

**Higgs and SUSY Particle Predictions
from SO(10) Yukawa Unification**

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Outline

- Introduction to SUSY GUTs

Gauge hierarchy problem

EWSB \iff Heavy top quark

Charge quantization

Gauge coupling unification

- SO(10) Yukawa unification
- χ^2 analysis and predictions

Yukawa unification is only consistent in a narrow region of SUSY parameter space with significant consequences for the Higgs and SUSY spectrum

- How robust are predictions?
- Conclusion

SUSY GUTs

- Charge quantization

— Georgi & Glashow; Pati & Salam; Georgi, Fritzsche & Minkowski

$$\text{SU}_5 \quad \left\{ Q = \begin{pmatrix} u \\ d \end{pmatrix} \bar{e} \bar{u} \right\} \subset 10$$

$$\left\{ \bar{d} \ L = \begin{pmatrix} \nu \\ e \end{pmatrix} \right\} \subset \bar{5}$$

$$H_u, H_d \subset 5_H, \bar{5}_H$$

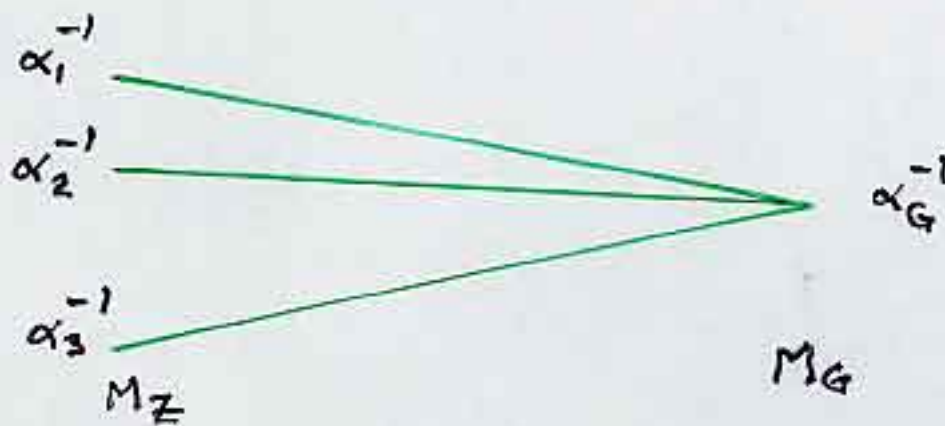
$$\text{SO}_{10} \quad 10 + \bar{5} + \bar{\nu}_{sterile} \subset 16$$

$$5_H, \bar{5}_H \subset 10_H$$

Gauge coupling unification *

— Dimopoulos, S.R. & Wilczek; Dimopoulos & Georgi

* Only evidence for SUSY



- Significant GUT threshold corrections from Higgs and GUT breaking sectors

Def: $M_G \iff \alpha_1(M_G) = \alpha_2(M_G) \equiv \alpha_G$

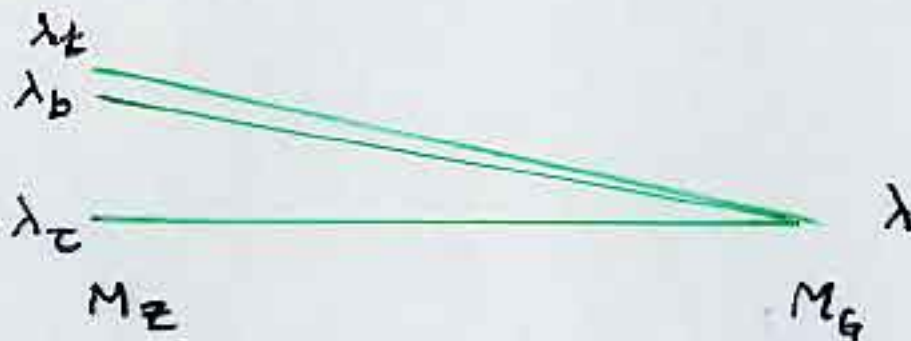
Good fit requires: $\epsilon_3 \equiv \frac{(\alpha_3(M_G) - \alpha_G)}{\alpha_G} \sim -4\%$



SO(10) Yukawa unification

$$SO_{10} \quad \lambda_b = \lambda_t = \lambda_\tau = \lambda_{\nu_\tau} = \lambda$$

— Banks; Olechowski & Pokorski; Ananthanarayan, Lazarides & Shafi; Dimopoulos, Hall & S.R.



$$\lambda \Leftrightarrow m_b/m_\tau \quad m_\tau = \lambda_\tau \frac{v}{\sqrt{2}} \cos\beta \implies \tan\beta \approx 50$$

$$m_t = \lambda_t \frac{v}{\sqrt{2}} \sin\beta \sim 170 \pm 20 \text{ GeV}$$



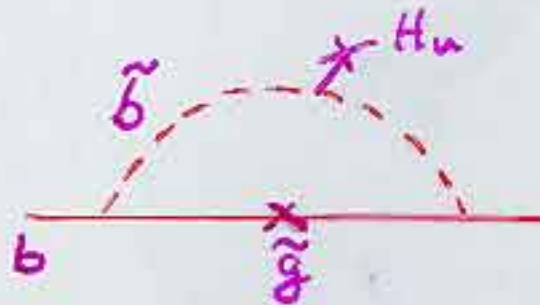
- Insignificant GUT threshold corrections from gauge and Higgs loop

Weak scale threshold corrections $\propto \tan\beta$

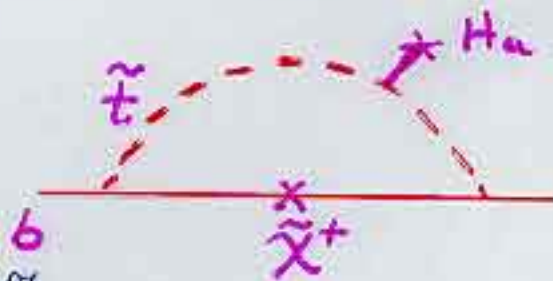
— Hall, Rattazzi, & Sarid; Carena, Olechowski, Pokorski & Wagner; Blazek, Pokorski & S.R.; Pierce, Bagger, Matchev & Zhang

$$\delta m_b / m_b = \Delta m_b^{\tilde{g}} + \Delta m_b^{\tilde{\chi}^+} + \Delta m_b^{\log} + \dots$$

$$\Delta m_b^{\tilde{g}} \approx \frac{2\alpha_3}{3\pi} \frac{\mu m_{\tilde{g}}}{m_b^2} \tan\beta$$



$$\Delta m_b^{\tilde{\chi}^+} \approx \frac{\lambda_t^2}{16\pi^2} \frac{\mu A_t}{m_t^2} \tan\beta$$



$$\Delta m_b^{\log} \approx \frac{\alpha_3}{4\pi} \log\left(\frac{\tilde{m}^2}{M_Z^2}\right) \sim 6\%$$



