

Summary of $|V_{ub}|$ Determinations

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Introductory Remarks

- * $|V_{ub}|$ can be measured in semileptonic B decays based on $b \rightarrow u l \nu$ transitions
- * **main obstacle** is the large background from $b \rightarrow c l \nu$ decays, whose elimination requires tight experimental cuts that are difficult to account for theoretically
 - \Rightarrow there is no single best way of measuring $|V_{ub}|$
- * study of **exclusive** and **inclusive** B decays allows extractions of $|V_{ub}|$ with very different kinds of theoretical uncertainties

$|V_{ub}|$ from Exclusive Decays

- * theory must provide predictions for form factors with controlled theoretical errors
 - least sophisticated approach
 - $\delta|V_{ub}|/|V_{ub}| = \delta F(q^2)/F(q^2)$ at any value of q^2

Different methods:

- lattice QCD
- light-cone QCD sum rules
- other techniques
- flavor symmetries

Form Factors from Lattice QCD

- * method of choice for a precise determination free of uncontrolled errors
- * recent progress in calculation of $B \rightarrow D^{(*)}$ form factors very encouraging
[→ A. Kronfeld]
- * but heavy-to-light transitions pose new challenges:
 - for $B \rightarrow \pi$ physics near zero recoil is affected by B^* pole contribution (B^*-B mass splitting is off by about a factor 2 in the quenched approximation)
 - for $B \rightarrow \rho$ need to get ρ lineshape right (problem of unstable particles in euclidean space)
 - calculation of chiral corrections (lattice–continuum matching) has its intricacies
- * on a timescale of 5 years, none of these appear as a show stopper

Form Factors from QCD Sum Rules

- * next best theoretical tool are light-cone QCD sum rules, which are specifically designed for describing heavy-to-light transitions near zero recoil
- * in past years, significant refinements have been achieved (NLO corrections, higher-twist effects, target-mass corrections)
 - although not entirely model-independent, sum rules make a large number of testable predictions, which can be confronted with data
 - state-of-the-art predictions can be trusted at level of 10–20%
- * matching between lattice and sum rule calculations in some range of q^2 values provides a powerful consistency check of both techniques

Other Techniques

- * soft-collinear effective theory implies novel relations between heavy-to-light form factors in the heavy-quark limit
 - have been implemented in relativistic quark models, which are much more sophisticated than earlier models such as BSW, KS, ISGW
- * constraints on shapes of heavy-to-light form factors can be derived using analyticity and unitarity (dispersive approach)
- * promising strategy is to **combine** information from lattice, sum rules, heavy-quark relations, and dispersive methods [→ Postler]

Form Factors using Flavor Symmetries

- * idea is to use a combination of heavy-quark and SU(3) symmetries to derive relations between various $B \rightarrow \pi, \rho$ and $D \rightarrow \pi, K, K^*$ form factors at the same value of recoil velocity
- * using charm data, the $B \rightarrow \pi, \rho$ form factors can then be predicted up to symmetry-breaking corrections
- * nice academic exercise, but reliance on Λ/m_c expansion and SU(3) symmetry is a severe drawback
 - both Λ/m_c and SU(3) breaking are expected to be large in the charm sector (don't trust $20\% \times 20\% = 4\%$ arithmetics!)
 - not clear that a “recoil window” exists where the symmetry relations work (limited kinematic range in D decays)
- * precision less than for any decent quark model (most models respect symmetry relations, and so spread between models is an estimator of theoretical errors)

$|V_{ub}|$ from Inclusive Decays

Are we in good shape, or are we in no shape?

$|V_{ub}|$ from Inclusive Decays

Are we in good shape, or are we in no shape?

- * Is it an advantage (really!) to avoid sensitivity to the shape function?
- * answer not clear!
 - trade-off between Fermi motion effects and duality violations and/or larger power corrections
 - **universality** of shape function in B decays into massless partons implies many powerful relations

