
The Future of Lattice Calculations

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Outline

- Brief discussion of systematics on lattice
- Issues with staggered fermions
- What quantities can be calculated well on the lattice?
- State of the art (focus on f_π , f_K — model for heavy-light calculations we want to do)
- One more important systematic error
- What precision can be expected, and when?

Systematic Errors

Chiral extrapolation:

- Computer time rises like large power of $1/m_{u,d}$ as these masses approach their physical values.
- Must work with larger masses and extrapolate; use chiral perturbation theory (χ PT) to guide extrapolation.
 - “Partially quenched” χ PT ($m_{\text{val}} \neq m_{\text{sea}}$) esp. useful.
 - Note: as long as there are 3 flavors of sea quarks, partially quenched world has real world as a subset.
- Essentially all interesting quantities with light (u, d) quarks have significant curvature $\sim m_\pi^2 \ln(m_\pi^2/\Lambda_\chi^2) \Rightarrow$ must get to small quark mass (probably $m_{u,d} \sim m_s/4$ to $m_s/8$) to control chiral logs in extrapolation.
- If only have large mass ($m_{u,d} \gtrsim m_s/2$), will be stuck with 10% or even 20% errors (**Kronfeld & Ryan**).

Systematic Errors

Quenching:

- Most lattice calculations to date have left out the effects of sea quarks.
- Reduces computer time by large factor: ~ 100 – 1000 or more.
- But an uncontrolled approx: errors 10%–20% or more.
- Need to move beyond quenching for lattice to fulfill promise & have an impact on phenomenology.
- Dynamical (sea quark) masses must be small: otherwise just trading quenching error for chiral extrap error.
- In near term only “staggered” quarks are viable:
 - very fast
 - have residual chiral symmetry that protects from “exceptional configurations”
⇒ can go to low quark mass

Systematic Errors

Discretization Effects (finite lattice spacing a):

- Order of errors in a depends on fermion action:
 - Wilson: $\mathcal{O}(a)$
 - Improved Wilson (“clover”) (perturbative): $\alpha_S^2 \mathcal{O}(a)$
 - Improved Wilson (“clover”) (nonperturbative): $\mathcal{O}(a^2)$
 - Staggered: $\mathcal{O}(a^2)$
 - Improved staggered (“Asqtad”): $\mathcal{O}(\alpha_S a^2)$
 - Domain wall: $\mathcal{O}(a^2)$
 - Overlap: $\mathcal{O}(a^2)$
- Now & foreseeable future: if few % accuracy is desired, discretization errors are not negligible at fixed a for any action.
- **Must** extrapolate to continuum ($a \rightarrow 0$) using known leading a dependence.

Systematic Errors

Finite Volume:

- Depends on how many hadrons are in initial or final state:
 - More than 1 hadron \Rightarrow probably out of reach in next five years for all but the most qualitative studies.
 - Single particles \Rightarrow currently feasible volumes are enough to reduce finite volume effects to few % level (without major sacrifice on discretization errors).
 - Typically $V \gtrsim (2.5\text{fm})^2$ will do the job.

Systematic Errors

Finite Volume:

- Can do even better for single-particle quantities whose mass dependence is determined by χ PT.
- χ PT predicts the volume dependence (for large volume), so can correct for finite volume effects \Rightarrow negligible errors ($\sim 0.1\%$).
- Ignore finite volume effects for single-particle states from here on.

Systematic Errors

“Setting the Scale”

- Lattice spacing a is determined after the fact by comparing the result for one dimensional quantity with experiment. (Equivalent to fixing Λ_{QCD} or α_S .)
- So lattice error in the quantity used to set the scale will infect all other dimensionful results.
- The best we can do today is probably from Υ splittings: 2S-1S or 1P-1S: gives about 2% scale error.
- The scale error is usually negligible on dimensionless ratios (like f_{B_s}/f_B), but is not strictly 0. (Error can enter indirectly through determination of quark masses or momenta.)