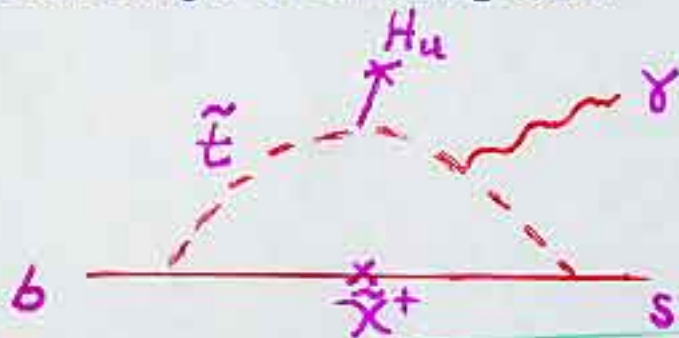
- $\Delta m_b^{\tilde{g}} \sim -\Delta m_b^{\tilde{\chi}^+} > 0$ for $\mu > 0$ [our conventions]

Can be $\sim 50\%$

- Good fits require $\delta m_b / m_b < -2\%$

Data favors $\mu > 0$

$b \rightarrow s\gamma$ — chargino term typically dominates and has opposite sign to the SM and charged Higgs contributions for $\mu > 0$, thus reducing the branching ratio.



- $\mu < 0$ is problematic.

$$\frac{\lambda_t^2}{16\pi^2} \frac{\mu A_t}{m_{\tilde{t}}^4} m_b \bar{b} \not{F}_{\mu\nu} \sigma_{\mu\nu} s \tan\beta$$

$$\frac{g_{\mu-2}}{2} \equiv a_{\mu}^{NEW} = \frac{25}{15} (16) \times 10^{-10}$$

— Muon $g-2$ Collaboration and Brookhaven E821 Collaboration 2000

Sign of a_{μ}^{NEW} correlated with sign of μ

— Chattopadhyay & Nath

- Favors $\mu > 0$.

χ^2 Analysis

- 11 input parameters at M_G

$$\begin{aligned} & \lambda, \alpha_G, M_G, \epsilon_3 \\ & m_{10}, A_0, \tan \beta(M_Z) \\ & D_X \text{ [D term splitting]} \\ & \text{(or } \Delta m_H^2 \text{ [} \text{Universal scalar masses} \text{])} \\ & \text{Just so splitting} \\ & m_{16}, \mu, M_{1/2} \text{ — fixed} \end{aligned}$$

- 9 observables — $X_i^{exp} (\sigma_i)$

$$\begin{aligned} & G_\mu, \alpha, \alpha_s(M_Z) = 0.118 (0.002) \\ & \rho^{NEW}, M_Z, M_W \\ & M_t = 174.3 (5.1) \\ & m_b(m_b) = 4.20 (0.20), M_\tau \end{aligned}$$

$$\chi^2 = \sum_{i=1}^9 \frac{(X_i^{exp} - X_i^{theory})^2}{\sigma_i^2}$$

Bottom Line

- Yukawa unification possible only in a narrow region of SUSY parameter space
- Since for $\mu > 0$ need $|\Delta m_b^{\tilde{\chi}}| > \Delta m_b^{\tilde{g}} \rightarrow \sim 6\%$

$$\Delta m_b^{\tilde{g}} \approx \frac{2\alpha_3}{3\pi} \frac{\mu m_{\tilde{g}}}{m_b^2} \tan\beta$$
$$\Delta m_b^{\tilde{\chi}^+} \approx \frac{\lambda_t^2}{16\pi^2} \frac{\mu A_t}{m_t^2} \tan\beta$$

- Implies

$$A_t \ll 0 \iff A_0 \ll 0$$

- Also
$$\begin{pmatrix} m_t^2 & m_t A_t \\ m_t A_t & m_{\tilde{t}}^2 \end{pmatrix}$$

$$m_{\tilde{t}} \ll m_{\tilde{b}}$$

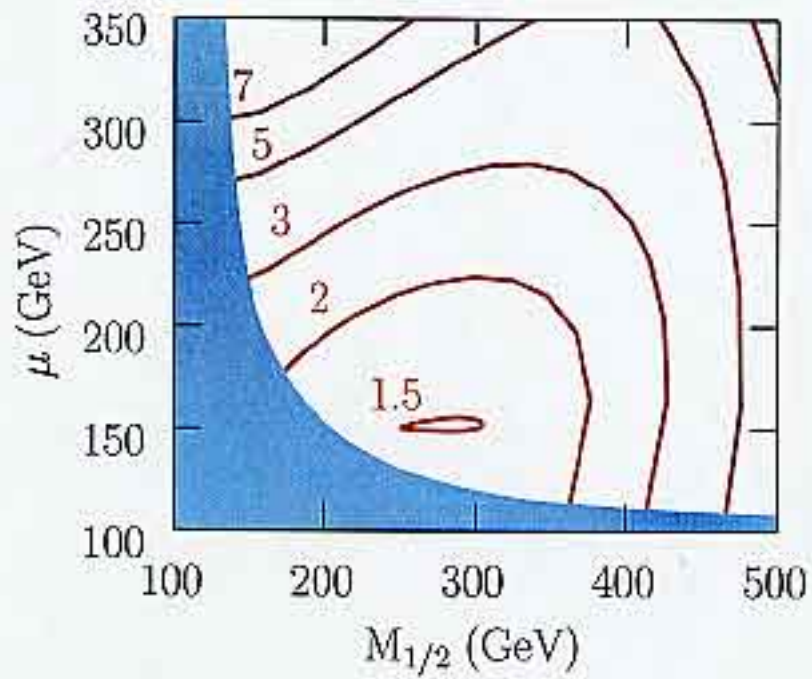
- Good fits require

$$A_0 \sim -1.9 m_{16}$$
$$m_{10} \sim 1.35 m_{16}$$
$$m_{16} > 1500 \text{ GeV}$$