$|V_{ub}|$ from Inclusive Decays

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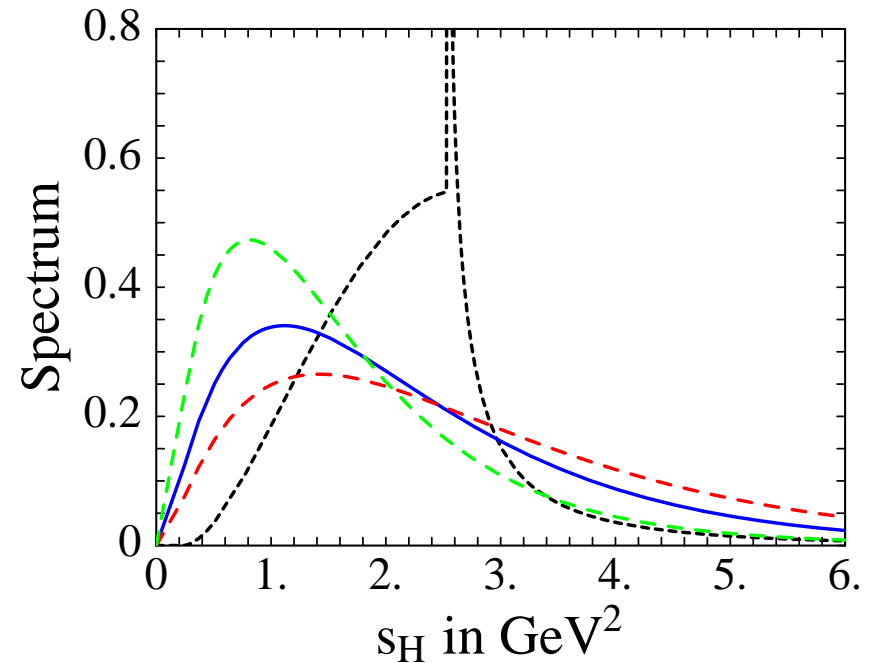
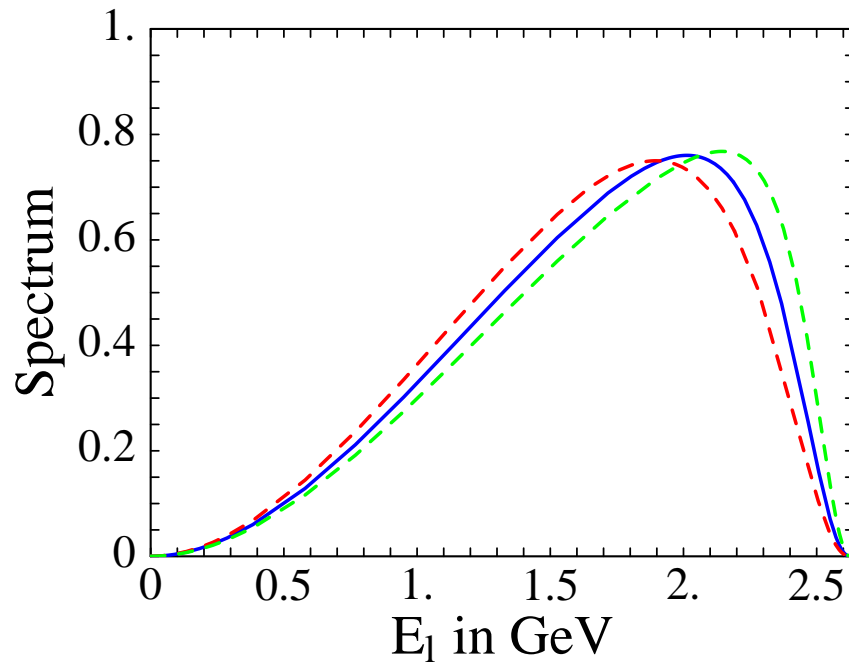
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Would you avoid using parton distribution functions in DIS and collider physics?

Some prominent strategies:

- * cut on charged-lepton energy: $E_l > \frac{m_B^2 - m_D^2}{2m_B}$
⇒ eliminates $\sim 90\%$ of signal events
- * cut on invariant hadronic mass: $M_H < m_D$
⇒ eliminates $\sim 20\%$ of signal events
- * cut on invariant leptonic mass: $M_{l\nu} > m_B - m_D$
⇒ eliminates $\sim 80\%$ of signal events
- * for last strategy, characteristic scale of heavy-quark expansion $\mu_c \approx 1 \text{ GeV}$ is much less than b -quark mass

- * for first two strategies, results are sensitive to **Fermi motion** of the b quark (“QCD smearing” of parton spectra, described by a nonperturbative **shape function**)



$|V_{ub}|$ using Radiative B Decays

- * at leading power in Λ/m_b , Fermi motion effects can be eliminated by taking a ratio of **weighted integrals** over the endpoint regions of the charged-lepton energy spectrum in $B \rightarrow X_u l \nu$ and the photon energy spectrum in $B \rightarrow X_s \gamma$:

[M.N., 1993; \rightarrow Rothstein]

$$\left| \frac{V_{ub}}{V_{cb}} \right|^2 \simeq \left| \frac{V_{ub}}{V_{tb} V_{ts}^*} \right|^2 = \frac{3\alpha}{\pi} K_{\text{pert}} \frac{\hat{\Gamma}_u(E_0)}{\hat{\Gamma}_s(E_0)} + O\left(\frac{\Lambda}{m_B}\right)$$

