



# Tracking and Fitting

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*High Energy Physics* —

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BaBar Detector Physics Series

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# Outline

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- ◆ BaBar tracking devices: SVT and DCH
- ◆ Track finding
  - SVT stand-alone
  - DCH stand-alone
  - An alternative approach
- ◆ Track fitting
  - Method of maximum likelihood
  - Kalman filtering
- ◆ Constrained fitting
  - vertexing



# What Exactly IS Tracking?

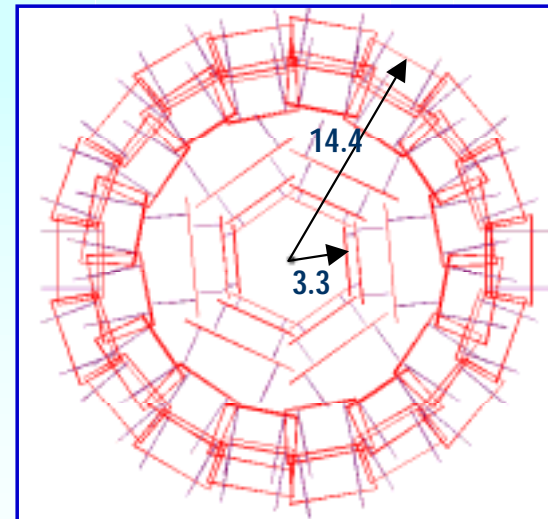
- ◆ Tracking as such consists of two parts:
  - Track **finding** (pattern recognition)
  - Track **fitting**
    - really comes down to minimizing  $\chi^2$ , which is a quantity that measures how close the measured parameters are to what they are assumed to be from a particular fit hypothesis (e.g., helical trajectory)
    - fitting would be trivial if it was not for complications arising because of multiple scattering, energy loss, non-uniform magnetic field, etc., and if we understood our detectors perfectly.



# BaBar Silicon Vertex Detector



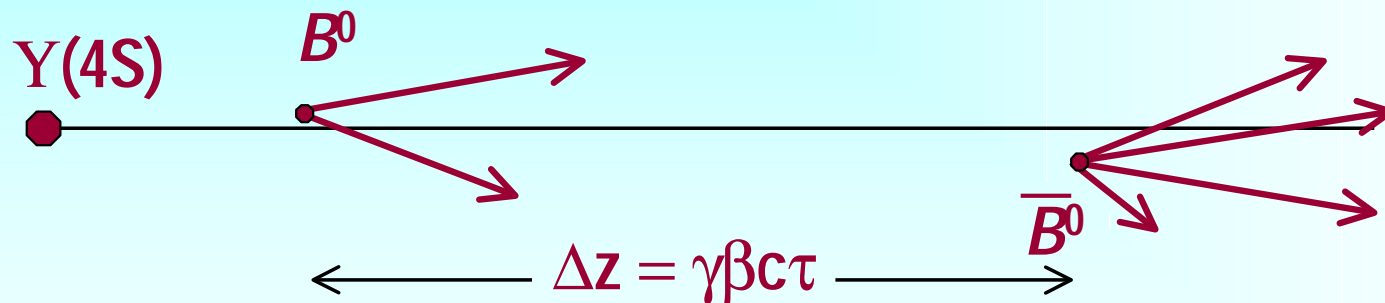
- ◆ 5 layers of double-sided silicon microstrip detectors
- ◆  $\sim 0.94 \text{ m}^2$  of silicon
- ◆  $\sim 150,000$  electronics readout channels





# Main Goal of the Vertex Detector

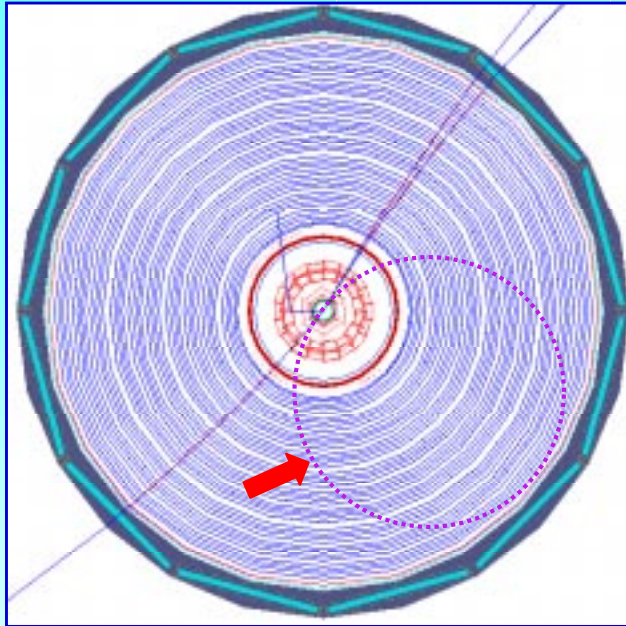
- ◆ Good vertexing is crucial for many analyses, and indispensable for CP violation studies.
- ◆ The main purpose of the BaBar vertex detector is to determine the separation between the two  $B$  decay vertices along the  $z$  axis:



- ◆ However, it is not its only function!



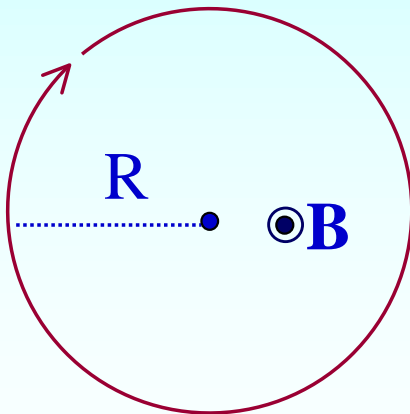
# Tracks in SVT and DCH



- ◆ For a charged particle in a uniform magnetic field  $\mathbf{B}$ :

$$|\mathbf{P}_T| \text{ (GeV/c)} = 0.3 |\mathbf{B}| \text{ (T)} R \text{ (m)}$$

where  $P_T$  = particle momentum component in a plane  $\perp$  to  $\mathbf{B}$ , and  $R$  is the trajectory radius





# Why 5 layers?

- ◆ Compare the following parameters for the three new  $B$ -factory detectors (all operating in a 1.5 T  $\mathbf{B}$ -field):

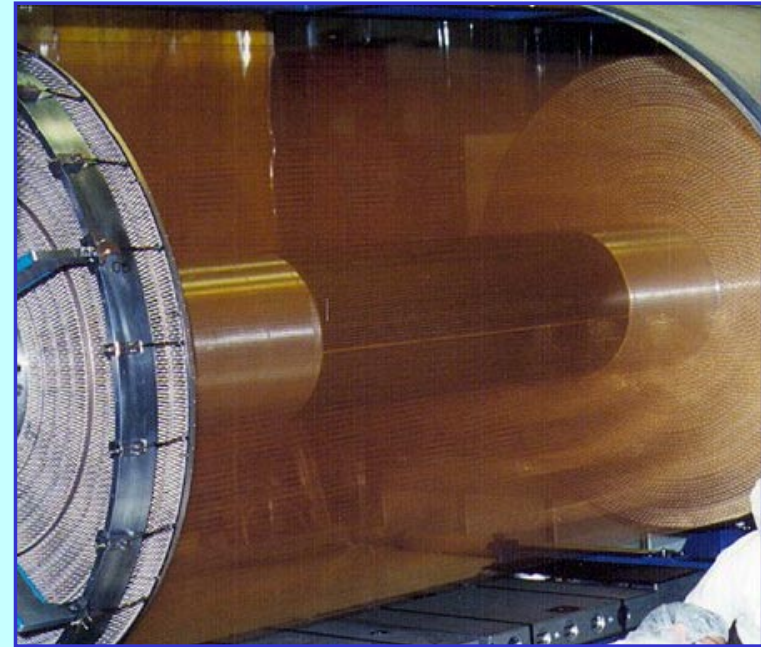
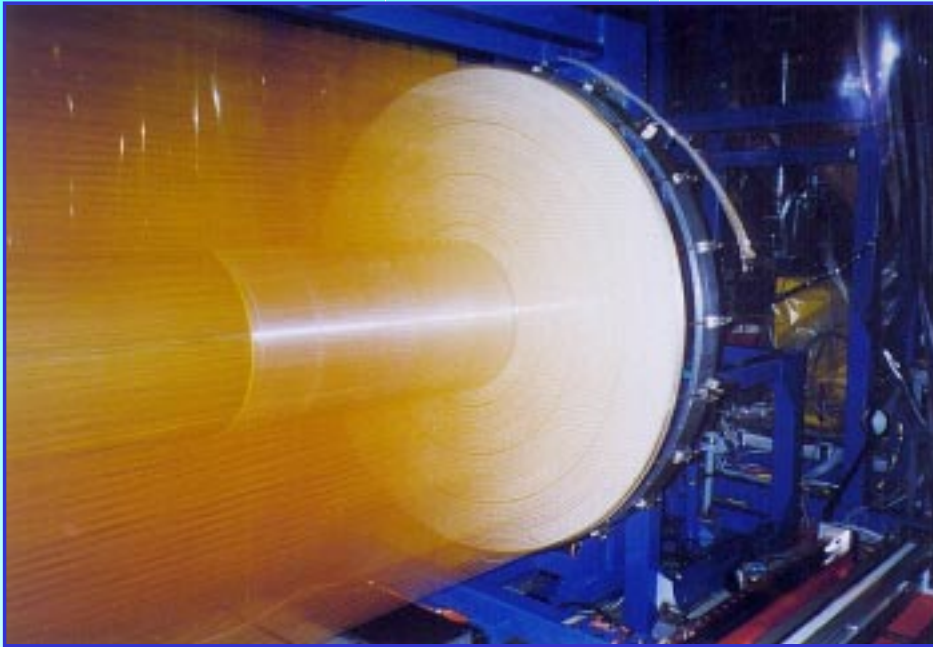
Experiment	# layers in SVT	Inner radius of DCH	Minimum $P_T$ at inner radius of DCH
CLEO-III	4	17.5 cm	39 MeV/c
Belle	3	8.5 cm	19 MeV/c
BaBar	5	22.5 cm	51 MeV/c