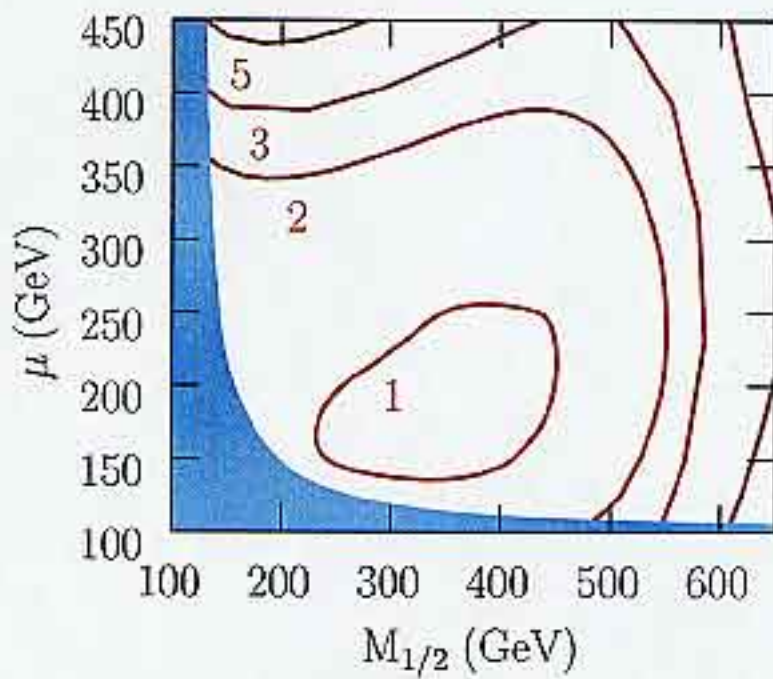
Just So Higgs splitting

- Higgs splitting with ~~Universal scalar masses~~

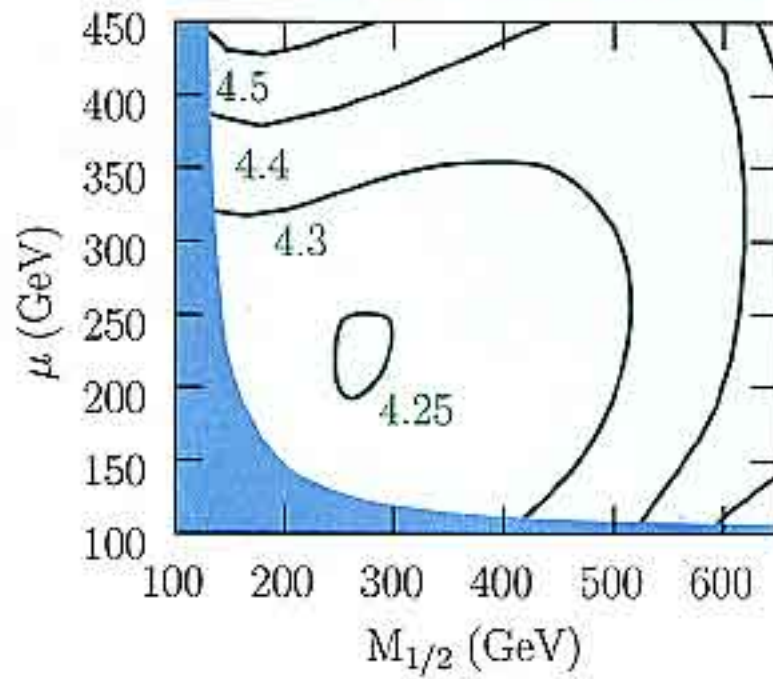
χ^2 for $m_{16} = 1.5$ TeV



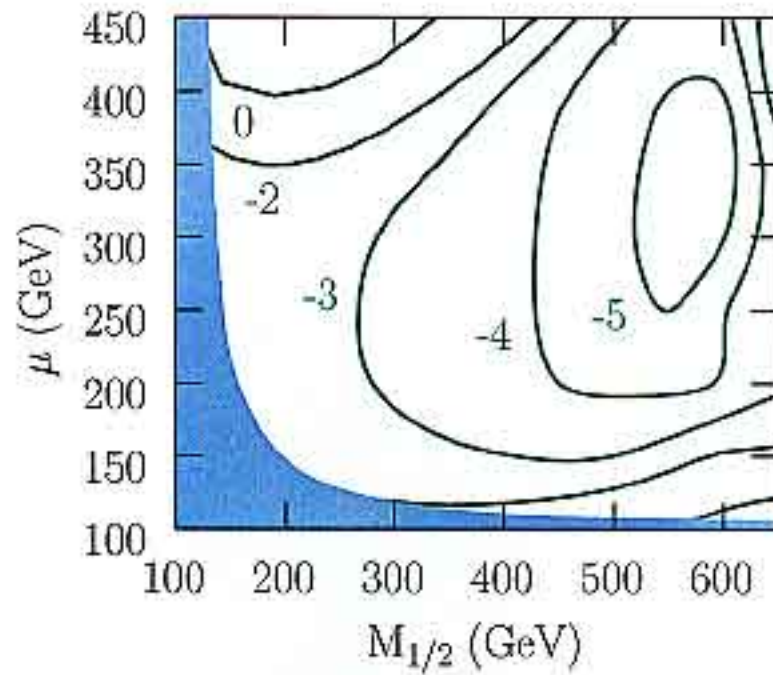
χ^2 for $m_{16} = 2$ TeV



m_b (GeV), $m_{16} = 2$ TeV

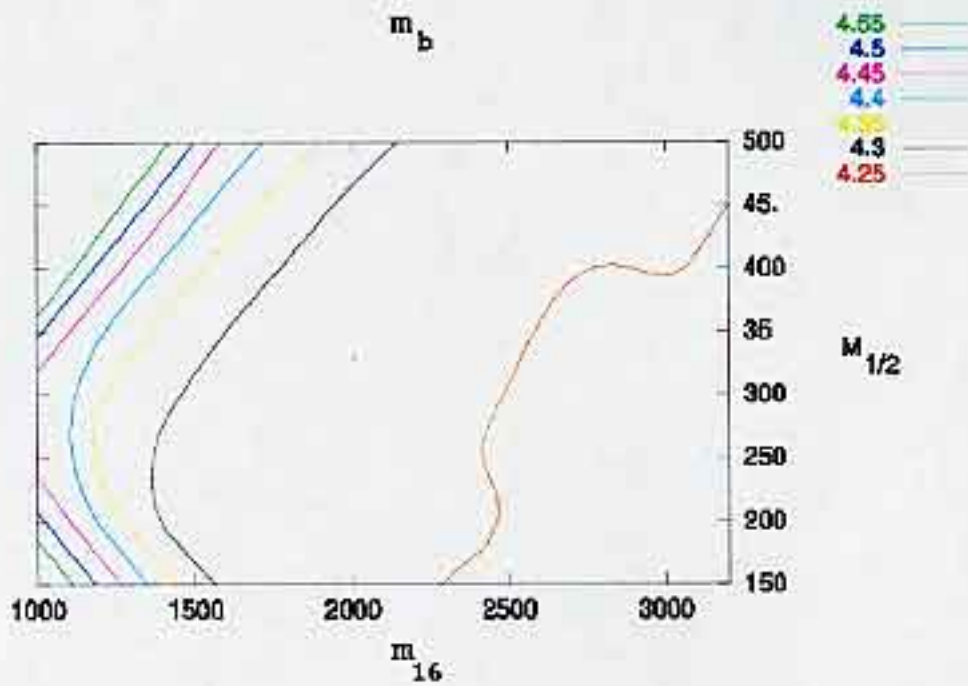
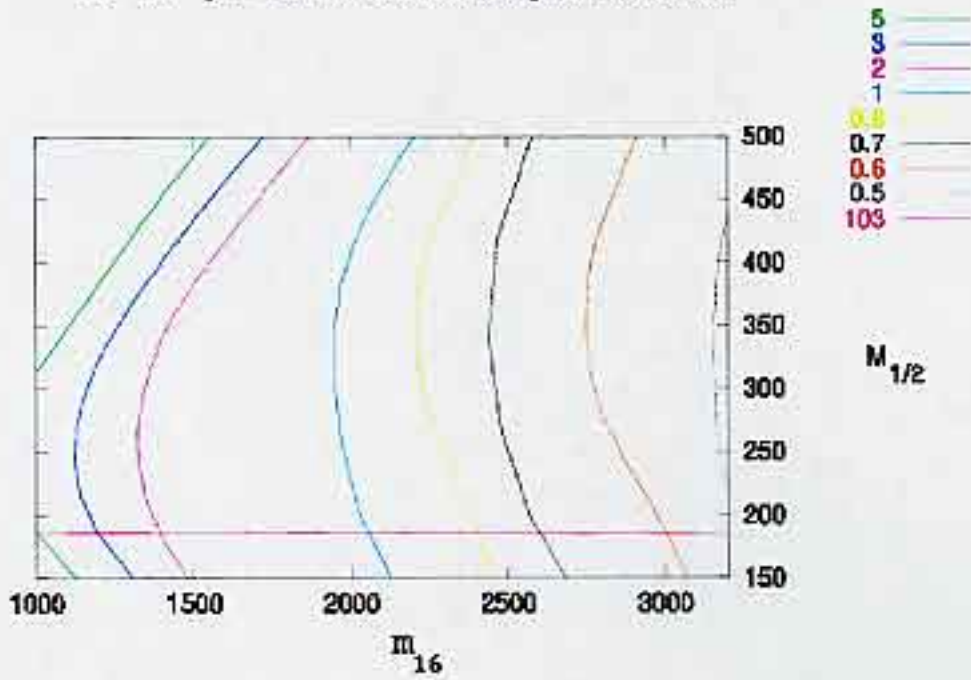


SUSY corr. to m_b (%)



$m_{16} \uparrow \quad \chi^2 \downarrow$

χ^2 for $\mu = 150$ GeV fixed with charging 103 GeV line



Results — SUSY spectra for $\chi^2 \leq 1$

$$\tilde{t}_1 \sim (175 - 250) \text{ GeV}$$

$$\tilde{b}_1 \sim (500 - 650) \text{ GeV}$$

$$\tilde{\tau}_1 \sim (250 - 500) \text{ GeV}$$