Staggered Baggage

Staggered quarks have an incomplete reduction of lattice doubling symmetry

- Each staggered lattice field (each staggered flavor) \Rightarrow 4 equivalent particles in continuum.
- The quantum number of this unphysical multiplication of d.o.f. is called “taste.”
- Taste symmetry, while presumably exact in the continuum, is violated on the lattice at $\mathcal{O}(\alpha^2 a^2)$ (for improved staggered).
- Although formally non-leading (leading discretization errors are $\mathcal{O}(\alpha a^2)$), taste violations are numerically important.
- Leads to both practical and theoretical issues.

Staggered Baggage

- Practical issue: taste violations must be taken into account in χ PT used for extrapolation in light quark mass.
 - Need to fit finite- a lattice data to “staggered chiral perturbation theory” ($S\chi$ PT).
 - $S\chi$ PT has been worked out (Lee & Sharpe; Aubin & C.B.) for π - K system.
 - Applying $S\chi$ PT to heavy-light mesons should be straightforward, but has **not** yet been done.
 - Will present future error estimates with & without assumption that heavy-light $S\chi$ PT will have been worked out.

Staggered Baggage

- Theoretical issue: need to eliminate taste degree of freedom for quarks in simulations (need 1 flavor – 1 quark).
 - For valence quarks, taste is simply an annoyance: just choose external quarks to have a single taste.
 - But for dynamical quarks, must reduce tastes to 1 per flavor by taking $\sqrt[4]{\text{Det}}$
 - At finite lattice spacing this is nonlocal operation. Could it introduce nonuniversal behavior; nonlocality in the continuum?
 - $\sqrt[4]{\text{Det}}$ is correct at every finite order of perturbation theory.
 - Also some significant nonperturbative evidence against this disaster (and no evidence for it), but many reasonable people remain unconvinced.

Staggered Baggage

- Estimates below on what can be done on lattice **assume** that staggered quarks are ok: give QCD in the continuum limit.
- If not, safer but slower methods can be used, but they will lead to **much** larger systematic errors for the quantities you want to know.

Gold-plated Quantities

So what quantities can the lattice compute well (few %) in the next ~ 5 years?

- At most one hadron in initial and final state
 - large enough volume, small enough a are then feasible for unquenched theory.
- Stable hadrons, not near thresholds
 - No ρ or K^* , sorry!
 - Unstable particles require very large volumes (and untested techniques) to treat decay products correctly
 - Decays like $B \rightarrow \rho$ **will** be calculated, but errors will **not** be very well controlled ($\sim 10\%$?) — model dependence in answers and error estimates.
 - $B \rightarrow D^* \ell \nu$ may be an exception because model dependence multiplies $\mathcal{F}(1) - 1$.

Gold-plated Quantities

So what quantities can the lattice compute well (few %) in the next ~ 5 years?

- The η also probably excluded. [Need to include η - η' mixing \Rightarrow disconnected graphs \Rightarrow difficult]
- No high momenta (need $|\vec{p}|a \ll 1$).
 - probably limited to $|\vec{p}| \lesssim 1\text{GeV}$
 - $\Rightarrow q^2 \gtrsim (17\text{ GeV})^2$ for $B \rightarrow \pi$ semileptonic form factors.
 - prob. also need $q < q_{\text{max}}$ to avoid B^* pole at end point. (Minimum lattice momentum for fixed lattice size may then require $|\vec{p}| \gtrsim 350\text{ MeV}$ or more.)
- Chiral extrapolation must be under good control
- We call such quantities “gold-plated” = “no excuses”

