit the center of the foil 'see' a longitudinal component  $P_{L,e^-}$  of the polarization of the electrons in the foil.

This  $P_{L,e^-}$  can either be parallel or anti-parallel to the positron polarization:

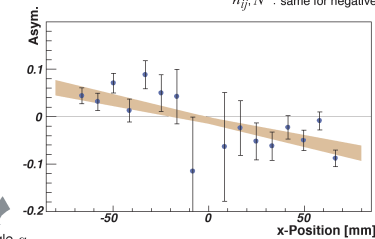


The annihilation cross section depends on the relative orientation of the spins of positrons and electrons in the foil. It is larger if both spins are anti-parallel.

$$\text{Asymmetry: } A_{ij} = \frac{n_{ij}^- - n_{ij}^+}{n_{ij}^- + n_{ij}^+} \quad \text{where } n_{ij}^- : \text{number of annihilations in bin } ij \text{ for positive foil polarization}$$

$$n_{ij}^+ : \text{total number of annihilations for positive polarization}$$

$$n_{ij}^- : \text{same for negative polarization}$$



- angle  $\alpha$
- electron polarization in foil ( $P_{e^-} = 7.2\%$ )
- analysing power of 0.79
- background factor of 0.75 (backgr. ratio 25%, mainly due to bremsstrahlung)

$$\langle P_L \rangle_E$$

### Preliminary Results

$P_T$  at the Time of Annihilation

Fitting the two perpendicular components  $P_1$  and  $P_2$  using a Log Likelihood parameter estimation

Use the diff. annihilation cross-section

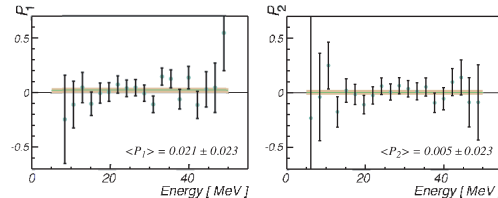
$$\frac{1}{\sigma_0} \frac{d\sigma}{d\Omega} = 1 + A \cdot \sin(\omega t + \alpha) =: f(P_1, P_2, E_{\nu_1}, E_{\nu_2}, \Psi, t)$$

where amplitude and phase are functions of the Energy and  $\Psi$

$$\mathcal{L}(P_1, P_2) := -\ln \prod_{i=1}^n f(P_1, P_2, E_{\nu_1}^i, E_{\nu_2}^i, \Psi^i, t^i)$$

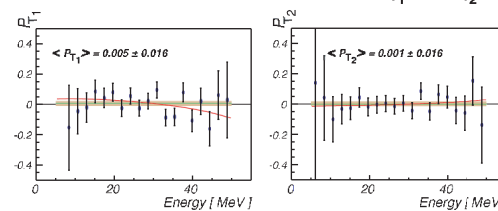
$$= -\sum_{i=1}^n \ln f^i(P_1, P_2) \quad \text{where } n \text{ is the number of 'good' annihilation events}$$

$\rightarrow$  for negative polarisation of the electrons in the magn. foil:



- $\rightarrow$  time zero from  $\mu$ SR Effect ( $\rightarrow$  orientation of muon spin relative to  $P_1$  and  $P_2$ )
- $\rightarrow$  rotation of transverse polarization components in the field of the spin precessing magnet (MC)
- $\rightarrow$  convolution with energy loss of the positron in the apparatus (MC)
- $\rightarrow$  sums of  $P_{T1}$  and  $P_{T2}$  for negative and positive foil polarization

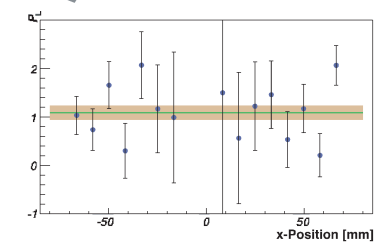
Transverse Polarization Components  $P_{T1}$  and  $P_{T2}$



### Longitudinal Polarization $P_L$

Asymmetry  $A_{ij}$

- angle  $\alpha$
- electron polarization  $P_{e^-}$
- analysing power
- background factor



$$\langle P_L \rangle_E = 1.09 \pm 0.15$$

### Summary of Preliminary Results and Implications

	General analysis	V - A + $g_{RR}^S$
$10^3 \times \langle P_{T1} \rangle_E$	$5 \pm 16$	$5 \pm 16$
$10^3 \times \langle P_{T2} \rangle_E$	$1 \pm 16$	$1 \pm 16$
$10^3 \times \eta$	$95 \pm 60$	$-4 \pm 14$
$10^3 \times \eta''$	$98 \pm 57$	$-10^3 \times \eta$
$10^3 \times \alpha'/A$	$-13 \pm 29$	0
$10^3 \times \beta'/A$	$8 \pm 16$	$1 \pm 7$
$10^3 \times \text{Re } g_{RR}^S$	-	$-8 \pm 28$
$10^3 \times \text{Im } g_{RR}^S$	-	$4 \pm 28$
$10^3 \times  g_{RR}^S $	-	$9 \pm 28$

Implications from these preliminary results:

1. No Evidence for Additional Scalar Couplings in Muon Decay
  - Green circle: Result of a general analysis including all possible left- and righthanded scalar, vector and tensor couplings.
  - Red circle: Only one additional coupling interferes with the lefthanded vector coupling in the SM.
  - $\eta = -0.004 \pm 0.014$
  - $\beta'/A = 0.001 \pm 0.007$
  - $\text{Im}(g_{RR}^S) = 0.004 \pm 0.028$
  - $\text{Re}(g_{RR}^S) = -0.008 \pm 0.028$
  - $|g_{RR}^S| = 0.009 \pm 0.028$

2. Model Independent Measurement of  $G_F$

The model independent Fermi coupling constant is:

$$G_F = \sqrt{192\pi^3 \frac{1}{\tau_\mu m_\mu^5} \left( 1 - 4\eta \frac{m_e}{m_\mu} \right)}$$

using the measured value of  $\eta = 0.095 \pm 0.061$  yields:

$$G_F = 1.16532 \pm 0.00069 \times 10^{-5} (\hbar c)^3 \text{ GeV}^{-2}$$

1) For illustrational purposes radiative corrections are not given.

### Outlook

Improve precision of previous experiment [1] by almost one order in magnitude to:

$$\Delta \langle P_{T1} \rangle = 0.004$$

$$\Delta \langle P_{T2} \rangle = 0.004$$

Assuming V - A and one additional coupling, this will reduce the limits for  $\eta$  and  $g_{RR}^S$  to:

$$\Delta \eta = 0.004$$

$$\Delta \text{Re} \{ g_{RR}^S \} = 0.009$$

$$\Delta \text{Im} \{ g_{RR}^S \} = 0.009$$

[1] H. Burkard et al., Phys. Lett. **160 B** (1985) 343

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