

Determine
The Running of
 α QED
from
Small Angle Bhabha
Scattering

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● The Problem of Luminosity at LEP

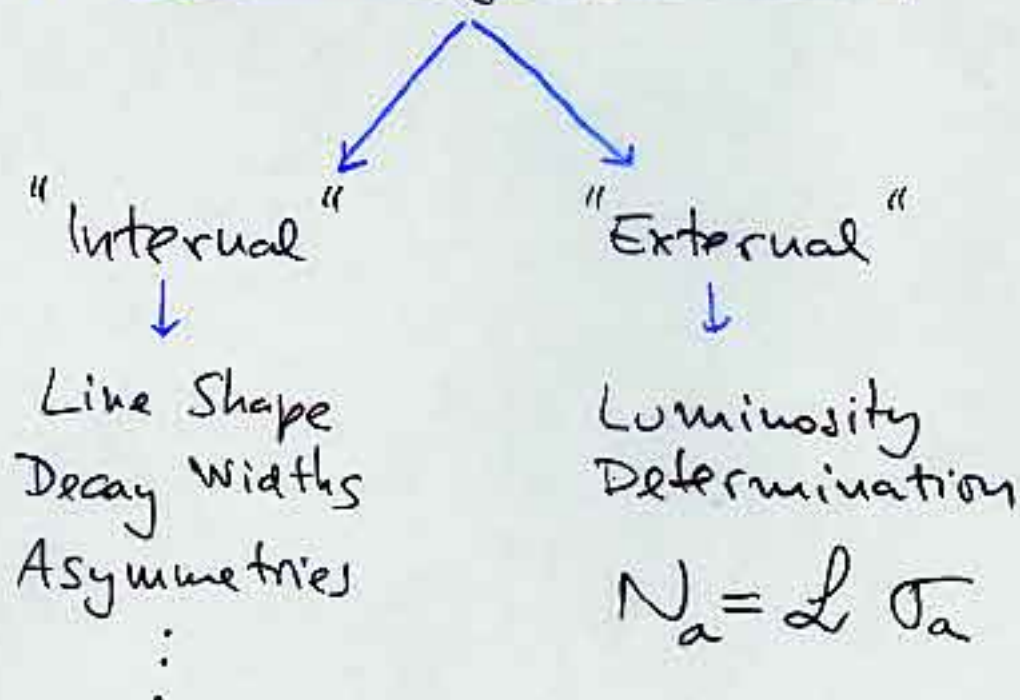
1989 Expected Experimental Accuracy
on the Luminosity $\frac{\Delta L}{L} \approx 2 \div 4\%$

1994 $\frac{\Delta L}{L} \approx 0.08 \div 0.1\%$

Dominated by the theoretical error

$$\sqrt{s} = M_{Z^0} \quad L = \frac{N_{\text{Bhabha}}}{\sigma_{\text{Bhabha}}^{\text{Th}}} \leftarrow \begin{array}{l} 0.09\% \\ 0.25\%? \end{array}$$