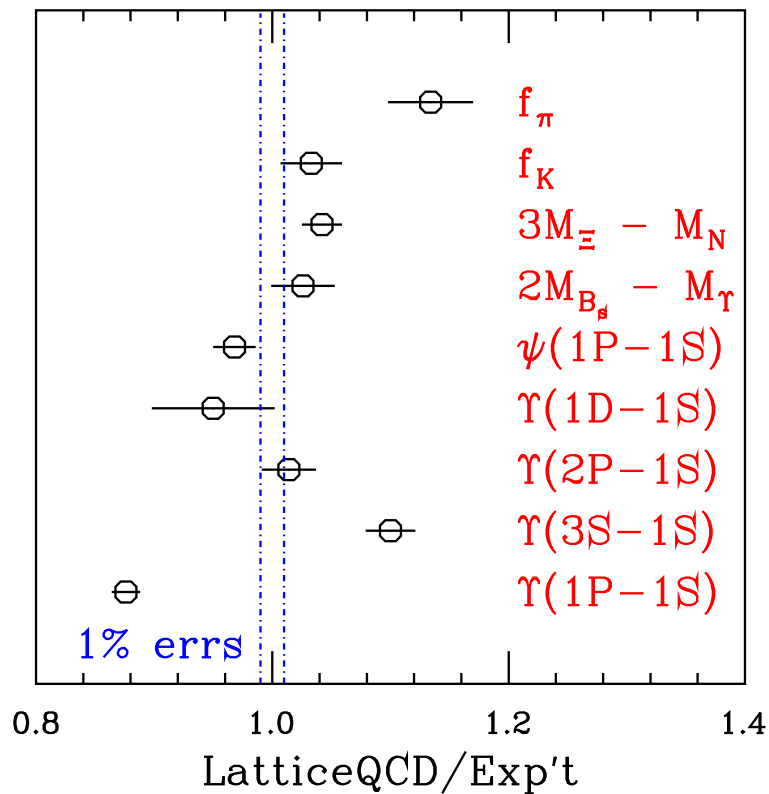
State of the Art

MILC configurations

- Since 1999, MILC collaboration has been generating lattice configurations including three sea-quark flavors (u, d, s), using improved staggered quarks (with $m_u = m_d$).
- Large sample at $a \approx 0.125$ fm and $L \approx 2.5$ fm.
 - Wide range of sea quark masses
 - Lowest $m_{u,d} \sim 10$ MeV (“ m_π ” ≈ 260 MeV).
- Smaller sample at $a \approx 0.09$ fm, $L \approx 2.5$ fm still being generated.
 - Two low-mass sets completed
 - Lowest $m_{u,d} \sim 15$ MeV (“ m_π ” ≈ 320 MeV).
 - Working on $m_{u,d} \sim 10$ MeV (“ m_π ” ≈ 260 MeV).
- Configurations publicly available:
<http://qcd.nersc.gov/>

State of the Art

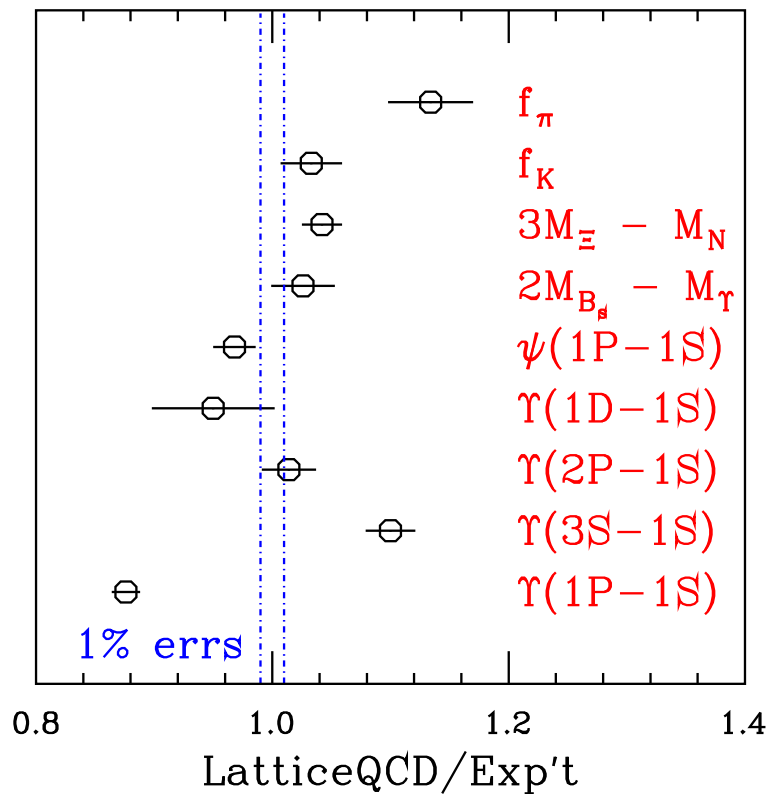
Results with existing MILC configurations (“MILC0 set”
(C. Davies, *et al.*, hep-lat/0304004)