$$\tilde{g} \sim (600 - 1200) \text{ GeV}$$

$$\tilde{\chi}^+ \sim (120 - 240) \text{ GeV}$$

$$\tilde{\chi}^0 \sim (75 - 160) \text{ GeV}$$

} $b \rightarrow s \gamma$?

] light

- $\tilde{\chi}^+$ visible at Tevatron? Needs new analysis !!

New benchmark pt.

- What about $\tilde{t}_1, \tilde{\tau}_1$??

Constrained by $b \rightarrow s \gamma$

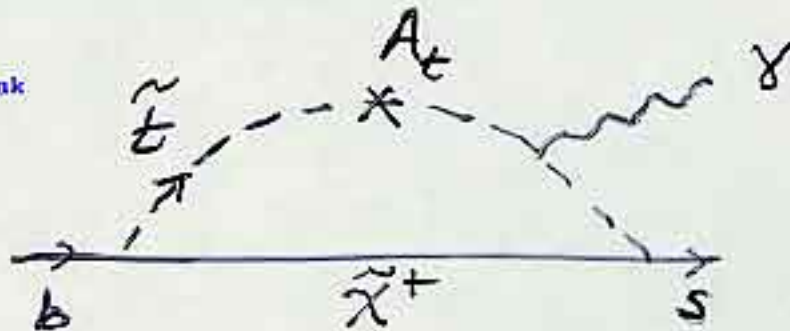
- What about Higgs — $\bar{A}_0, h_0, H_0, H_{\pm}$??

Constrained by $B_s \rightarrow \mu^+ \mu^-$

Two additional experimental constraints

$B \rightarrow s\gamma$

Gambino and Misiak

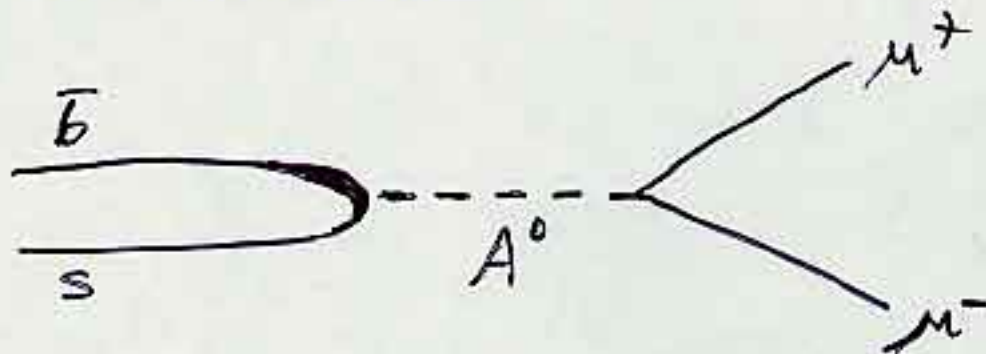


$\sim \lambda_t^2 \frac{\mu A_t \tan \beta}{(m_{\tilde{t}_1})^4} \Rightarrow m_{\tilde{t}_1} \uparrow$