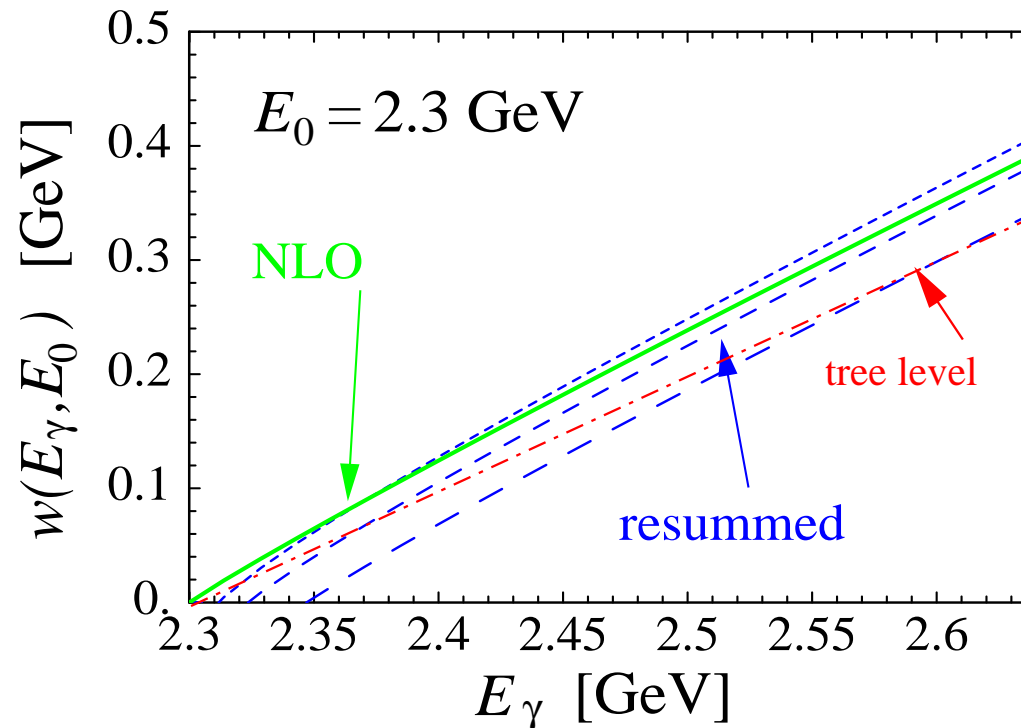
where:

$$\hat{\Gamma}_u(E_0) = \int_{E_0}^{m_B/2} dE_l \frac{d\Gamma(B \rightarrow X_u l \nu)}{dE_l}$$

$$\hat{\Gamma}_s(E_0) = \frac{2}{m_B} \int_{E_0}^{m_B/2} dE_\gamma w(E_\gamma, E_0) \frac{d\Gamma(B \rightarrow X_s \gamma)}{dE_\gamma}$$

- * short-distance correction K_{pert} and weight function $w(E_\gamma, E_0)$ known at NLO
- * results are stable:

$$K_{\text{pert}} = 0.134^{+0.007}_{-0.009} [\text{scale}]^{+0.007}_{-0.006} [m_c] \pm 0.010 [\alpha_s^2]$$



- * theoretically mature, little room for improvements
- * results allow determination of $|V_{ub}|$ with decent theoretical uncertainty:

$$\delta|V_{ub}| = 5\% \text{ [pert.]} \oplus 10\text{--}20\% \text{ } [\Lambda/m_b]$$

Main advantage:

- Fermi motion effects cancel (at leading order)

Main limitations:

- duality violations (strong weighting towards resonance region)
 - Λ/m_b corrections difficult to control
- * importance of these effects can be probed by checking that result for $|V_{ub}|$ is independent of the experimental cutoff E_0 employed