● Precision Physics at LEP



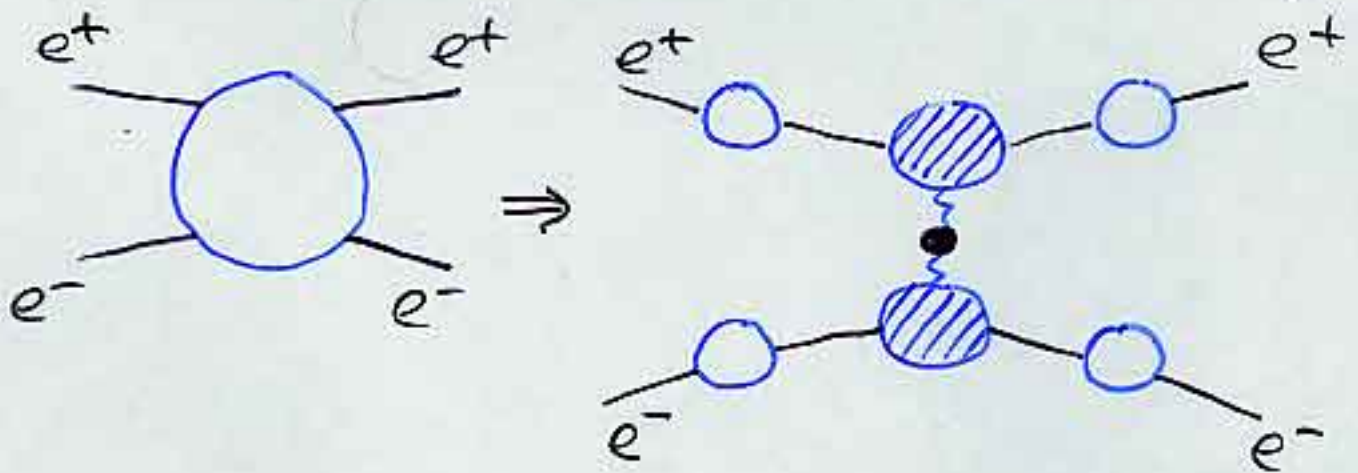
Luminosity \leftrightarrow Small Angle Bhabha Scattering

LEP

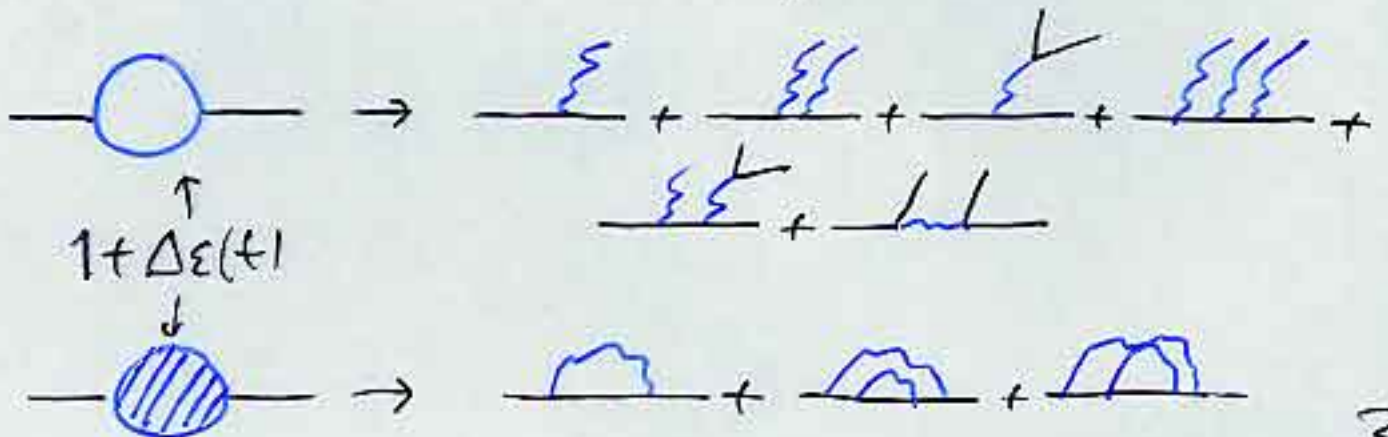
1994 \rightarrow Increasing Accuracy to evaluate Radiative Corrections

- A Next-to-Next-to-Leading accurate determination of the Small Angle Bhabha Cross Section

$$\sigma^{\text{Bhabha}} = \sigma^{(0)} \left(1 + \sigma_{\theta, W} + \sigma^{\delta} + \sigma^{2\delta} + \sigma^{e^+e^-} + \sigma^{3\delta} + \sigma^{e^+e^-} \right)$$



Photon vacuum polarization



$1 + \Delta\epsilon(t)$

1995 → Theoretical Accuracy $\approx 0.11\%$
 Small Angle Bhabha

$$\frac{d\sigma}{dt} = \frac{d\sigma^{(0)}}{dt} \left(\frac{\alpha(t)}{\alpha(0)} \right)^2 (1 + \Delta\epsilon(t))$$

$$\alpha(q^2) = \frac{\alpha(0)}{1 - \Delta\alpha(q^2)}$$



Vacuum
Polarization

• $\Delta\alpha = \Delta\alpha^{\text{leptons}} + \Delta\alpha^{\text{quarks}} *$ (* see figures)