$B_s \rightarrow \mu^+ \mu^-$

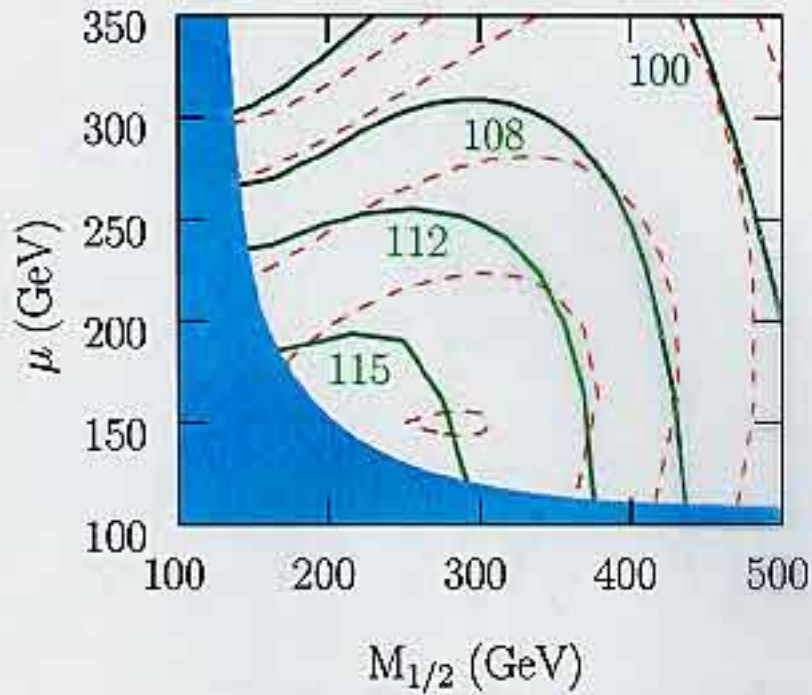
Babu and Kolda; Chankowski and Slawianowska; Bobeth et al.;