courtesy Doug Roberts

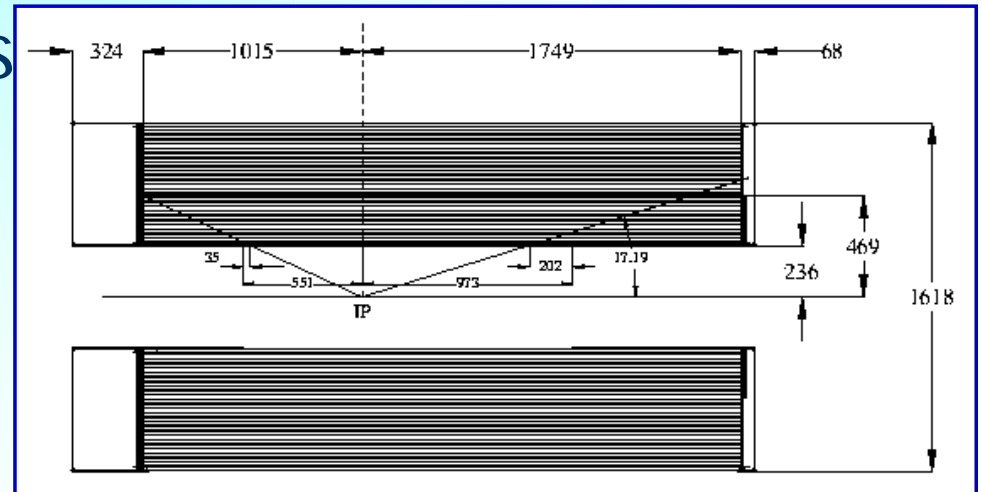
- ➔ BaBar SVT must not only do vertexing, but also perform tracking for low momentum tracks (up to  $P_T \sim 120$  MeV/c).



# BaBar Drift Chamber



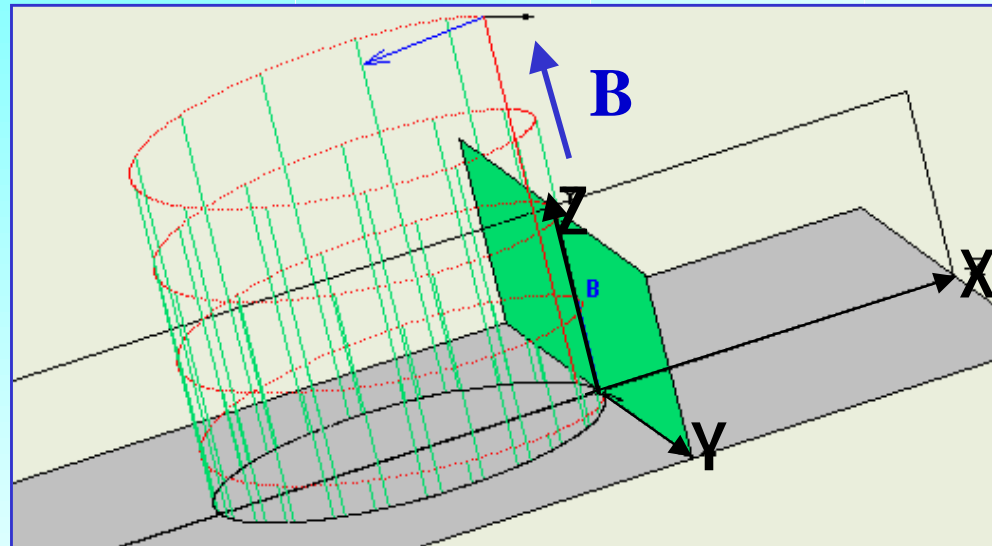
- ◆ 40 layers with 7,104 drift cells
- ◆ Layers organized into 10 superlayers, following the pattern AUVAUVAUVA (*axial (A)-stereo(U,V)*)





# Particle in a Uniform B-field

- ◆ Recall that a charged particle in a uniform magnetic field moves along a helix.



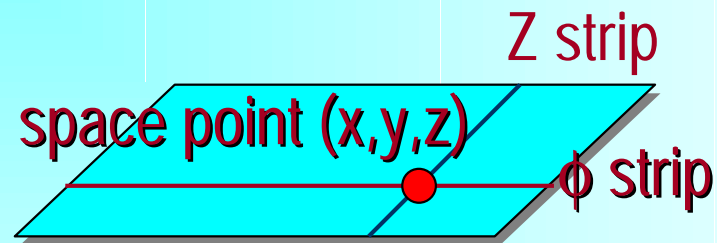
$$\mathbf{B} \parallel \mathbf{z}$$

- ◆ But this can be decoupled into
  - moving along a circle in the  $xy$ -plane (need **3** points to define it)
  - and moving along a straight line in  $z$  (need **2** points to define it)

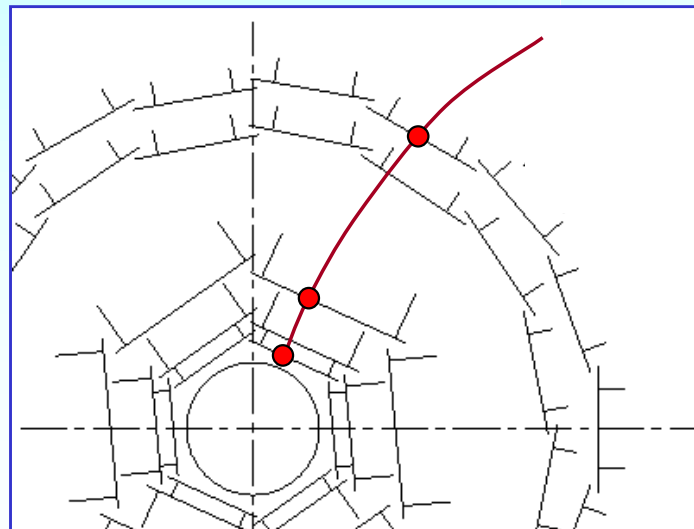


# Track Finding in the SVT

- ◆ SVT stand-alone pattern recognition:
  - form **space points** from matching  $\phi$  and Z hits:



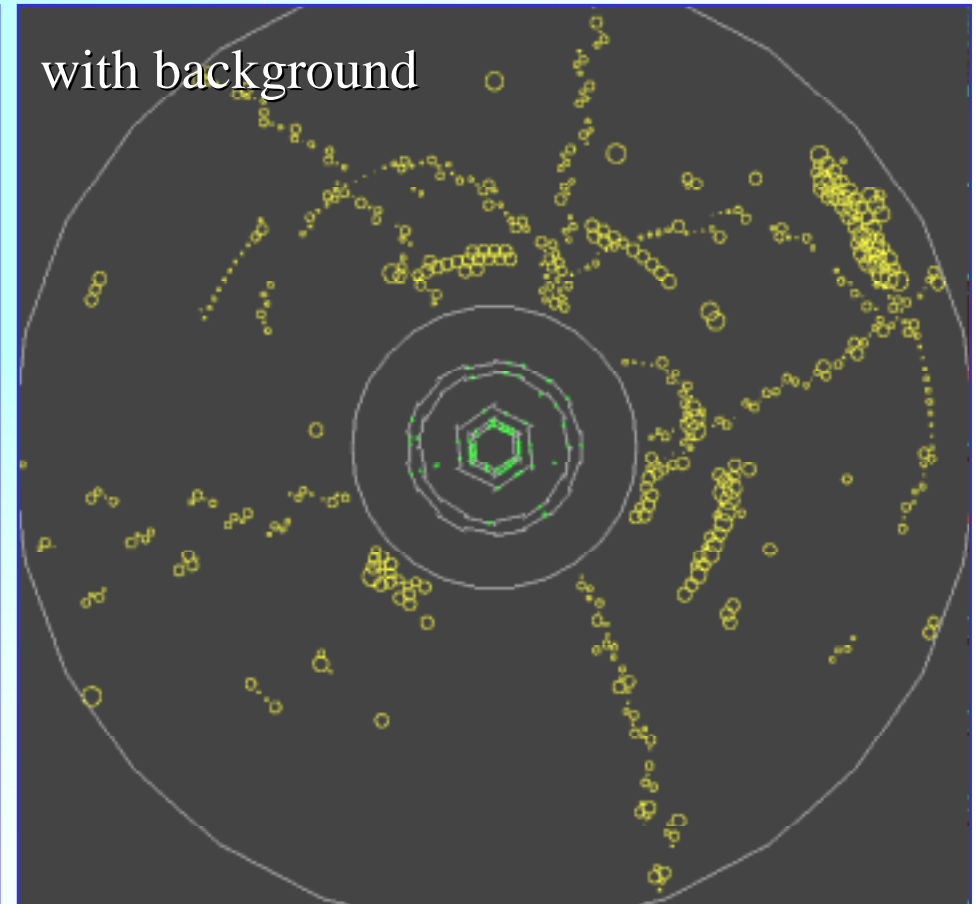
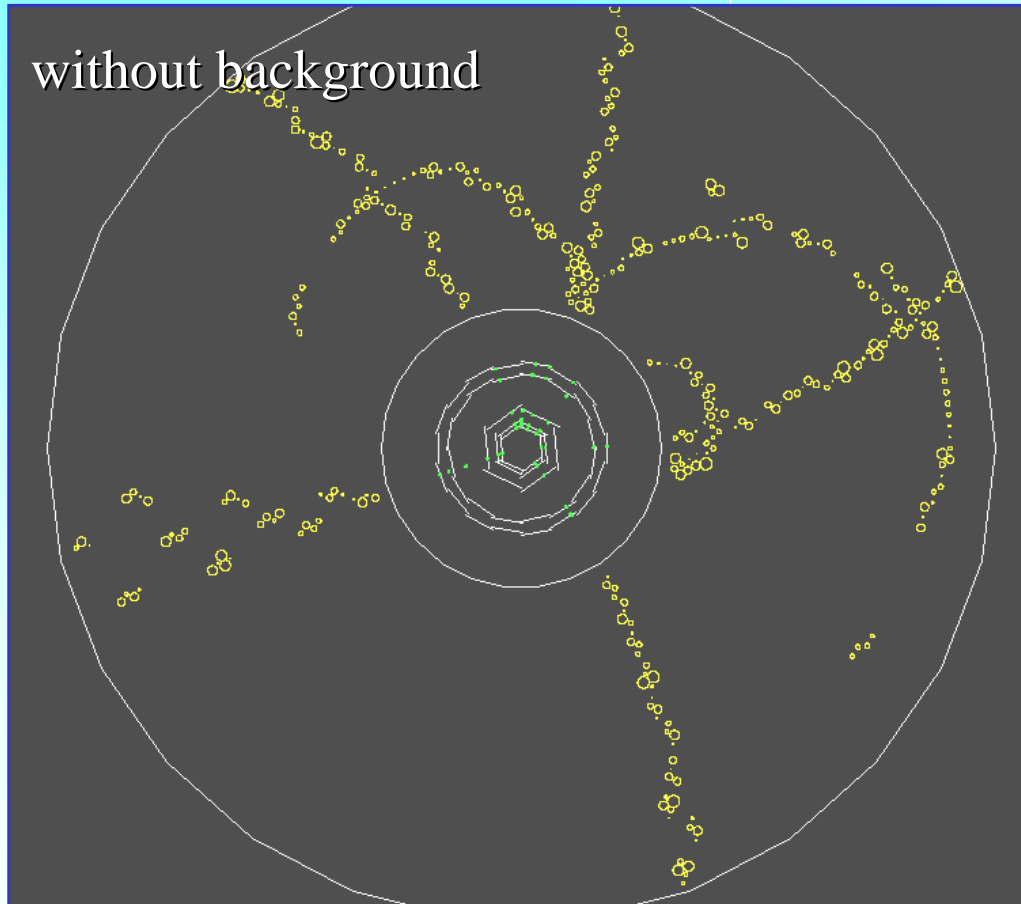
- find 3 space points on different layers that might form a track:





# Pattern recognition

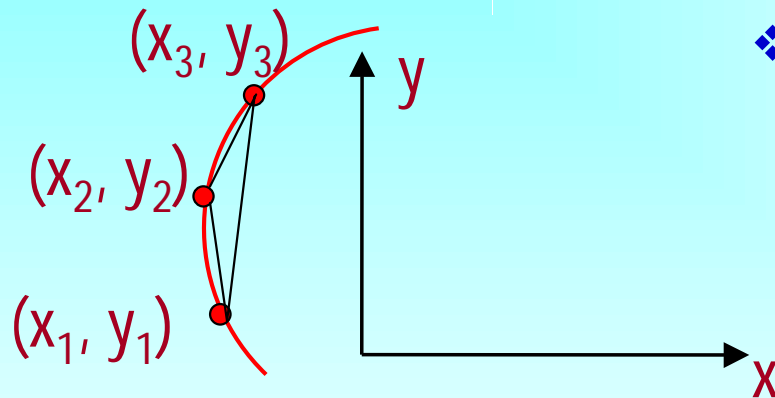
- ◆ Pattern recognition is sometimes easier for the eye than for the computer...





# Track Fitting in the SVT

- ◆ Once three suitable points are found, it is possible to determine the  $r$ - $\phi$  parameters of a circular track



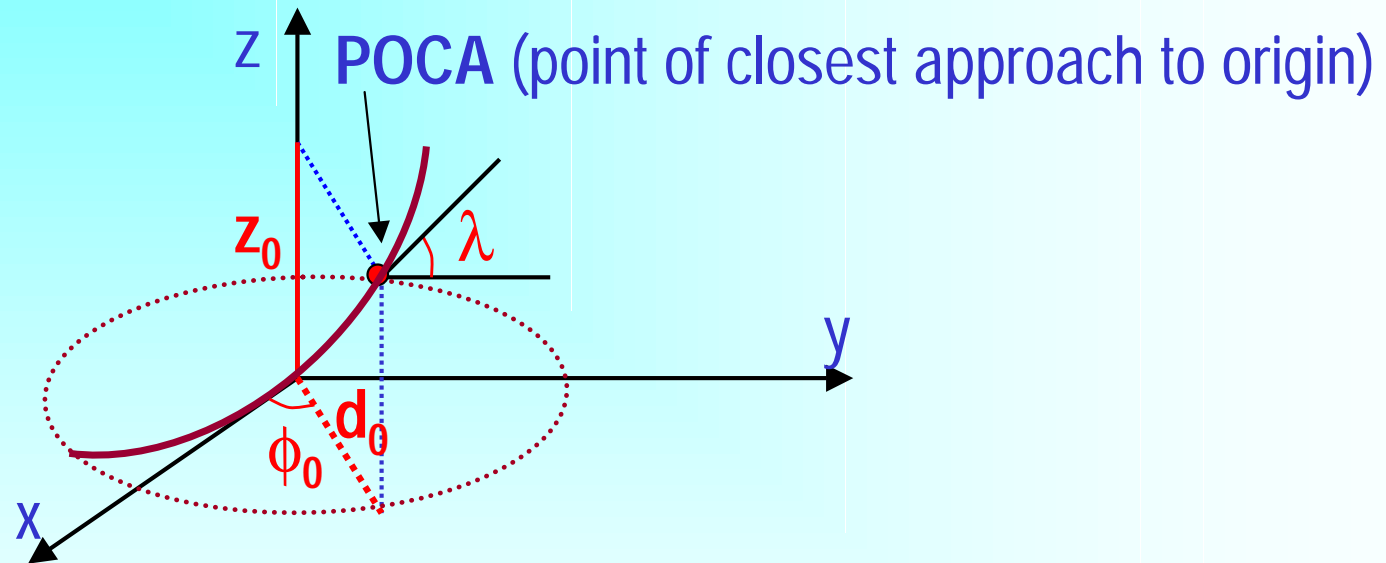
- ❖ The points are required:
  - to be close enough together azimuthally
  - to have consistent times
  - to lie in a "road" defined as  $a + b/p_T$   
(2nd term accounts for multiple scattering)

- ◆ Then add  $z$  information (for at least two hits) to the circle... and turn it into a helix!



# Track Fitting in the SVT (cont.)

- ◆ At the end, we end up with tracks fit to helices with 5 parameters: ( $d_0$  (cm),  $\phi_0$ (rad),  $\omega$ (cm<sup>-1</sup>),  $z_0$ (cm),  $\tan\lambda$ )



$\omega$  = *geometrical curvature*

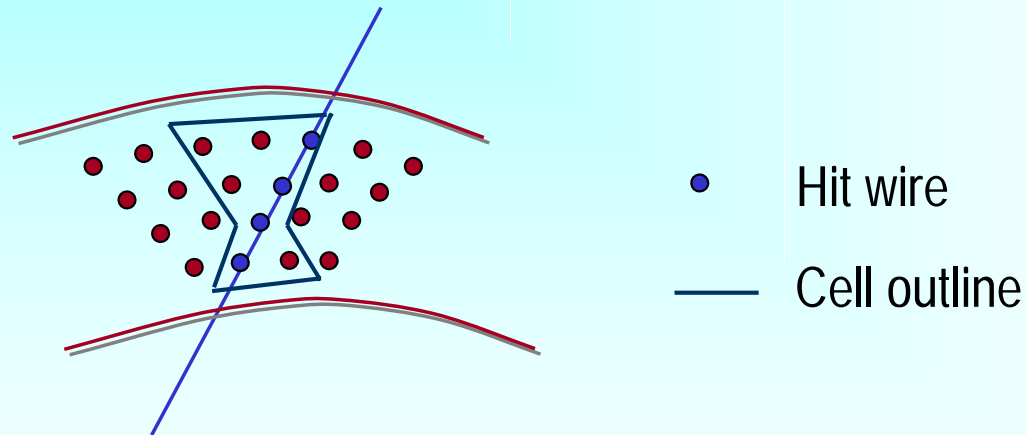
$r$  = *radius of curvature*

$$r = \frac{1}{\omega} = \frac{10^{13}}{qB_z c} p_T = \frac{333.567 p_T}{qB_z}$$



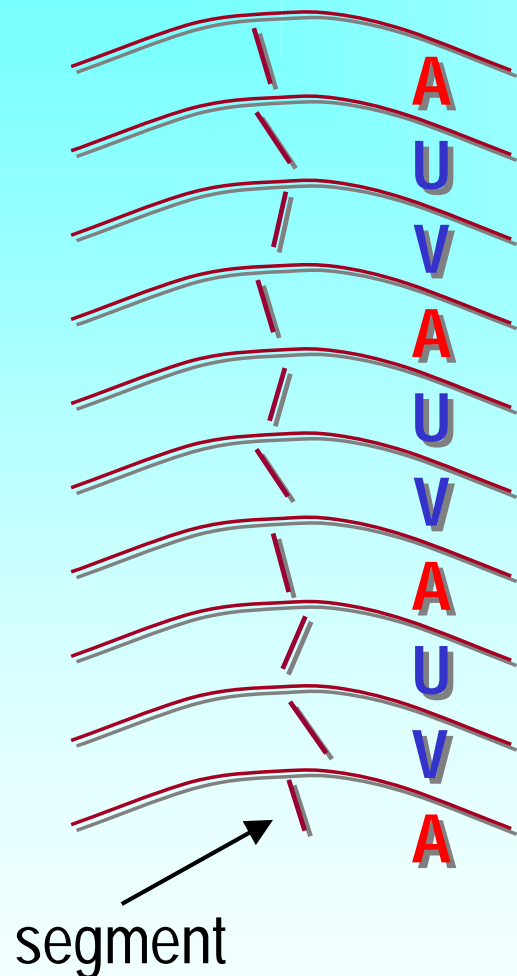
# Track Finding in the DCH

- ◆ Tracking in the DCH is based on **segments**
- ◆ Segment finder looks for *patterns* of hit wires and calls the valid ones segments.
- ◆ Here is an example of a valid pattern:





# Stereo/Axial Layers in the DCH

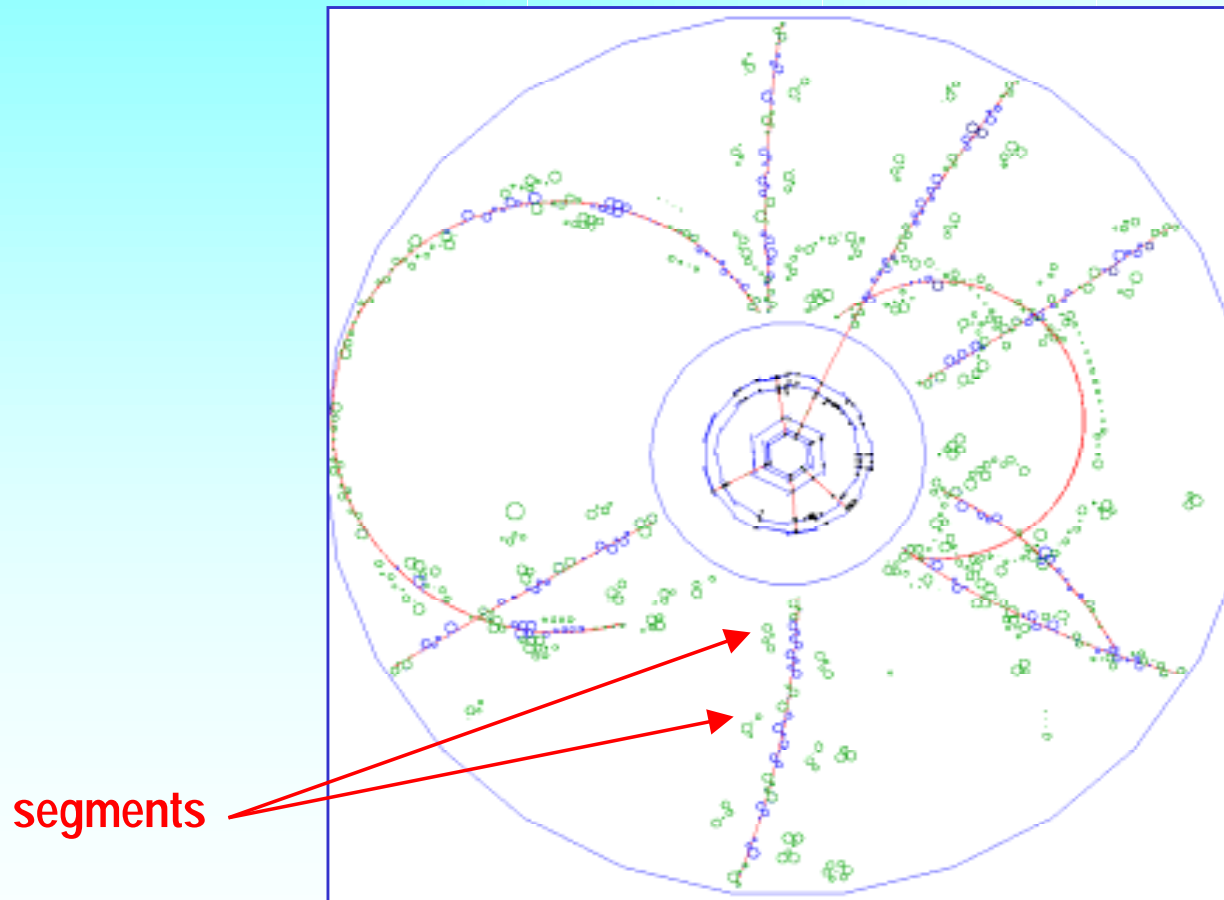


- ◆ Three segments in the **axial** layers are used to form a circle in the xy plane
- ◆ Then z-measurements from the **stereo** layers are added... and we have a helix!



# BaBar DCH Track Finders

- ◆ There are two independent segment-based track finders run in series
  - DchTrackFinder for tracks coming radially from origin ( $|d_0| < 1\text{cm}$ )
  - DcxTrackFinder for all tracks



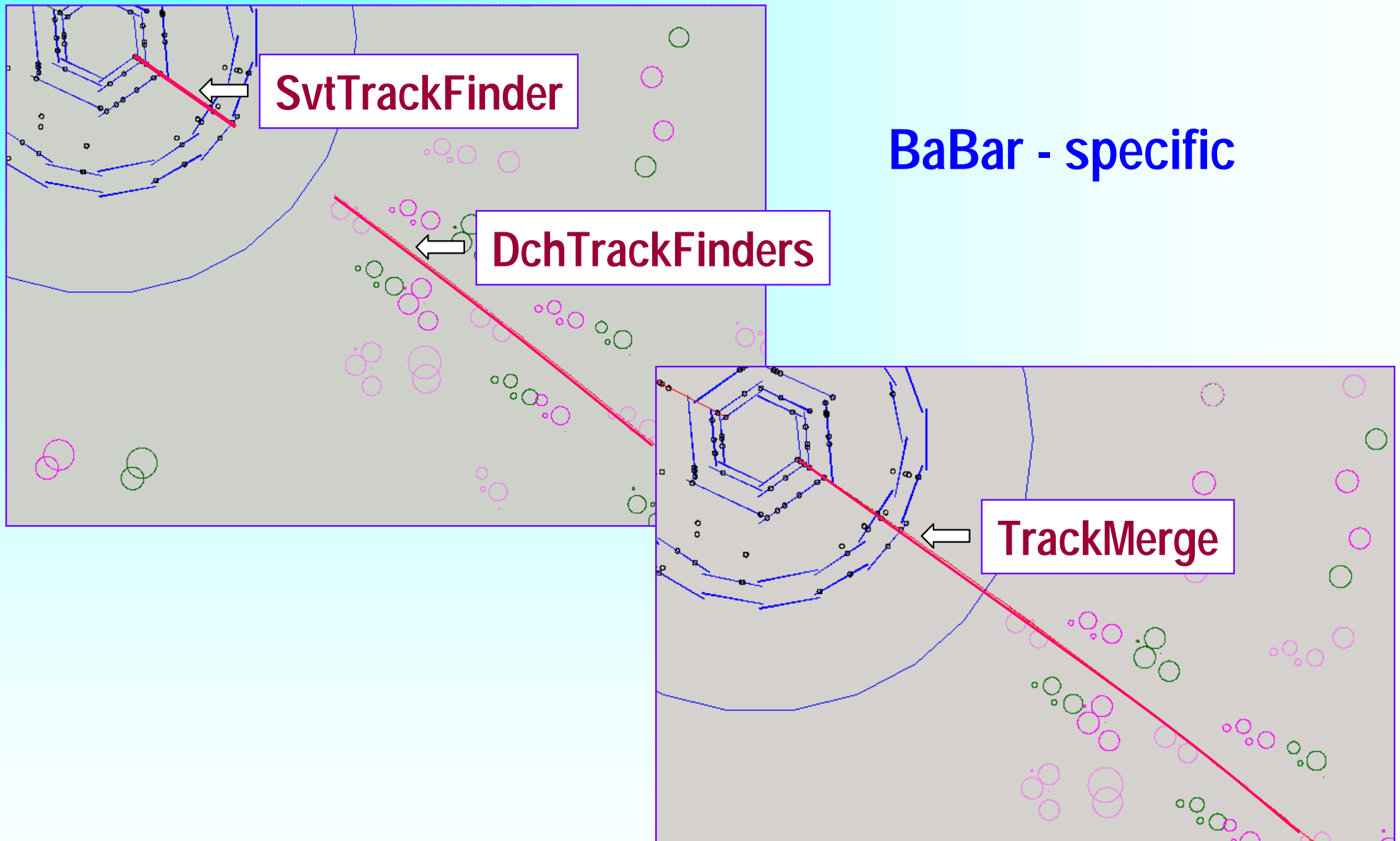
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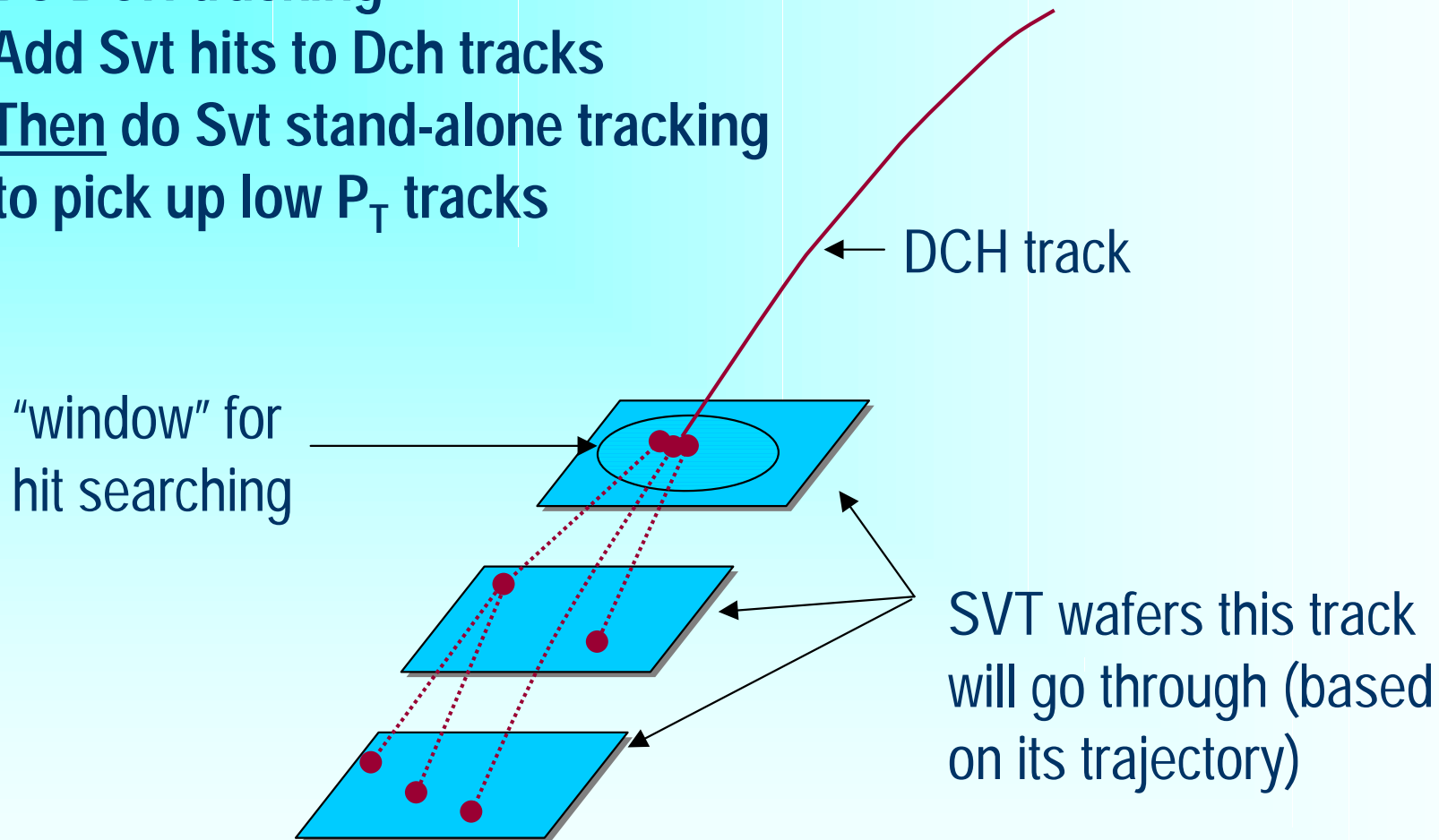
# Combining SVT and DCH Tracks