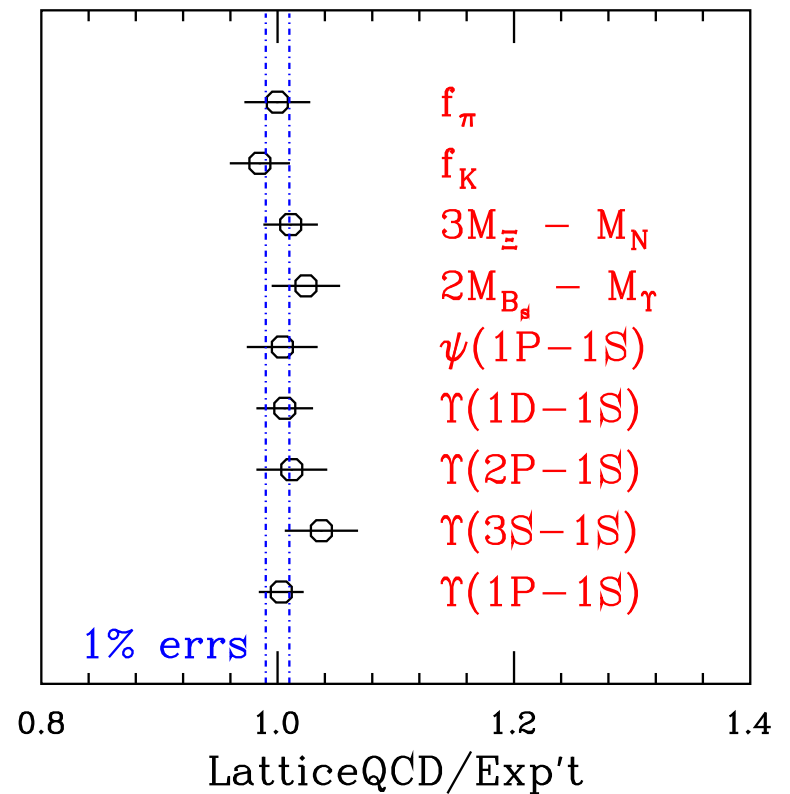
Quenched Approximation

State of the Art

Results with existing MILC configurations (“MILC0 set”
(C. Davies, *et al.*, hep-lat/0304004)



Quenched Approximation



Full QCD ($n_F = 3$)

[scale from Υ 2S-1S; $m_{u,d}, m_s, m_c, m_b$ from $m_\pi, m_K, m_{D_s}, m_\Upsilon$]