Dedes, Dreiner and Nierste; Isidori and Retico

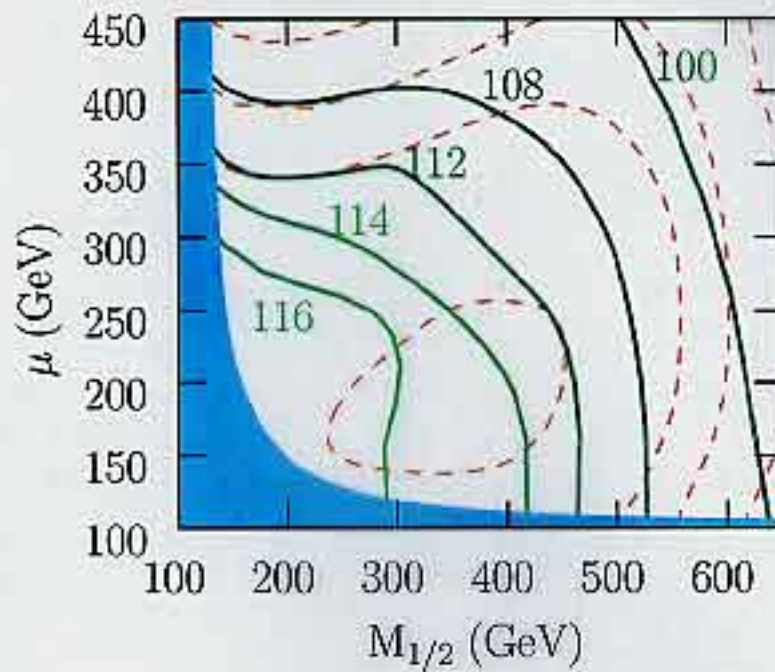


$m_{A^0} \gtrsim 200 \text{ GeV}$

m_{h^0} (GeV), $m_{16} = 1.5\text{TeV}$

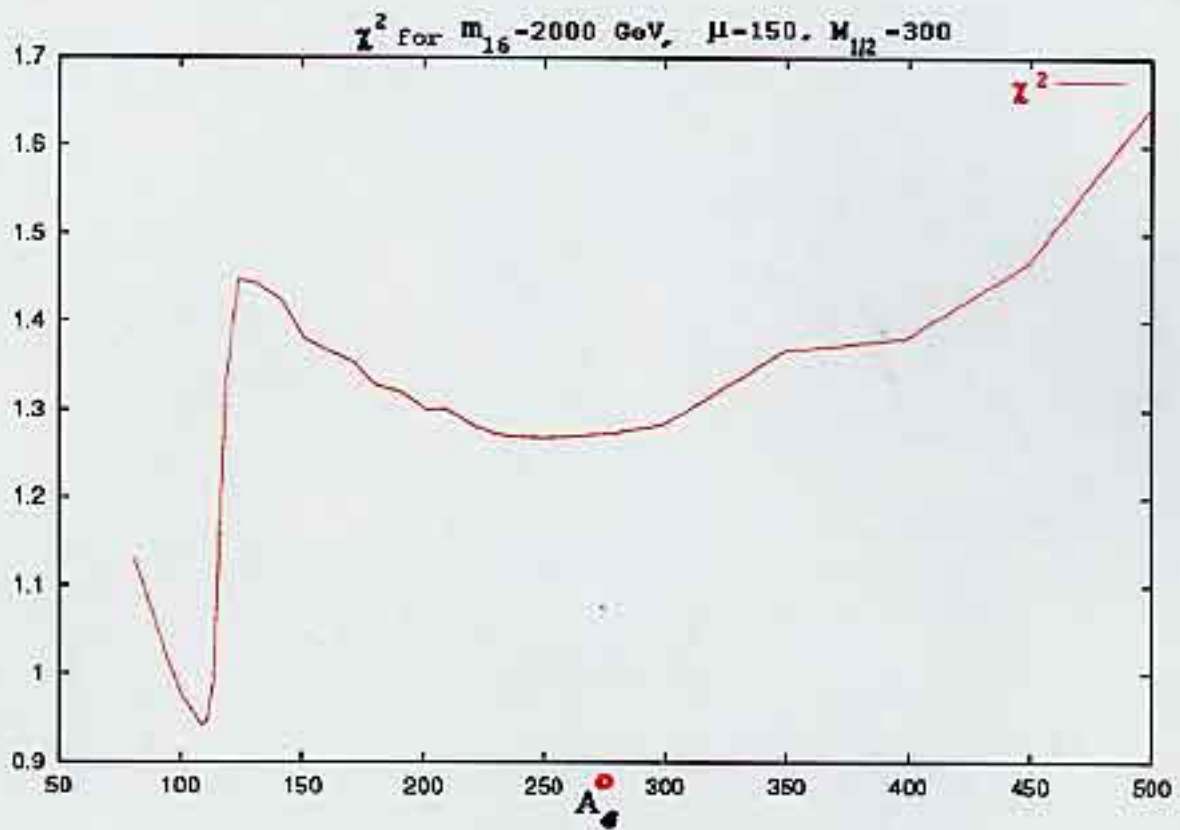


m_{h^0} (GeV), $m_{16} = 2\text{TeV}$

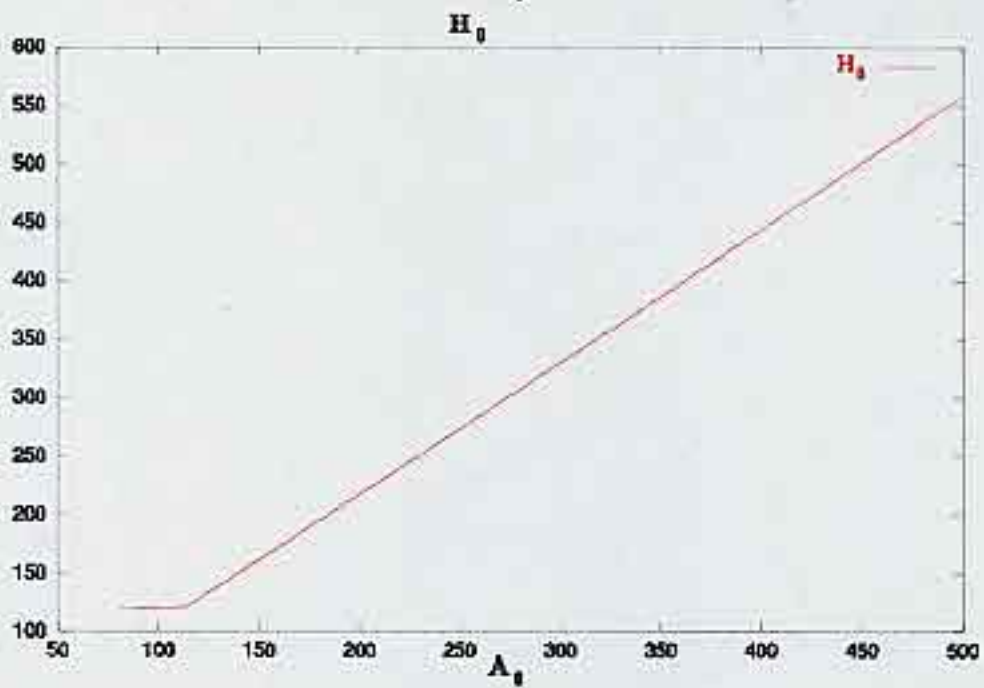
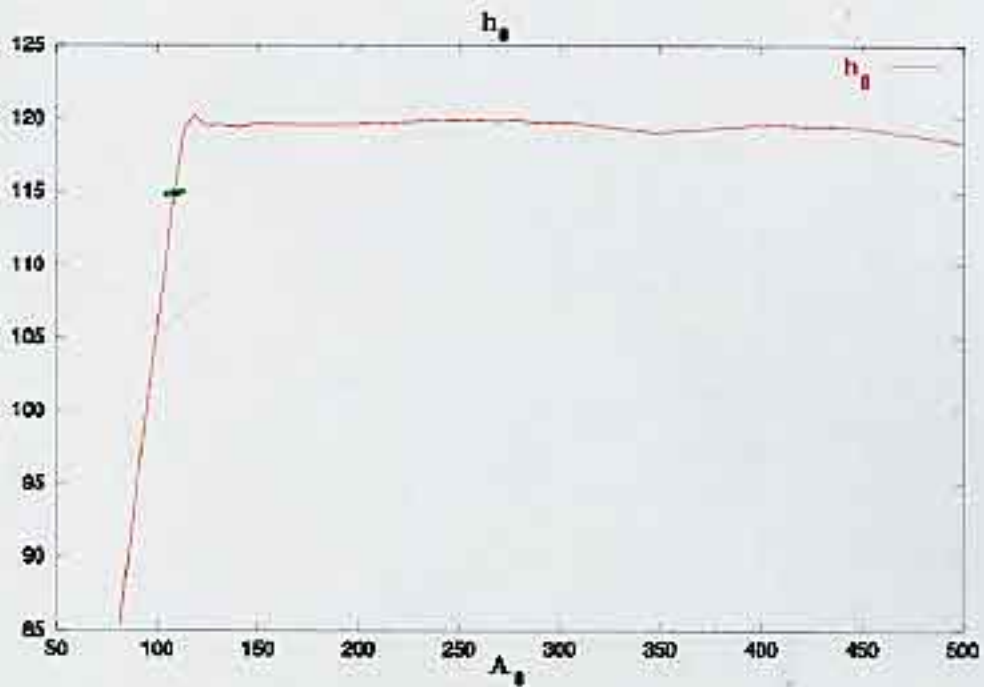


$m_{h^0} \approx 115 \pm 3 \pm 3$
 $\chi^2 \leq 1$ theor.

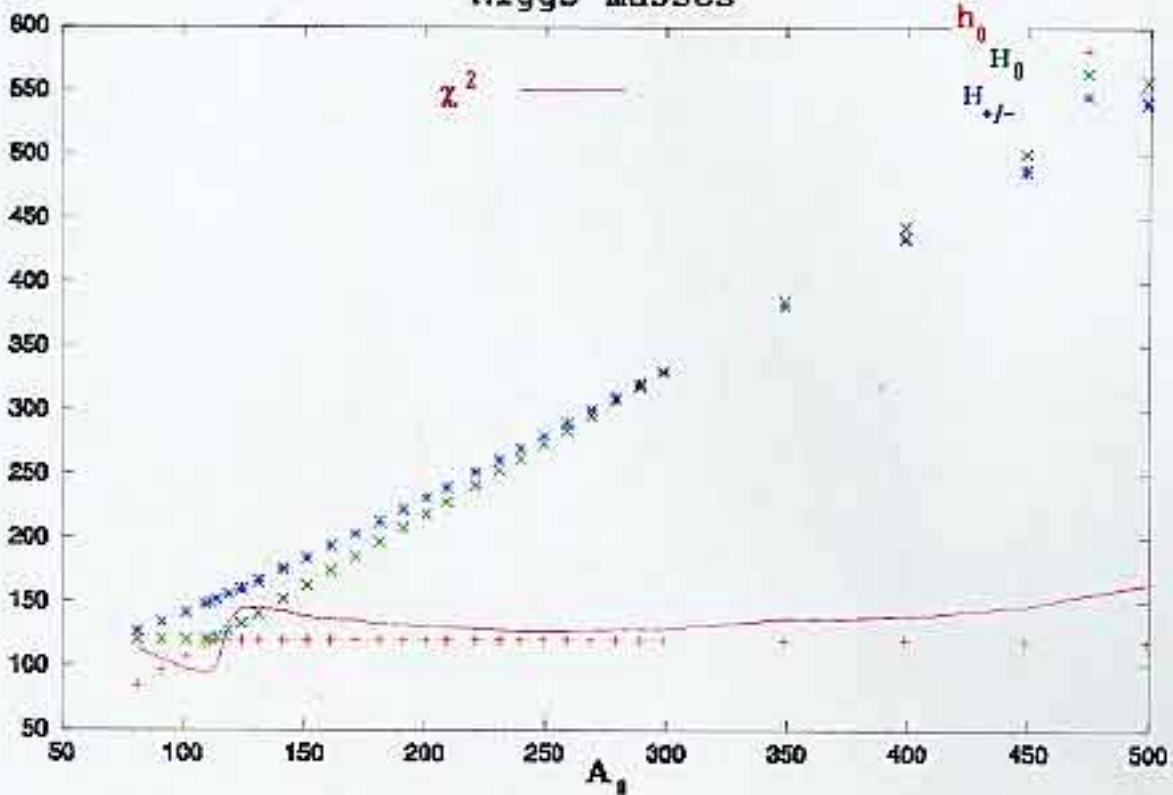
Sensitivity of χ^2 to A_0 mass



Sensitivity of Higgs masses



Higgs masses



Conclusion

• $SO(10)$ Yukawa unification makes definite predictions for Higgs and Sparticle spectrum