# An Alternative Approach

- ◆ Do DCH tracking
- ◆ Add Svt hits to Dch tracks
- ◆ Then do Svt stand-alone tracking to pick up low  $P_T$  tracks





# What Defines This Window?

- ◆ One can show that:
  - if we have  $n$  measuring planes extending to max radius  $R$
  - each with intrinsic spatial resolution  $\varepsilon$
  - and evenly spaced
- ◆ then the track parameter resolutions will be:

$$\sigma_{d_0} = \sqrt{\frac{9}{n}} \varepsilon$$

$$\sigma_{\phi_0} = \sqrt{\frac{192}{n}} \frac{\varepsilon}{R}$$

$$\sigma_{z_0} = \sqrt{\frac{4}{n}} \frac{\varepsilon}{\mu}$$

stereo angle

- ◆ so, for  $n = 40$ ,  $\varepsilon = 140 \mu\text{m}$ ,  $R = 81 \text{ cm}$ ,  $\mu = 50 \text{ mrad}$ , we get:

$$\sigma_{d_0} = 66 \mu\text{m}, \quad \sigma_{\phi_0} = 0.4 \text{ mrad}, \quad \sigma_{z_0} = 885 \mu\text{m}$$



# What about the SVT?

- ◆ The SVT gives much better resolution on  $d_0$ ,  $z_0$ ,  $\tan\lambda$ !
  - The intrinsic resolutions are of the order of 10-15  $\mu\text{m}$  in xy plane, and 15-30  $\mu\text{m}$  in z.
  - The track parameter resolutions are dominated by multiple scattering, not by intrinsic resolution, for most of the momentum range.
  - at some  $r$  ( $=x^2+y^2$ ):

$$\sin(\phi - \phi_0) = \frac{(b/r)(1 + \omega d_0) + \omega r}{1 + 2\omega d_0}$$

- ➔ For angles, need good resolution on the momentum, which comes (mostly) from the DCH!



# Momentum Resolution in the DCH

- ◆ For many ( $n > 10$ ) position measurements, the **curvature resolution** is:

$$\sigma_{\omega} = \frac{\varepsilon}{L'^2} \sqrt{\frac{720}{n+4}}$$

projected length of track  
onto bending plane

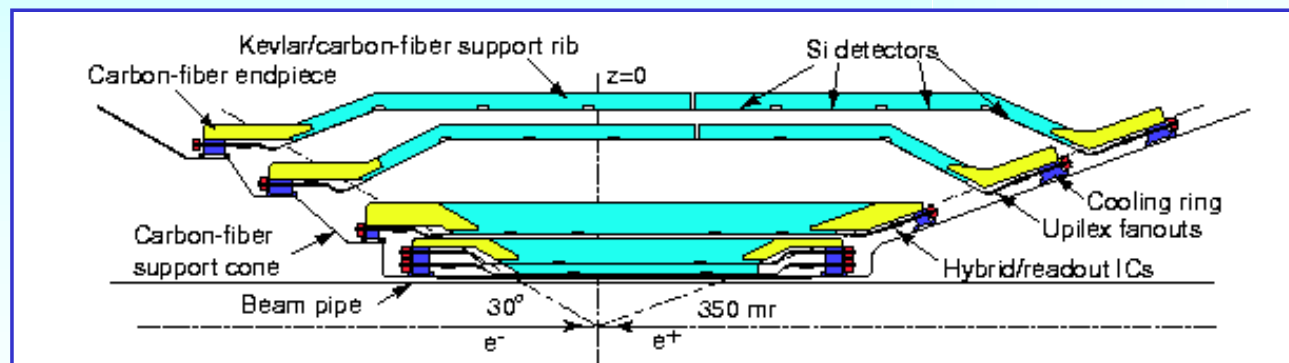
- ◆ Assuming intrinsic resolution  $\varepsilon = 140 \mu\text{m}$ ,  $n = 40$  measurements, and  $L' = 52.2 \text{ cm}$ , we get

$$\sigma(P_T)/P_T = 0.46\% P_T$$



# The more measurements the better?

- ◆ One can see that the resolutions  $\sim 1/\sqrt{n}$
- ◆ Therefore, the more measurements one makes the better?
  - Not always! If there is significant **multiple scattering**, adding measurements may *degrade* the resolution on some track parameters.
    - e.g., when one fits for the impact parameters, the **innermost** precision measurement from the SVT provides most useful information





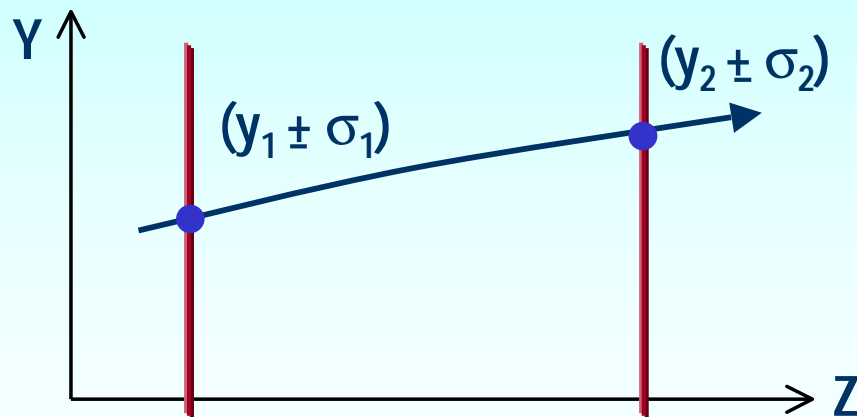
# Track Parameter Resolution

- ◆ The DCH information dominates the **momentum** resolution
- ◆ The SVT information dominates the **impact parameter** resolution, both in  $xy$  and  $z$  directions.
- ◆ But how exactly does one obtain track parameter resolutions?
  - = How does fitting really work?



# Least Square Track Fitting: an Example

- ◆ In order to start fitting a track, one needs two things:
  - a **model** which approximates the track's trajectory
  - an understanding of the **detector accuracy** (resolution)
- ◆ Let's consider an example: two measurements to be fit to a straight line



Assume that the two points  
can be fit to

$$f_i = a + z_i b, i = 1, 2$$



# Some General Considerations

- ◆ Suppose we have  $n$  data values  $y_l$  ( $l=1, \dots, n$ )
  - assume they are functions of  $m$  variables  $\alpha_i$  ( $m < n$ ):
$$y_l = f_l(\alpha_i)$$
  - assume each measurement  $y_l$  has a Gaussian measurement error  $\sigma_l$ .