π - K System

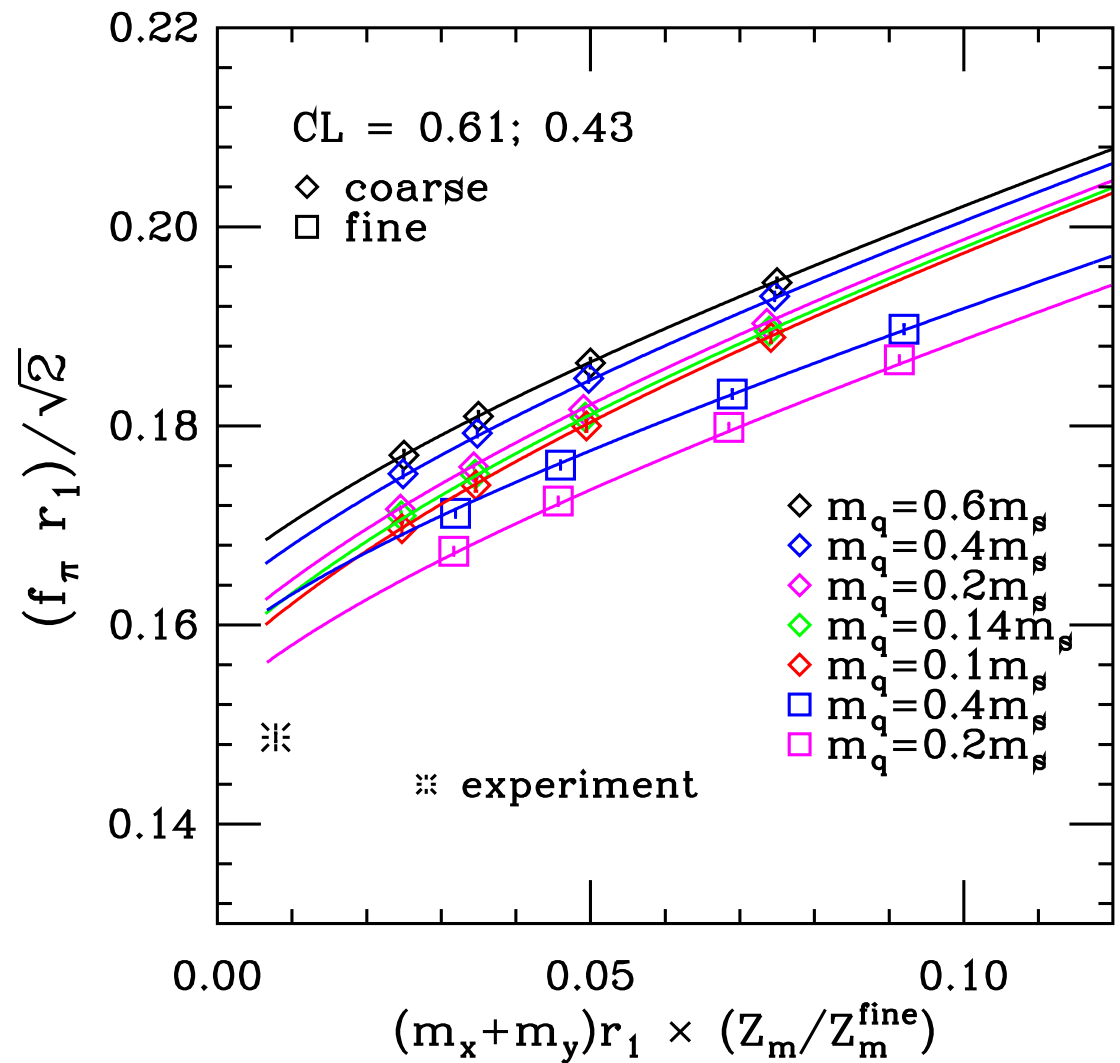
*Current MILC calculation
of f_π and f_K*

- Closest existing example to future calculations of heavy-light leptonic and semileptonic decay constants.
- Basis of the error estimates to follow.