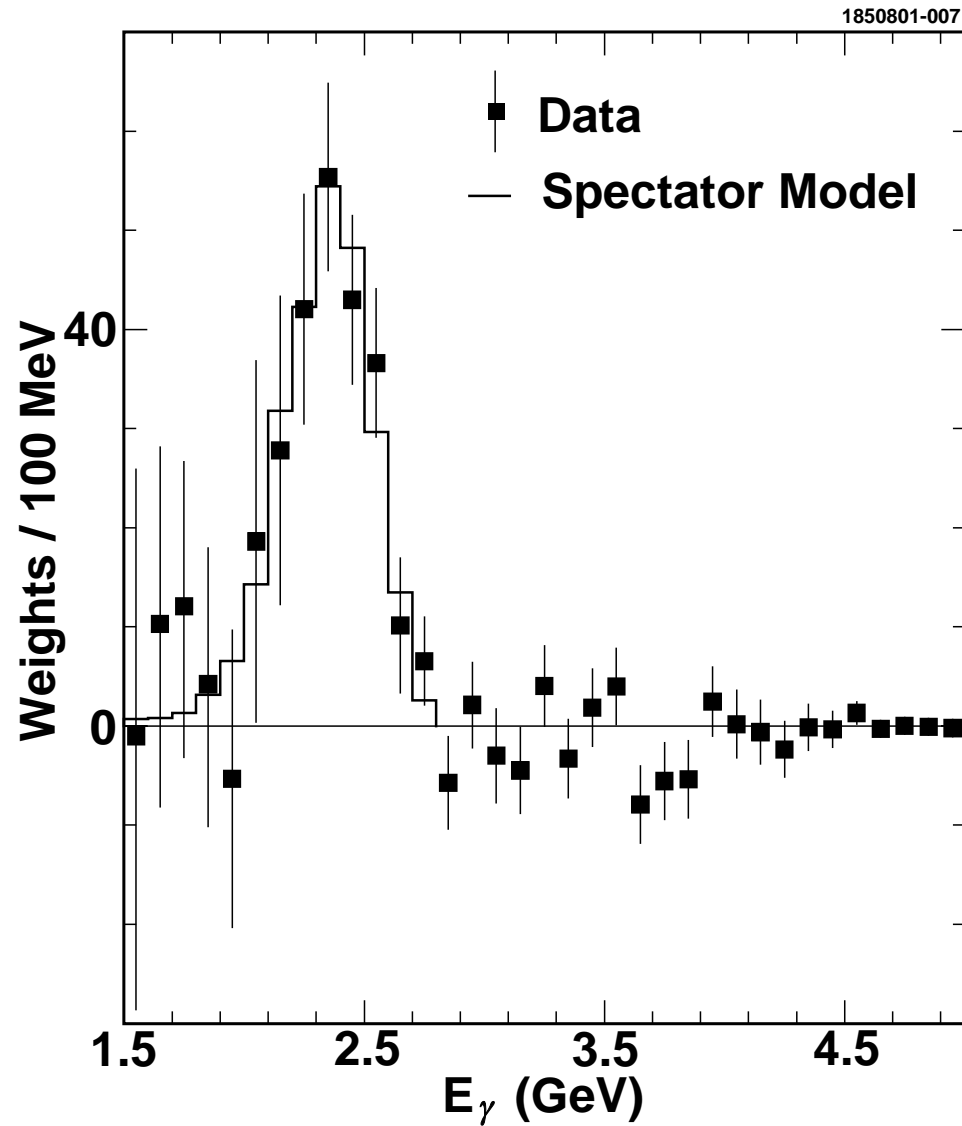
$|V_{ub}|$ via Extraction of the Shape Function

* shape function can be determined with good precision from $B \rightarrow X_s \gamma$ photon spectrum (requires a deconvolution, but experimenters do that all the time)

* facilitated by fact that parton spectrum is, to a good approximation, a spike:
 $d\Gamma/dE_\gamma \propto \delta(E_\gamma - m_b/2)$

(this fact is also used in the global V_{cb} , $\bar{\Lambda}$, λ_1 fit [\rightarrow Becher])

* here is the shape function (almost...) measured by CLEO:



- * once shape function is known, **all** $B \rightarrow X_u l \nu$ decay distributions can be predicted over their **entire** kinematic range
 - allows one to use $E_l, M_H, E_H, M_{l\nu}$ spectra or combinations thereof
 - analog of what is routinely done in DIS and collider physics
 - shape function parameterization can be refined as the data get more precise
 \Rightarrow no intrinsic limitation

Advantages:

- optimal use of data
- least susceptible to duality violations
- first-order Λ/m_b corrections switch themselves off as kinematic regions are enlarged

Limitation:

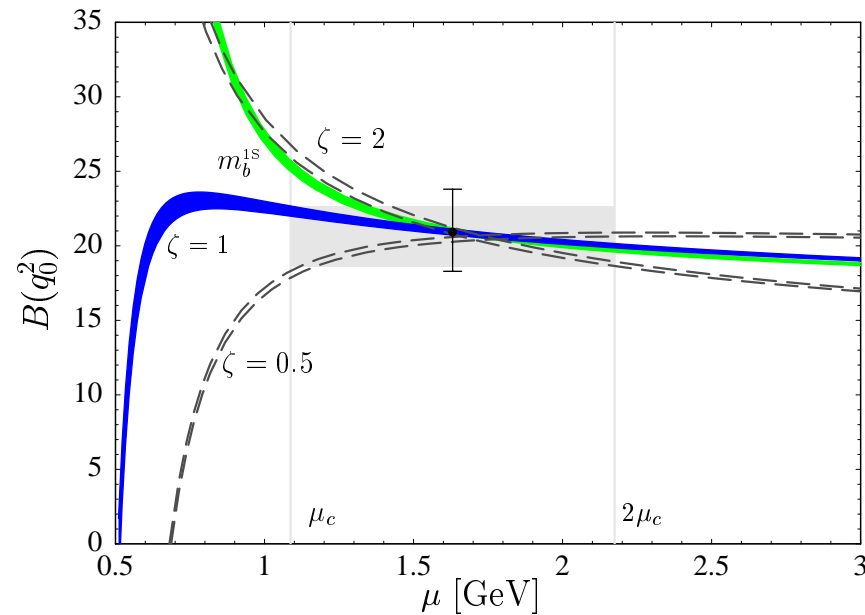
- mostly experimental (quality of $B \rightarrow X_s \gamma$ data)

Avoiding the Shape Function

- * sensitivity to Fermi motion effects can be eliminated by cutting out the “collinear region” $E_H \gg M_H$
 - experimental cut $E_H < m_D$ converts “shape function terms” $\Lambda E_H / M_H^2 \sim 1$ into power corrections $\sim (\Lambda / \mu_c)$
 - formally, this yields a **local** OPE in powers of Λ / μ_c
 - some technical complexities arise from the simultaneous expansion in small ratios μ_c / m_b and $\Lambda_{\text{QCD}} / \mu_c$ (“hybrid expansion”)

* define, with $q_0^2 = (m_B - m_D)^2$:

$$\text{Br}(B \rightarrow X_u l \nu) |_{q^2 > q_0^2} \equiv B(q_0^2) \times |V_{ub}|^2 \left(\frac{\tau_B}{1.6 \text{ ps}} \right)$$



* result:

$$B(q_0^2) = 20.9_{-2.6}^{+2.9} [m_b]_{-2.3}^{+1.8} [\text{pert.}] \pm 2.0 [(\Lambda/\mu_c)^3] = 20.9 \pm 4.0$$

⇒ allows determination of $|V_{ub}|$ with $\sim 10\%$ accuracy

Main advantage:

- formally a **local** OPE, controlled by expansion in powers of Λ/m_c

Main limitations:

- ditto!
- dominant source of error is the **guestimate** of $(\Lambda/m_c)^3$ corrections (known to be huge for charm lifetimes)
- with $\Lambda = 0.5$ GeV these have a 5% effect on $|V_{ub}|$; with $\Lambda = 0.7$ GeV they have a 15% effect; with $\Lambda = 1$ GeV they have an 40% effect
- will be difficult to reduce or eliminate this uncertainty

Main advantage:

- formally a **local** OPE, controlled by expansion in powers of Λ/m_c

Main limitations:

- ditto!
 - dominant source of error is the **guestimate** of $(\Lambda/m_c)^3$ corrections (known to be huge for charm lifetimes)
 - with $\Lambda = 0.5$ GeV these have a 10% effect on $|V_{ub}|$; with $\Lambda = 0.7$ GeV they have a 30% effect; with $\Lambda = 1$ GeV they have an 80% effect
 - hard to see how this uncertainty could be reduced or eliminated
- * to paraphrase Zoltan Ligeti (HEP 2001):

Unfortunately, it is much easier to write down a Λ/m_c expansion than to test its convergence!

Combination of Different Cuts

- * combinations such as $M_{l\nu}$ cut \oplus M_H cut appear very promising [\rightarrow Bauer, Menke]
 - **Caveat:** need to include $O(\alpha_s)$ corrections to $d\Gamma/dM_H$ spectrum for a realistic study of shape function effects
- * typically, a trade-off between Fermi motion effects and duality violations and/or large Λ/m_c corrections
- * optimization should take into account both theoretical and experimental realities
- * ultimately, such combinations are the best bet to reducing the error on $|V_{ub}|$ below the 10% level

Summary

- * clean determinations of $|V_{ub}|$ are difficult and hindered by the necessity to eliminate the large charm background
- * variety of proposals exist, each with their own theoretical limitations
- * combination of various methods will help to reduce the error on $|V_{ub}|$ below the 10% level and get confidence in the reliability of the theoretical methods