- ◆ Then the probability density for the measurements is:

$$g(y_1, y_2, \dots, y_n) = \prod_l \frac{1}{\sqrt{2\pi\sigma_l}} \left( e^{-\sum_l \frac{(y_l - f_l)^2}{2\sigma_l^2}} \right)$$
$$\equiv \prod_l \frac{1}{\sqrt{2\pi\sigma_l}} \left( e^{-\sum_l \chi^2} \right), \text{ where } \chi^2 \equiv \sum_l \frac{(y_l - f_l(\alpha))^2}{\sigma_l^2}$$

- maximizing the probability density = minimizing the  $\chi^2$



# Fitting to a line

- ◆ We need to minimize

$$\chi^2 \equiv \frac{(y_1 - f_1(a, b))^2}{\sigma_1^2} + \frac{(y_2 - f_2(a, b))^2}{\sigma_2^2} = \frac{(y_1 - a - bz_1)^2}{\sigma_1^2} + \frac{(y_2 - a - bz_2)^2}{\sigma_2^2}$$

- ◆ So:

$$\frac{\partial \chi^2}{\partial a} = \frac{y_1 - a - bz_1}{\sigma_1^2} (-2) + \frac{y_2 - a - bz_2}{\sigma_2^2} (-2) = 0$$

$$\frac{\partial \chi^2}{\partial b} = \frac{y_1 - a - bz_1}{\sigma_1^2} (-2z_1) + \frac{y_2 - a - bz_2}{\sigma_2^2} (-2z_2) = 0$$

Two equations with two unknowns (a, b) give a unique solution:

$$b = (y_1 - y_2)/(z_1 - z_2)$$

$$a = (z_1 * y_2 - y_1 * z_2)/(z_1 - z_2)$$

(who would have thought?!)



# Generalization

- ◆ n independent measurements  $\mathbf{Y}$  which we would like to fit to function  $\mathbf{F}$  with  $l$  parameters  $\Theta$ :

$$\mathbf{Y} = \begin{pmatrix} y_1 \\ \dots \\ y_n \end{pmatrix} \quad \mathbf{F} = \begin{pmatrix} f_1 \\ \dots \\ f_n \end{pmatrix} \quad \Theta = \begin{pmatrix} \vartheta_1 \\ \dots \\ \vartheta_l \end{pmatrix}$$

measurements ( $y_i$ )    predictions ( $f_i$ )    parameters ( $a, b$ )

$$V \equiv \langle \delta y_i \delta y_j \rangle = \begin{pmatrix} \sigma_1^2 & & 0 \\ & \dots & \\ 0 & & \sigma_n^2 \end{pmatrix} \quad A = \begin{pmatrix} a_{11} & & a_{1l} \\ & \dots & \\ a_{n1} & & a_{nl} \end{pmatrix}$$

covariance or error matrix    coefficients ( $z_i, 1$ )

- ◆ The solution is:  $\Theta = (A^T V^{-1} A)^{-1} A^T V^{-1} \mathbf{Y}$   $\rightarrow$  need  $V^{-1}$ !



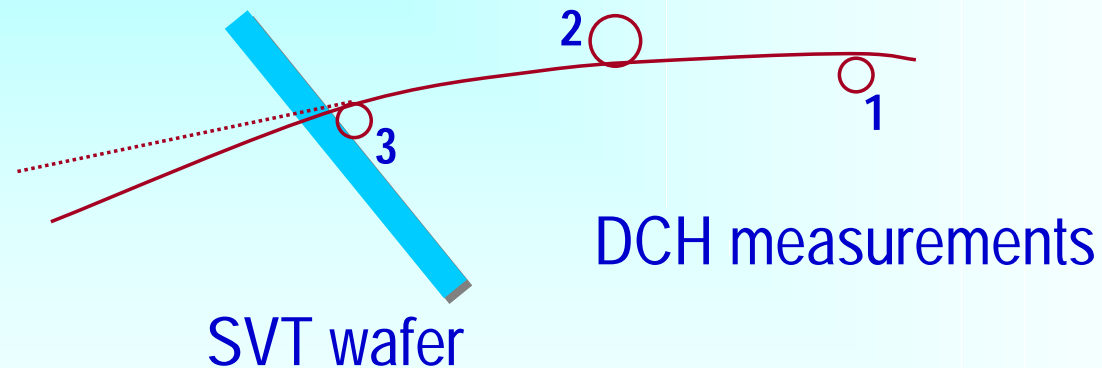
# Problems With This Approach

- ◆ This method is **global** in the sense that it fits all the measurements at the same time
- ◆ If all measurements are independent of each other, the execution time is  $\sim n$  (the # of measurements)
- ◆ But what if we have correlations between measurements?
  - the covariance matrix will contain non-diagonal terms
  - and inverting it becomes VERY time consuming for large  $n$ 
    - e.g., a track has 40 hits in the DCH and 5 in the SVT
    - that's  $40 + 2 \cdot 5$  (2 views in SVT) +  $2 \cdot 20$  (~20 pieces of material) = 90!
    - inverting a **90 x 90** matrix is no fun!



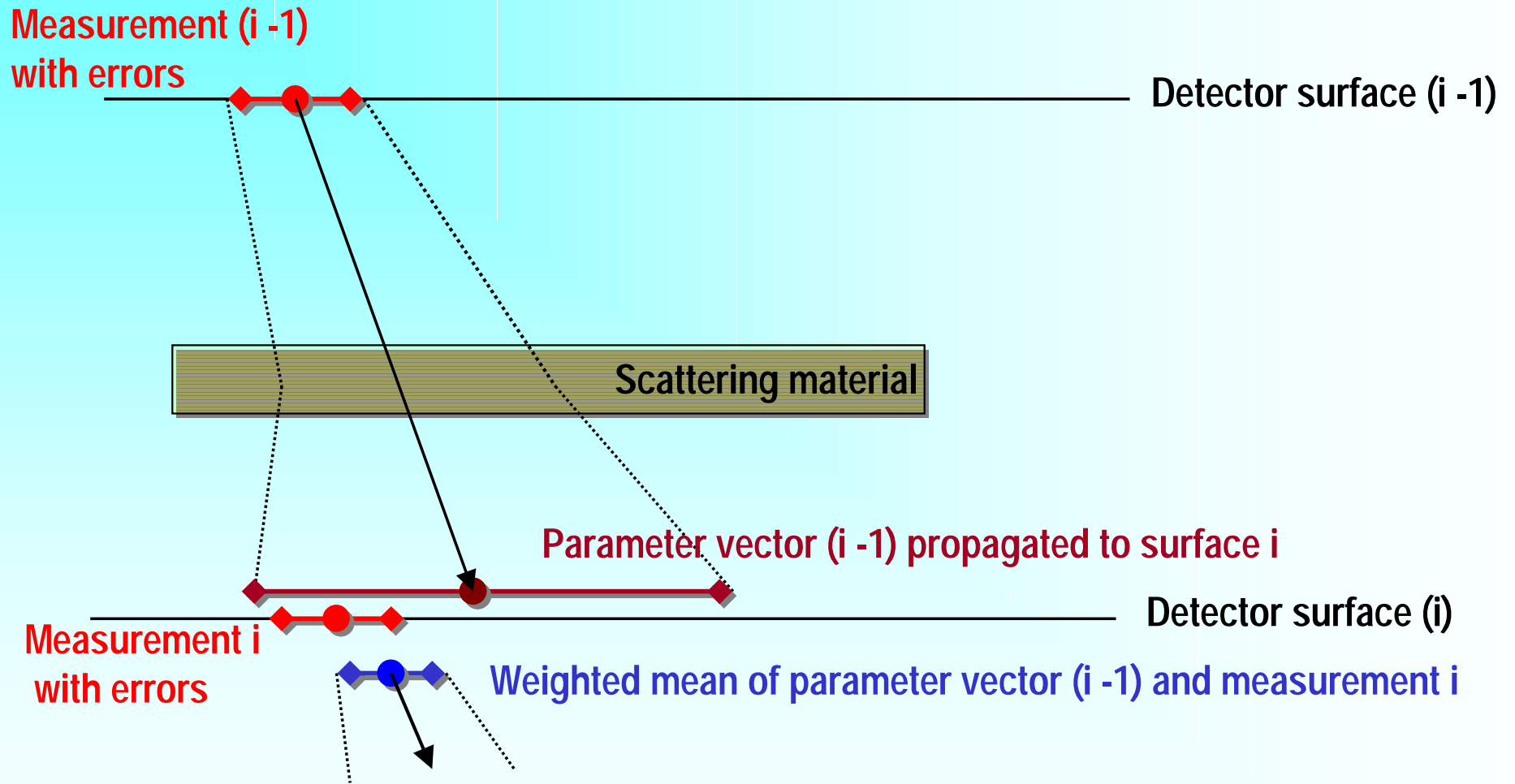
# The Solution: Kalman Filtering

- ◆ The idea:
  - estimate track parameters at every given point using previously obtained information + guess about the contributions from various physical processes (multiple scattering, energy loss,...)