π - K System

Current MILC calculation
of f_π and f_K

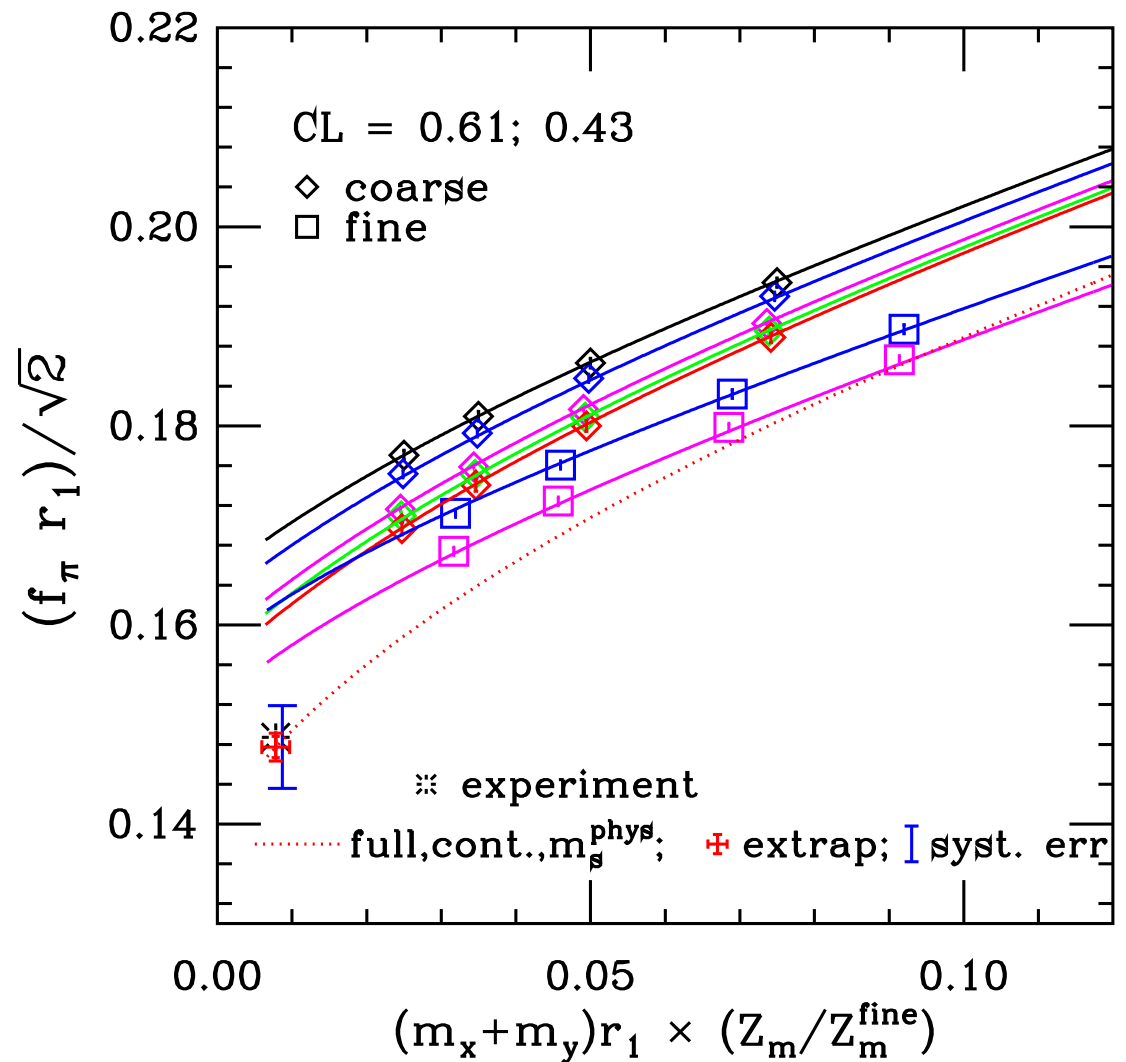
- Fit partially quenched data to $PQ_\chi PT$.