$1 + \Delta\epsilon(t)$ All the Radiative corrections

$$O(\alpha) \quad \alpha L^2, \alpha L, \alpha \text{ const.}$$

$$O(\alpha^2) \quad \alpha^2 L^4, \alpha^2 L^3, \alpha^2 L^2, \alpha^2 L$$

• Neglected terms $\approx O(10^{-4})$

- Weak Radiative Corrections $O(10^{-5})$
- QED corrections to weak contrib. $O(10^{-6})$
- Interference $S \leftrightarrow t \quad O(10^{-5})$
- Semi-collinear configurations of pair production $O(10^{-4})$
- Higher orders $n \text{ photons } n > 4$

⋮

The Standard Model

$$SU(2) \otimes U(1)$$

contains QED in an
Effective Way

$$\text{scales } q^2, t \ll M_Z^2, M_W^2$$

- Small angle Bhabha scattering is a QED Process effectively*
(* see table)

- The $\alpha(t)$ corresponds to the QED coupling evaluated at the scale t

for small angles

$$t = -\frac{s}{2}(1 - \cos\theta) \approx -\frac{s}{2}\left(-\frac{\theta^2}{2}\right) \approx \frac{s\theta^2}{4}$$

$$\sqrt{s} = 91.187$$

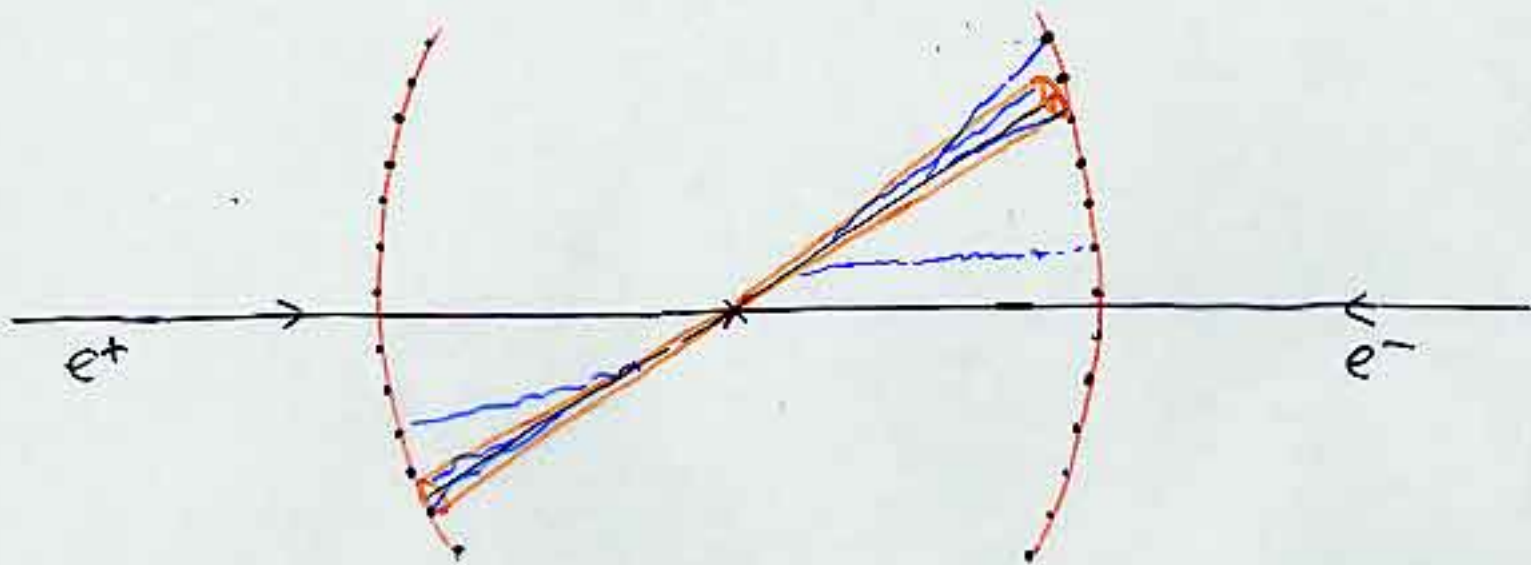
$$\theta = 30 \text{ mrad}$$

$$t = 2.2 \text{ GeV}^2$$

$$\theta = 150 \text{ mrad}$$

$$t = 30 \text{ GeV}^2$$

The Luminosity Data Sample



• Monte Carlo Codes

- NLLBHA Semianalytical ($\alpha^2 L$)
- LABSMC Monte Carlo for ($\alpha^2 L$)
Small Angle Bhabha