- $\tilde{\chi}_1^\pm \sim 120 - 240 \text{ GeV}$ Chargino visible at Tevatron ?
- $\tilde{\chi}_1^0 \sim 75 - 160 \text{ GeV}$ LSP
- $(m_{\tilde{t}_i})_{MIN} \sim 450 \text{ GeV}$ and $m_{\tilde{t}_1} \ll m_{\tilde{b}_1}$
- $m_h \sim 114 \pm 5 \pm 3 \text{ GeV}$

• Soft SUSY breaking parameters satisfy \star

- $A_0 \sim -1.9 m_{16}$,
- $m_{10} \sim 1.35 m_{16}$,
- $m_{16} > 1200 \text{ GeV}$

$$\Rightarrow \alpha_\mu^{SUSY} < 16 \times 10^{-10}$$

\star See also Bagger, Feng, Polonsky, and Zhang, *Phys. Lett. B* **473**, 264 (2000)

(Heavy 1st & 2nd; Light 3rd generation scalars : avoids large flavor violation)

and H. Baer and J. Ferrandis, *Phys. Rev. Lett.* **87**, 211803 (2001)

• We should know soon!

EWSB with large $\tan\beta$ needs

$$m_{H_u}^2 < m_{H_d}^2$$

- D_X term vs. Just So splitting

D_X splitting — $SO(10) \rightarrow SU(5) \times U(1)_X$

$$m_{(H_u, H_d)}^2 = m_{10}^2 \mp 2D_X$$

$$m_{(Q, \bar{u}, \bar{e})}^2 = m_{16}^2 + D_X$$

$$m_{(\bar{d}, L)}^2 = m_{16}^2 - 3D_X$$

$$\Rightarrow m_{\bar{t}}^2 \leq m_{\bar{b}}^2 \quad \text{Bad}$$

Just So splitting —

$$m_{(H_u, H_d)}^2 = m_{10}^2 (1 \mp \Delta m_H^2)$$

$$m_{(Q, \bar{u}, \bar{e})}^2 = m_{16}^2$$

$$m_{(\bar{d}, L)}^2 = m_{16}^2$$

$$\Rightarrow m_{\bar{t}}^2 \ll m_{\bar{b}}^2 \quad \text{Good}$$

- Higgs special

μ term — $\mu H_u H_d$

Doublet-Triplet splitting

ν_τ contribution to Higgs splitting

Yukawa term :

$$\lambda_{\nu_\tau} \bar{\nu}_\tau L H_u + \frac{1}{2} \bar{\nu}_\tau M_{\nu_\tau} \nu_\tau$$

$$\lambda_{\nu_\tau} = \lambda_t = \lambda_b = \lambda_\tau \equiv \lambda$$

- ν_τ contribution distinguishes H_u and H_d
- $\Delta m_{H_u}^2 \approx \frac{\lambda^2}{16\pi^2} (2m_{16}^2 + m_{10}^2 + A_0^2) \log\left(\frac{M_{\nu_\tau}^2}{M_G^2}\right) + \dots$

Taking typical GUT values:

$$\lambda = 0.7,$$

$$M_{\nu_\tau} = 10^{13} \text{ GeV}, M_G = 3 \times 10^{16} \text{ GeV} (\implies (\Delta m_\nu^2)_{\text{atm}})$$

$$A_0^2 \approx 2m_{10}^2 \approx 4m_{16}^2$$

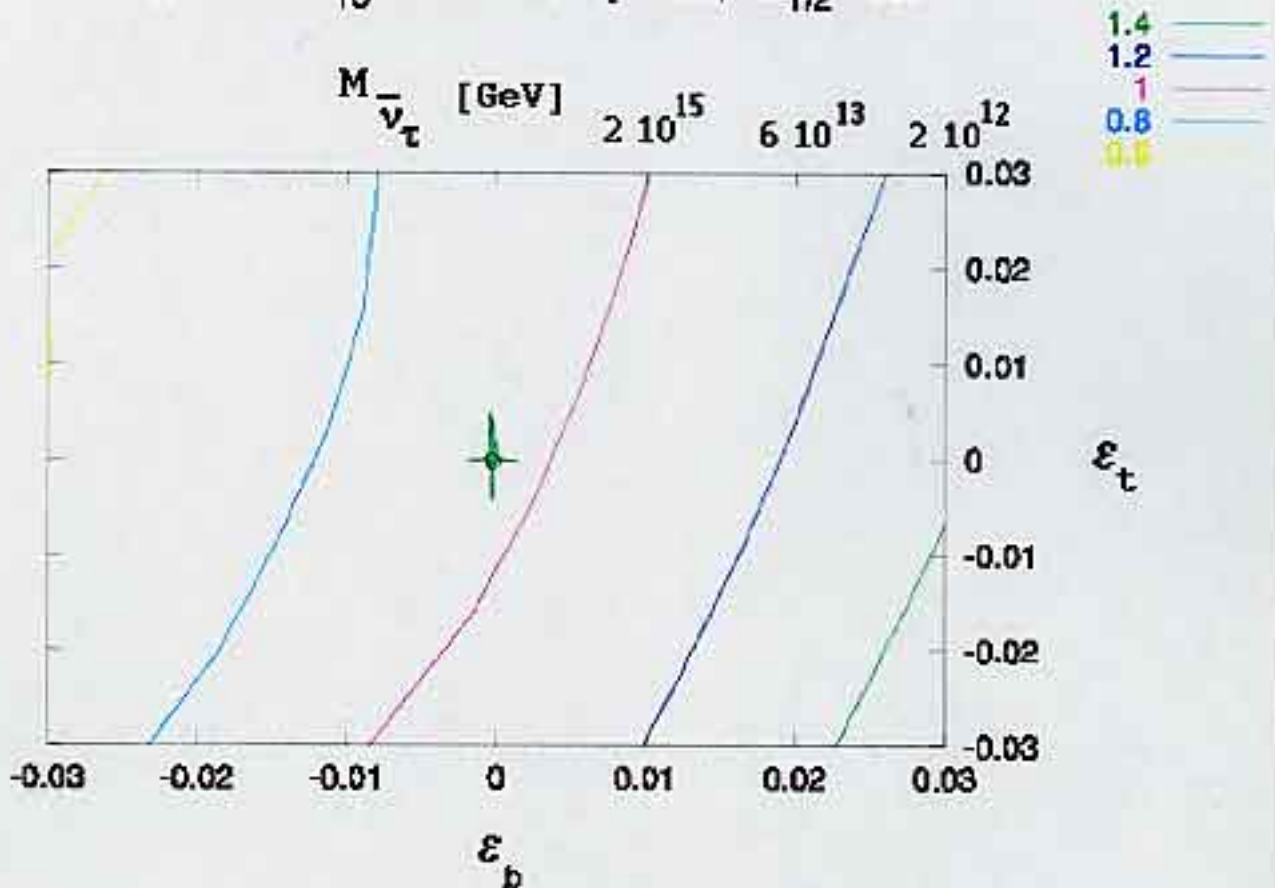
$$\Delta m_H^2 \equiv \frac{1}{2} \Delta m_{H_u}^2 / m_{10}^2 = .10$$