π - K System

Current MILC calculation of f_π and f_K

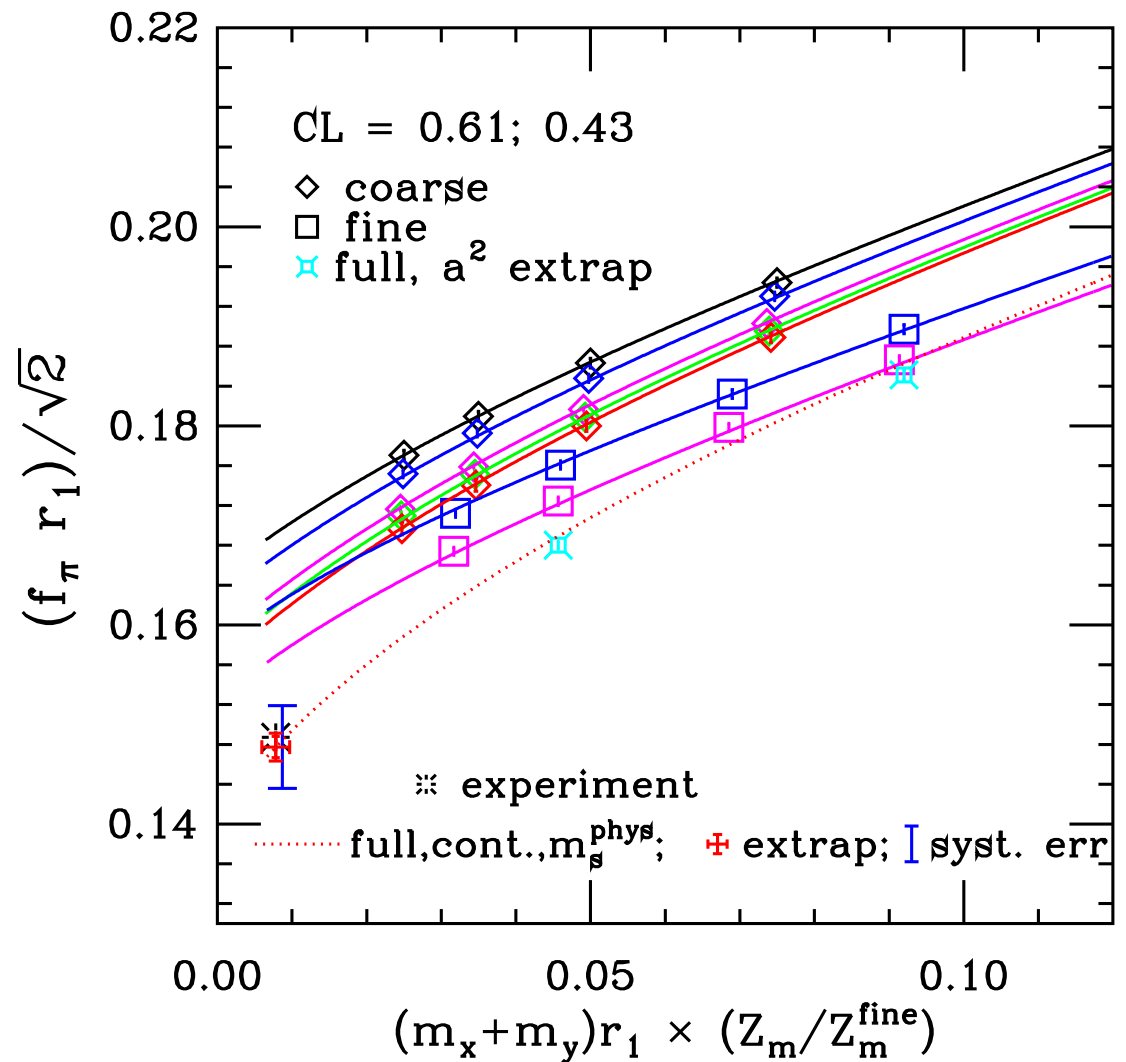
- Extrapolate fit params to continuum
- Set $m_{sea} = m_{val}$ and plot a function of m_{val} (.....)
- Set $m_{val} = m_{u,d}^{phys}$ (determined from m_π).



π - K System

Current MILC calculation of f_π and f_K

- Consistency check: extrapolate points with $m_{val} = m_{sea}$ to continuum at fixed quark mass.



Systematic Errors (*cont. . .*)

Perturbation theory:

- For f_π , f_K , we are lucky: staggered lattice PCAC \Rightarrow light-light axial current not renormalized \Rightarrow lattice & continuum currents are the same.
- But many interesting gold-plated quantities (*e.g.*, heavy-light leptonic & semileptonic decay constants) require perturbative calculation (or equivalent nonperturbative lattice computation) to match lattice current to continuum current.
- In new era of unquenched, light mass simulations, one-loop calculations have unacceptably large errors ($\sim 10\%$).
- But lattice perturbation theory very messy (complicated actions; no Lorentz invariance).

Systematic Errors (*cont. . .*)

Perturbation theory:

- Probably need “automated perturbation theory,” (Trottier, Lepage, & collabs.), based on work by Lüscher & Weisz.
 - My reading: no difficulties in principle with going to 2 or even 3 loops.
 - However, still some practical problems that need to be overcome:
 - Issues of IR regulation important
 - currently use “twisted boundary conditions” on lattice
 - to match to continuum, need to use same twisted boundary conditions there.
 - continuum perturbation theory (dim reg?) with twisted b.c. hard

Systematic Errors (*cont. . .*)

Perturbation theory:

- Will present future error estimates with & without assumption that 2-loop perturbative matching will exist.
- Luckily, many interesting quantities, *e.g.*, ratios like f_{B_s}/f_B , are independent or nearly independent of perturbation theory.