to be compared with

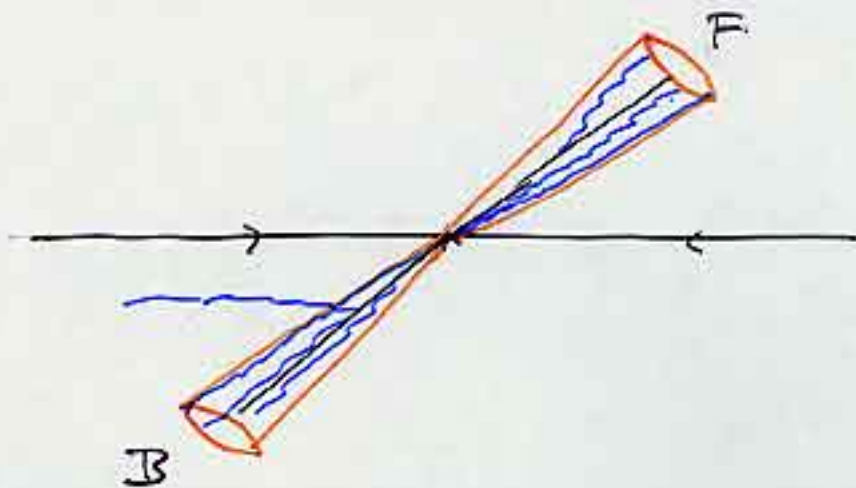
- BHLUMI

used to determine the Luminosity by the LEP Collaboration

{ NLLBHA Contains EW Matrix Elements
 LABSMC Born + O(α) QED
 BHLUMI

Describe all the small angle region with $\theta \gg \frac{m_e}{E_{beam}}$

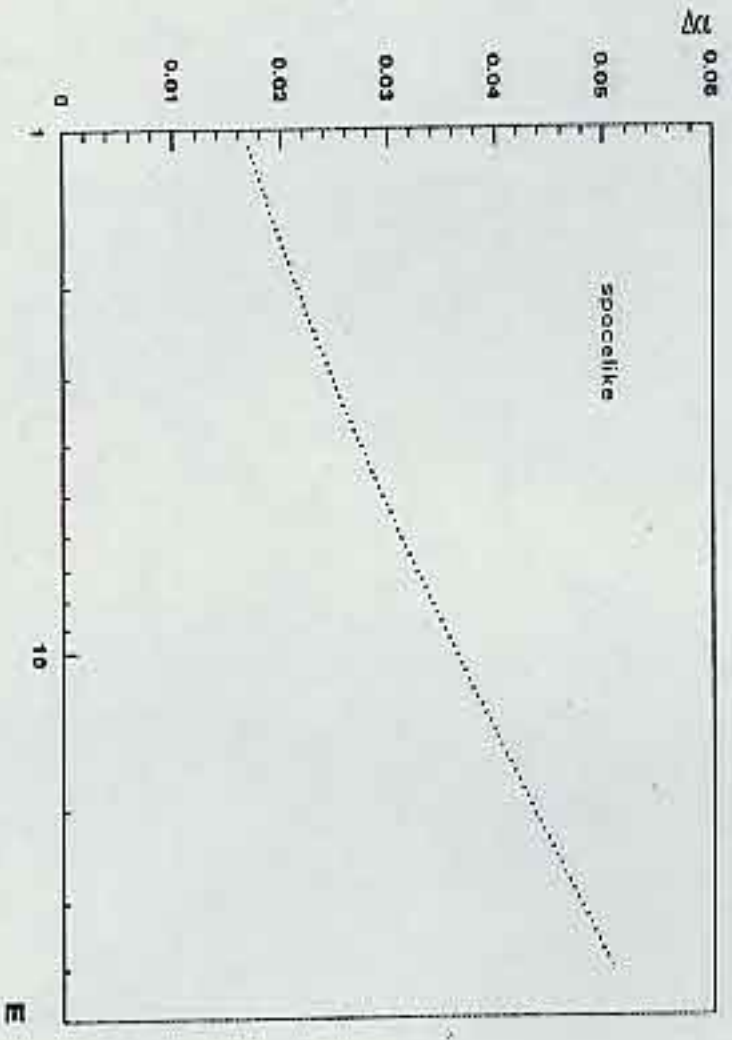
Calorimetric type Measurement



1. The cluster is reconstructed
2. Find the cluster center
3. Consider only the highest energy cluster in each hemisphere F, B
4. Energy Cuts $\min(E_F, E_B) > 0.65 E_{\text{beam}}$
 $\max(E_F, E_B) > 0.95 E_{\text{beam}}$