# Kalman Filter





# Kalman Filter (cont.)

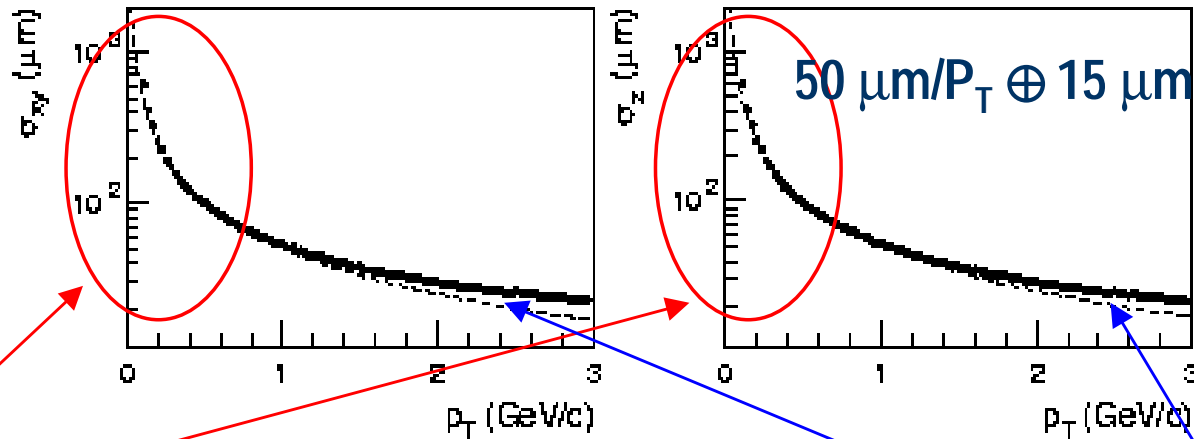
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- ◆ Kalman fitting is really about taking weighted averages!
- ◆ It's the simplest way to get the track errors right
  - with global fits, the errors cannot always be trusted
    - e.g., multiple scattering is most often ignored for speed...
    - ...and recall that in BaBar, the particles come from B's decaying almost at rest in the CM frame, so they have low momenta -- and multiple scattering becomes a very serious effect!

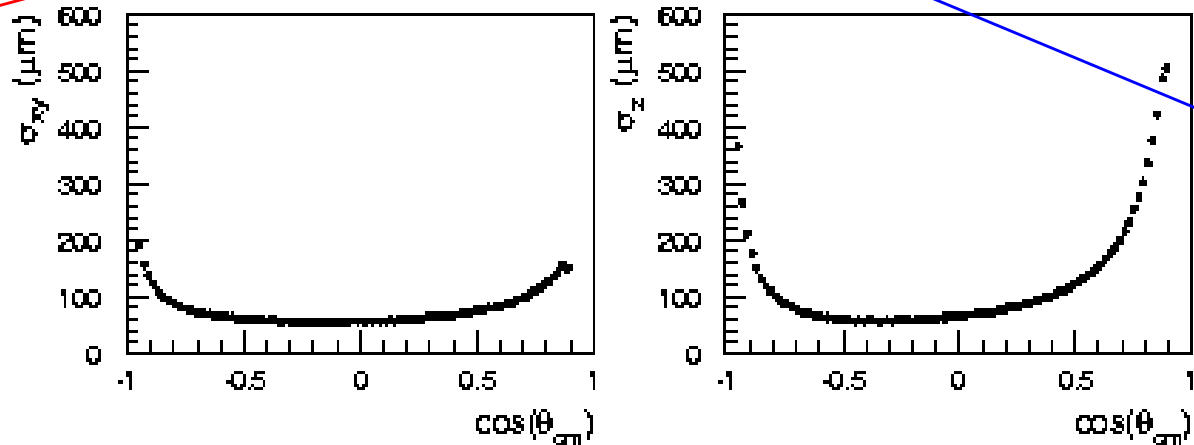


# Track Parameter Resolutions

## Kalman-fit tracks



multiple scattering degrades resolution!



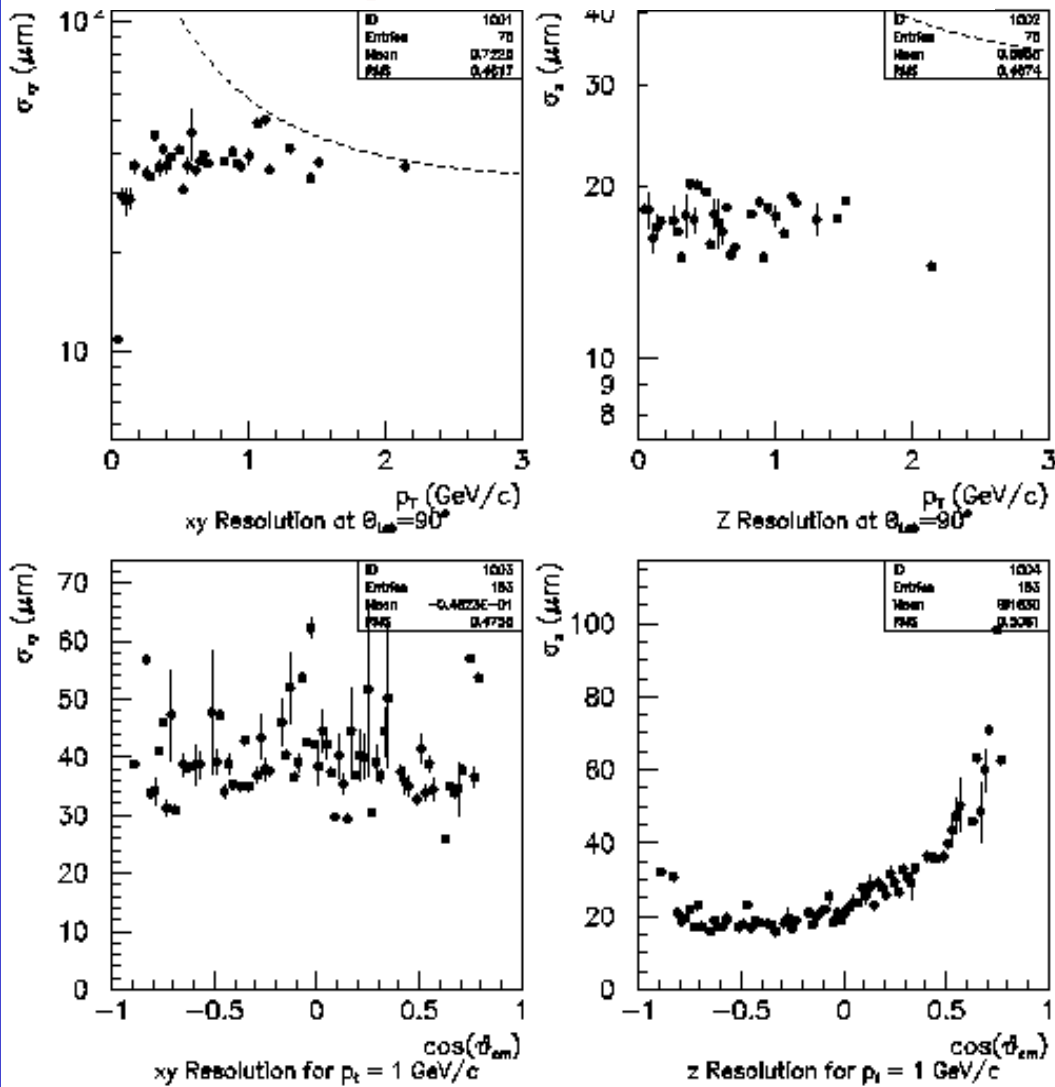
intrinsic resolution becomes limiting factor

courtesy Doug Roberts



# Track Parameter Resolution (cont.)

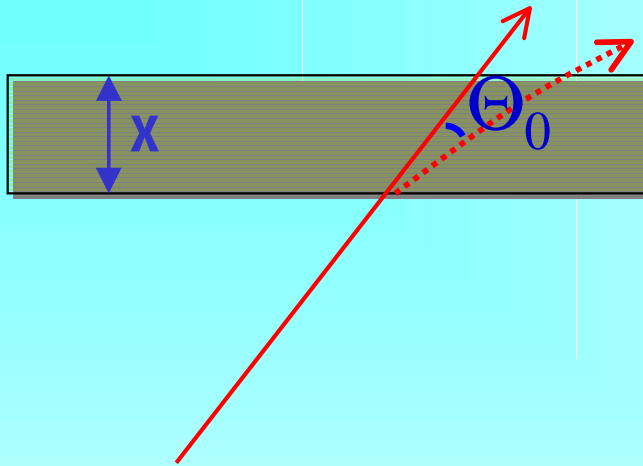
## Simple helix-fit tracks



→ multiple scattering is ignored, the errors are wrong!!



# Multiple Scattering



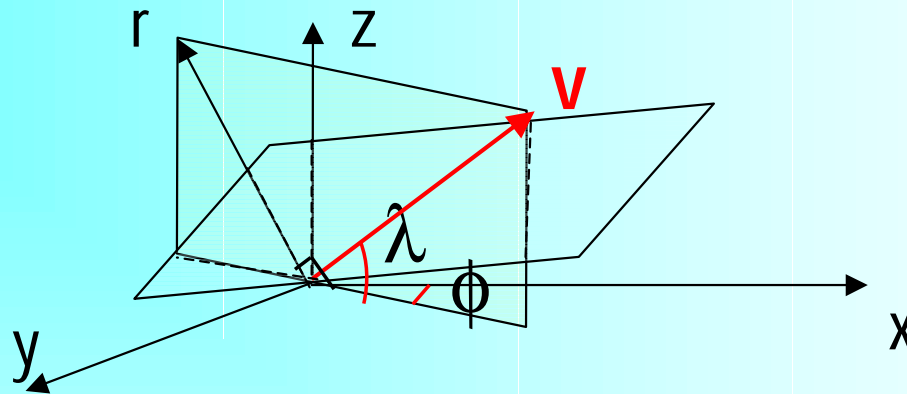
- ◆ Mostly due to Coulomb scattering from nuclei
- ◆ For small angles roughly Gaussian distribution:

$$\Theta_0 = \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{x / X_0} \left[ 1 + 0.038 \ln(x / X_0) \right]$$

►  $x/X_0$  is the thickness of the scattering material in radiation lengths



# Multiple Scattering in Track Fitting



- ◆ Multiple scattering does not affect the track's momentum, nor the track's fit parameters
- ◆ It only affects the error matrix:

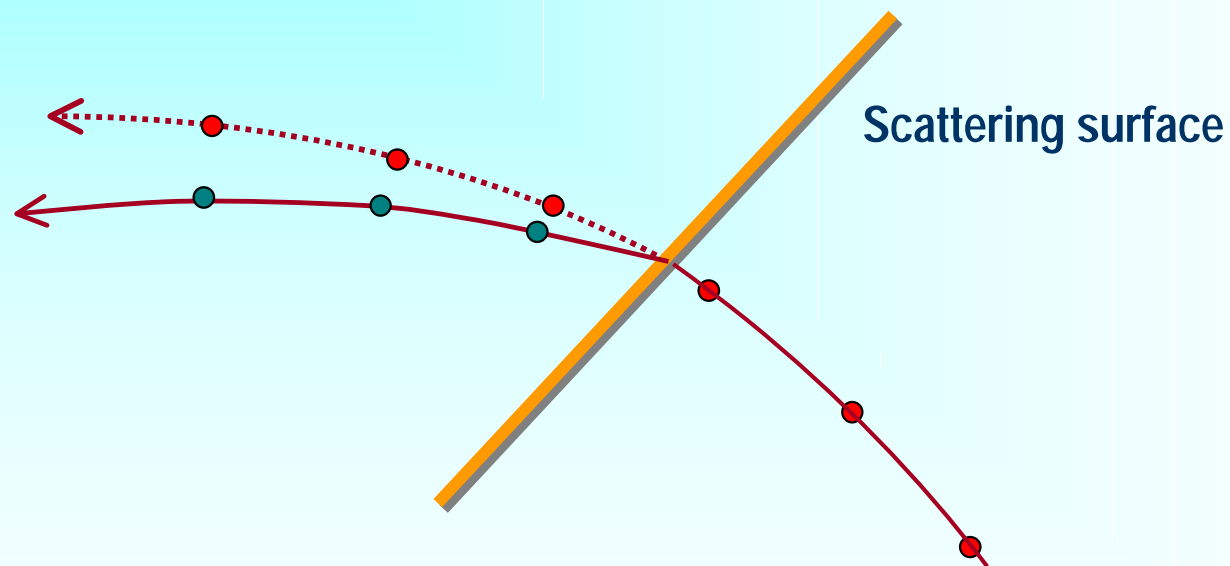
$$V_{\lambda\lambda} \rightarrow V_{\lambda\lambda} + \Theta_0^2 \quad V_{\omega\lambda} \rightarrow V_{\omega\lambda} + \omega \Theta_0^2 \sin\lambda / \cos^3\lambda$$

$$V_{\phi\phi} \rightarrow V_{\phi\phi} + \Theta_0^2 / \cos^2\lambda \quad V_{\omega\omega} \rightarrow V_{\omega\omega} + \omega^2 \tan^2\lambda \Theta_0^2$$



# Correlations

- ◆ Scatter in one point, and all points from then on get shifted, so in this sense they become **correlated**.
- ◆ This is reflected in the non-zero off-diagonal elements of the covariance matrix





# Energy Loss

- ◆ The processes:
  - ionization described by Bethe-Bloch formula:

$$-\frac{dE}{dx} = Kz^2 \frac{Z}{A} \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta}{2} \right]$$

- for  $e^-$ 's, *bremsstrahlung* is significant

- ◆ Basis of particle ID (see DCH talk)



# Energy Loss in Track Fitting

- ◆ Energy loss affects the track parameters (momentum):

$$\omega' \rightarrow \omega p / \sqrt{p^2 + 2E\Delta E + \Delta E^2}$$

- ◆ It also affects the error matrix (process is not deterministic):

$$V_{\omega\omega} \rightarrow V_{\omega\omega} + \omega^2 E^2 \delta E^2 / P^4$$



# Residuals

- ◆ **Residual** for track parameter  $\alpha$ :

$$r = \alpha_{\text{meas}} - \alpha_{\text{track}}$$

where  $\alpha_{\text{track}}$  is the result of the fit

- Fitters can typically provide two types of hit residuals:
  - when the hit is included in the track fit
  - when it's excluded
- ◆ Note that the  $\chi^2$  can be written as

$$\chi^2 \equiv \sum_l \frac{r_l^2}{\sigma_l^2}$$

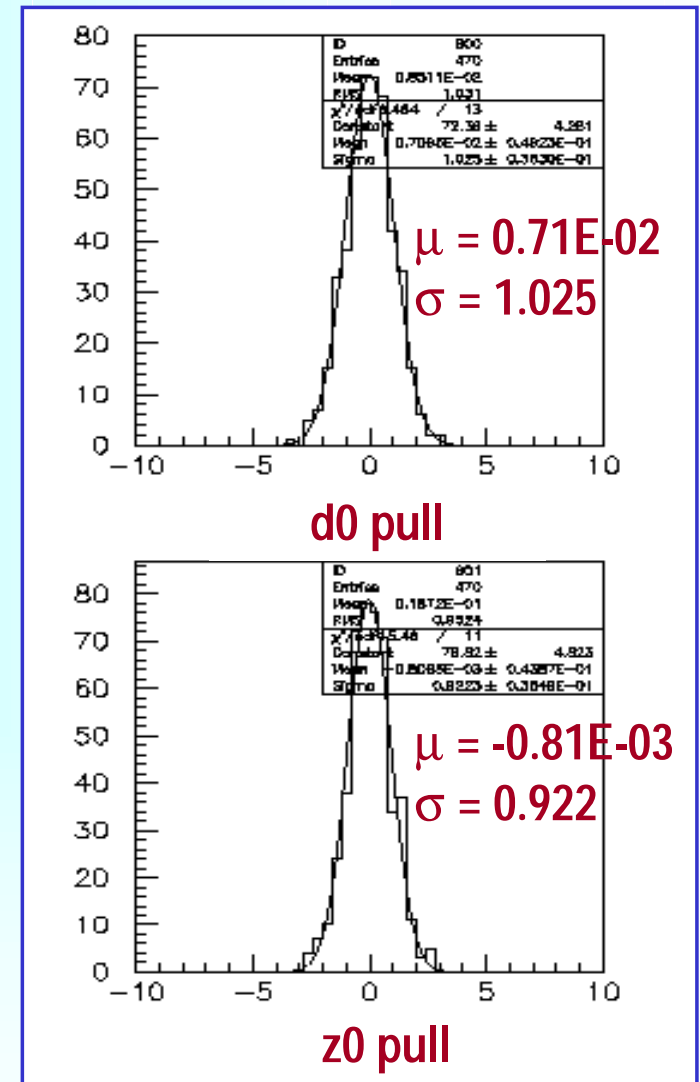


# Pulls

- ◆ A **pull** for a track parameter  $\alpha$  is defined as:

$$pull_{\alpha} \equiv \frac{\alpha_{meas} - \alpha_{track}}{\sigma_{\alpha}}$$

- ◆ If a fit is reasonable and errors are estimated correctly, expect a Gaussian with  $\sigma = 1$  and  $\mu = 0$ 
  - So by plotting pulls can see if errors are correct or over/under estimated!

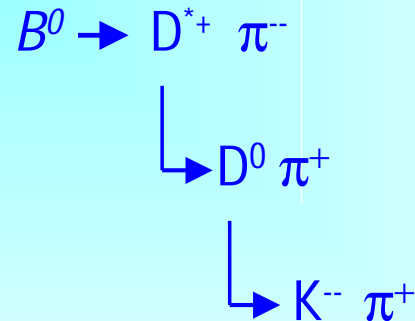




# Kinematic Fitting

- ◆ The idea of kinematic fitting is to use the known properties (**constraints**) of a given physical process to improve the measurements describing the process.

◆ E.g.:



- $D^0$  mass constraint for  $K, \pi$  (1)
- $D^0$  vertex (1)
- $D^{*+}$  mass constraint for  $D^0, \pi$  (1)
- $B$  vertex (3)
- $E_{\text{total}} = \text{beam energy}$  (1)

**7 constraints total**



# Lagrange Multiplier Method

- ◆ The general algorithm for constrained fitting is based on the **Lagrange multiplier** method.
  - The idea is to incorporate the process' constraints into the calculation of the parameters of interest.
  - Simple example: back to back particles



Constraint equations:

$$p_{x1} + p_{x2} = 0$$

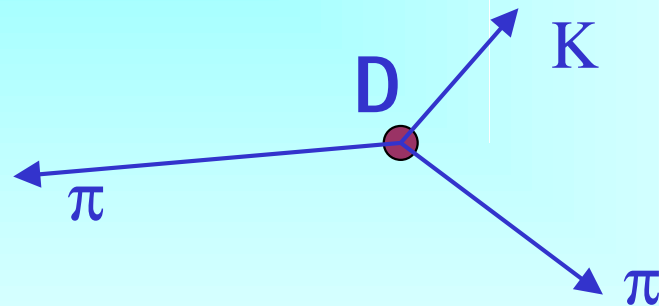
$$p_{y1} + p_{y2} = 0$$

$$p_{z1} + p_{z2} = 0$$



# Vertexing

- ◆ A particular application of constrained fitting is **vertexing** when the tracks are required to come from the same point in space



- we know that  $K, \pi, \pi$  must come from the same vertex
- this fact can be used to improve the mass resolution of the  $D$



# Conclusion

- ◆ The goal of tracking is to determine the **5 helical parameters** of a track as precisely as possible in the presence of complicated physical effects (multiple scattering, energy loss, non-uniform **B**-field, etc.).
- ◆ BaBar tracking devices, the SVT and the DCH, complement each other and allow us to get tracks over a large range of  $p_T$  (from  $\sim 40$  MeV/c to a few GeV/c).