- Just So splitting of just the right size

Just So Higgs splitting

~~Chained at Scales~~

χ^2 for $m_{16} = 2000 \text{ GeV}$, $\mu = 150$, $M_{1/2} = 300$



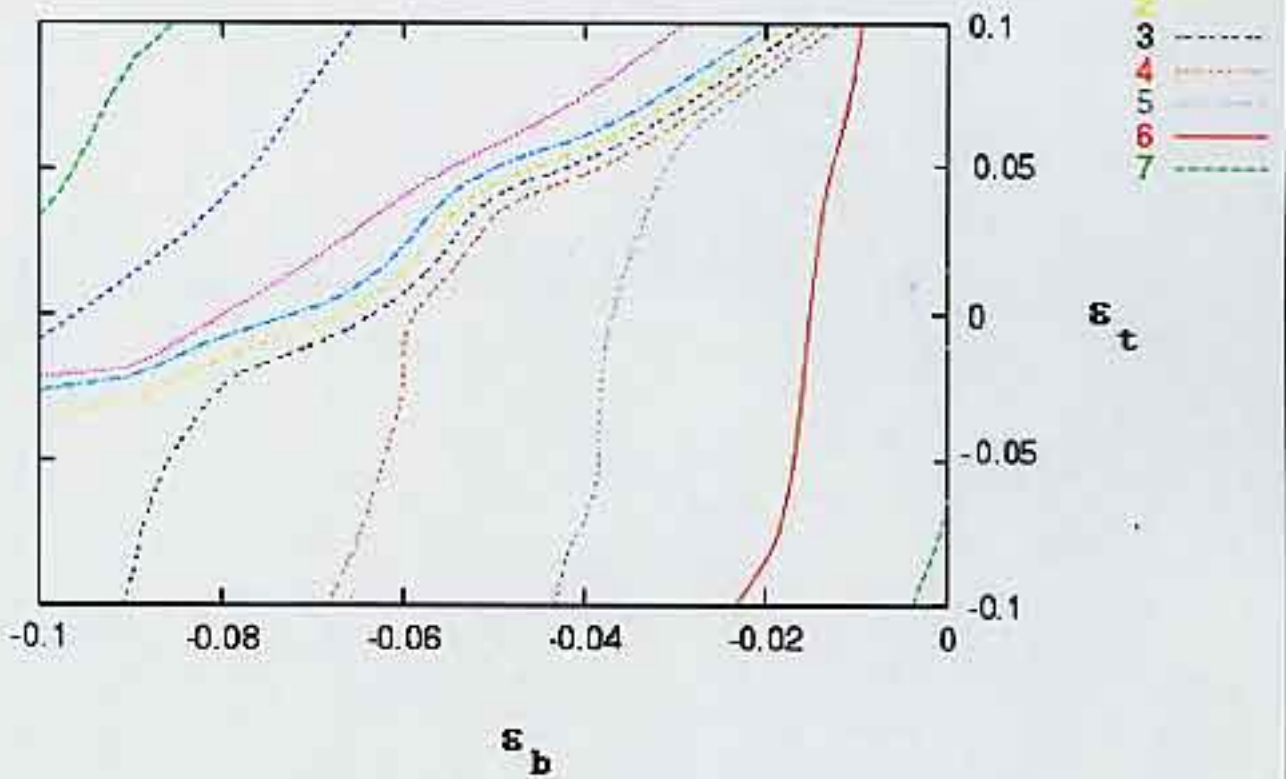
$$\lambda_\tau = \lambda$$

$$\lambda_i = \lambda (1 + \epsilon_i), \quad i = b, t$$

D_x term splitting

χ^2 for D terms

$m_{16} = 2000 \text{ GeV}, \mu = 150, M_{1/2} = 300$



Data points		with $B(b \rightarrow s\gamma)$ and $m_{A^0} = 200$ GeV	without
Input parameters			
α_G^{-1}		24.72	24.66
$M_G \times 10^{-16}$		3.00	3.07
ϵ_3		-0.040	-0.040
L		0.63	0.67
m_{16}		2000	2000
m_{10}/m_{16}		1.32	1.35
Δm_H^2		0.13	0.13
$M_{1/2}$		350	350
μ		200	200
$\tan \beta$		52.5	50.5
A_0/m_{16}		-1.71	-1.87
χ^2 observables	Exp (σ)		
M_Z	91.188 (0.091)	91.18	91.14
M_W	80.419 (0.080)	80.42	80.45
$G_\mu \times 10^5$	1.1664 (0.0012)	1.166	1.166
α_{EM}^{-1}	137.04 (0.14)	137.0	137.0
$\alpha_s(M_Z)$	0.118 (0.002)	0.1172	0.1176
$\rho_{new} \times 10^3$	-0.200 (1.10)	0.228	0.460
M_t	174.3 (5.1)	173.8	174.6
$m_b(m_b)$	4.20 (0.20)	4.46	4.27
M_τ	1.7770 (0.0018)	1.777	1.777
$B(b \rightarrow s\gamma) \times 10^3$	0.296 (0.035)	0.297	NA
TOTAL χ^2		2.1	0.9
h^0		118	116
H^0		217	121
A^0		200	110
H^\pm		229	148
$\tilde{\chi}_1^0$		130	130
$\tilde{\chi}_2^0$		190	190
$\tilde{\chi}_1^\pm$		178	178
\tilde{g}		909	913
\tilde{t}_1		509 (-30)	222
\tilde{b}_1		749 (-42)	588
$\tilde{\tau}_1$		459 (+47)	420
$\alpha_\mu^{SUSY} \times 10^{10}$	25.6 (16)	5.8	5.5