this suppresses the effects of the initial state radiation
5. Take the angle of the center of the most energetic cluster
 $\Rightarrow \theta$
6. This for several "rings" segmented detector to reproduce the differential distribution
7. Can be extended to NLC energies



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Figure 1: $\Delta\alpha$ versus $E = \sqrt{-t}$

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DESY qa-07

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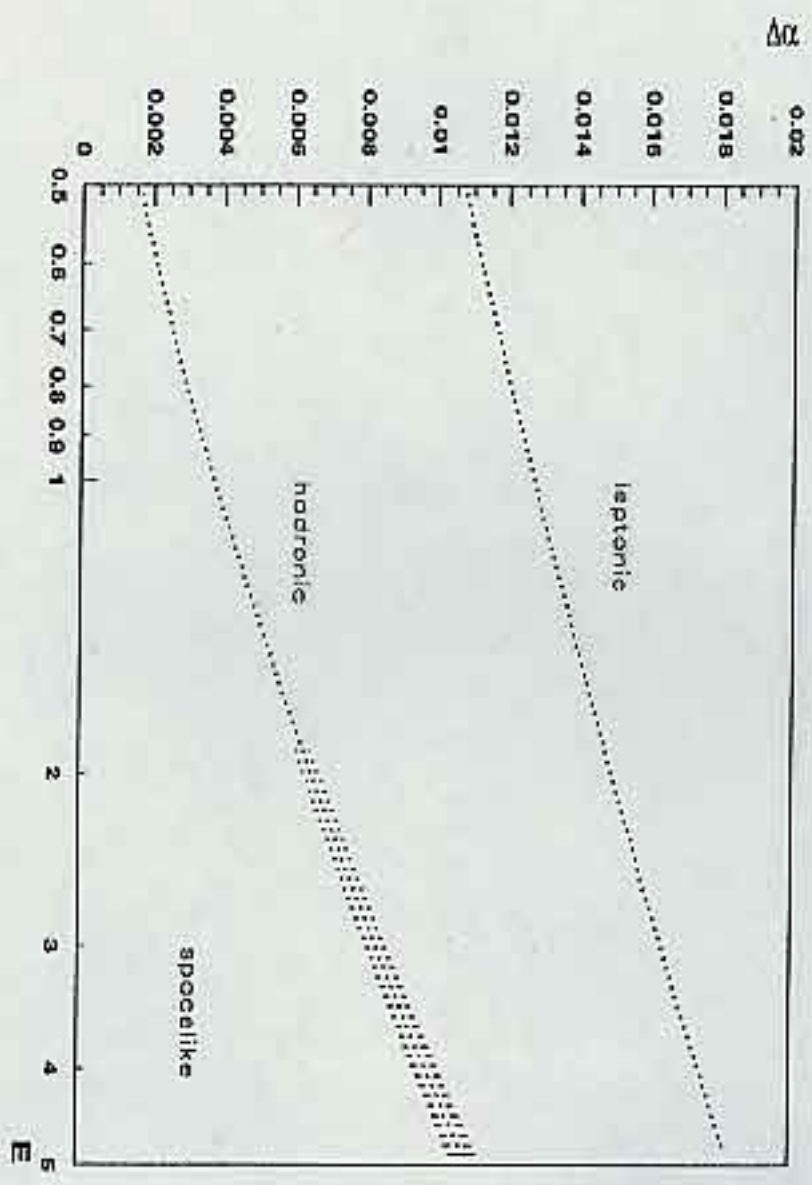


Figure 2: $\Delta\alpha$ for leptons and hadrons versus $E = \sqrt{-t}$

The cross section is integrated over angular range 45–110 mrad.

E_{CM}	91.187 GeV	91.2 GeV	189 GeV	206 GeV
QED	51.428	51.413	11.971	10.077
QED _t	51.484	51.469	11.984	10.088
EW	51.436	51.413	11.965	10.072
EW+VP _t	54.040	54.016	12.743	10.745
EW+VP	54.036	54.013	12.742	10.744

add a column for 800 GeV and angular range 5 to 50 mrad

NLC, TESLA

$$\frac{d\sigma}{dt} = \frac{d\sigma^0}{dt} \left(\frac{\alpha(t)}{\alpha(0)} \right)^2 (1 + \Delta\epsilon(t))$$

In Table 2 we give the theoretical predictions for the ingredients of Eq. (3), averaged over the 7 rings:

$$\sigma_i^0 = \int^{R_i} dt \frac{d\sigma^0}{dt}, \quad \langle (\alpha(t)/\alpha(0))^2 \rangle = \int \frac{dt}{t_{\max} - t_{\min}} \left(\frac{\alpha(t)}{\alpha(0)} \right)^2, \quad (1)$$

$$\langle 1 + \Delta\epsilon(t) \rangle = \frac{1}{\sigma_i^0 \langle (\alpha(t)/\alpha(0))^2 \rangle} \int^{R_i} dt \frac{d\sigma}{dt}.$$

Table 1: Theoretical predictions for the averages over the rings.

No. of ring	1	2	3	4	5	6	7
$E_{\text{CM}} = 91.2 \text{ GeV}$							
σ_i^0	20.078	7.8711	3.8737	2.1866	1.3534	0.8951	0.6224
$\langle (\alpha(t)/\alpha(0))^2 \rangle$	1.0425	1.0475	1.0516	1.0551	1.0582	1.0609	1.0634
$\langle 1 + \Delta\epsilon(t) \rangle$	0.9426	0.9439	0.9411	0.9395	0.9241	0.8918	0.7986
$E_{\text{CM}} = 189 \text{ GeV}$							
σ_i^0	4.6745	1.8323	0.9016	0.5088	0.3148	0.2081	0.1446
$\langle (\alpha(t)/\alpha(0))^2 \rangle$	1.0554	1.0613	1.0661	1.0702	1.0736	1.0767	1.0794
$\langle 1 + \Delta\epsilon(t) \rangle$	0.9376	0.9387	0.9355	0.9322	0.9155	0.8846	0.7885
$E_{\text{CM}} = 200 \text{ GeV}$							
σ_i^0	13.117	5.1422	2.5309	1.4284	0.8839	0.5846	0.4064
$\langle (\alpha(t)/\alpha(0))^2 \rangle$	1.0565	1.0625	1.0673	1.0714	1.0749	1.0780	1.0807
$\langle 1 + \Delta\epsilon(t) \rangle$	0.9374	0.9384	0.9347	0.9324	0.9150	0.8835	0.7881

← $\pm 80(10^{-4})$

Conclusions

- A well established region of the phase space exists where it is possible to test QED and where weak effects are negligible
- The running of α can be measured directly within a pure QED framework at LEP and at NLC energies
- $\Delta\alpha = \Delta\alpha^{\text{lept}} + \Delta\alpha^{\text{had}}$, if the precision of the data allows, opens the possibility of separately determine $\Delta\alpha^{\text{had}}$ in an extended range of energies