Future Lattice Precision

- Fermilab/MILC collaboration now beginning computations of heavy-light semi-leptonic form factors and leptonic decay constants:
 - Dynamical (sea quark) configurations are the existing improved staggered $n_F = 3$ MILC ones (“MILCO”).
 - Valence light quarks will also be staggered, so chiral regime reachable.
 - CPU requirements are relatively small since dynamical configurations already exist.
 - With time for analysis, I estimate this calculation can be completed in 1-2 years.
 - Estimates for future precision will be much more reliable once this 0th generation project is completed.

Future Lattice Precision

- Next steps will require new dynamical configurations (lighter quark mass, smaller lattice spacing):
 - MILC1
 - ~ 6 Teraflop-years
 - Machines now being built: Columbia QCDOC; large clusters (SciDAC project)
 - ETA: 3-5 years (including analysis time)
 - MILC2
 - ~ 50 -100 Teraflop-years
 - Next generation of machines
 - ETA: 5-8 years (including analysis time)

Future Lattice Precision

- Next steps will require new dynamical configurations (lighter quark mass, smaller lattice spacing):
 - **DWF1**: Dynamical domain wall fermions (or equivalent) at comparable mass and spacings to **MILC1**
 - Comparable precision to **MILC2** because no taste violations
 - Safe from all staggered baggage
 - ~ 1000 Teraflop-years
 - “Next next” generation of machines
 - ETA: 10 years (including analysis time)

Future Lattice Precision

- Error estimates to follow are based on DOE planning document prepared by S. Sharpe, C.B., A. El-Khadra, P. Mackenzie, and R. Sugar.
- However, I have tweaked those estimates using my own expectations/prejudices.

MILCO error estimates

| quantity | statist. | scale | light q | | heavy Q | pert. th. | |
|------------------------------|----------|-------|----------------|-------------|-----------|-----------|--------|
| | | | no S_χ PT | S_χ PT | | 1-loop | 2-loop |
| f_B | 3 | 2 | 5 | 2.5 | 3 | 7.5 | 2 |
| $f_B \sqrt{B_B}$ | 4 | 2 | 5.5 | 3 | 3 | 8.5 | 2.5 |
| f_{B_s}/f_B | 1 | – | 5 | 2.5 | 1 | – | – |
| ξ | 2 | – | 5.5 | 3 | 1 | – | – |
| $B \rightarrow \pi \ell \nu$ | 4.5 | 1 | 6 | 3 | 3 | 7.5 | 2 |
| $B \rightarrow D \ell \nu$ | 1 | 0.5 | 2 | 1 | 1 | 2.5 | 0.7 |

Estimated per cent errors. “Light q ” includes light quark chiral and discretization errors. “Heavy Q ” means heavy quark discretization errors. Semileptonic form factors are for restricted range $0.5 \text{ GeV} \lesssim \vec{p}_\pi \lesssim 1 \text{ GeV}$ (in B rest frame), but can have any bilinear current. $B \rightarrow D^* \ell \nu$ **may** be comparable to $B \rightarrow D \ell \nu$.

MILC1 error estimates

| quantity | statist. | scale | light q | | heavy Q | pert. th. | |
|------------------------------|----------|-------|----------------|-------------|-----------|-----------|--------|
| | | | no S_χ PT | S_χ PT | | 1-loop | 2-loop |
| f_B | 2 | 1.5 | 3 | 1.5 | 2 | 7.5 | 2 |
| $f_B \sqrt{B_B}$ | 3 | 2 | 3.5 | 2 | 2 | 8.5 | 2.5 |
| f_{B_s}/f_B | 0.8 | – | 3 | 1.5 | 0.8 | – | – |
| ξ | 2 | – | 3.5 | 2 | 1 | – | – |
| $B \rightarrow \pi \ell \nu$ | 3 | 0.7 | 4 | 2 | 2 | 7.5 | 2 |
| $B \rightarrow D \ell \nu$ | 0.6 | 0.5 | 2 | 1 | 0.6 | 2.5 | 0.7 |

Estimated per cent errors. “Light q ” includes light quark chiral and discretization errors. “Heavy Q ” means heavy quark discretization errors. Semileptonic form factors are for restricted range $0.35 \text{ GeV} \lesssim \vec{p}_\pi \lesssim 1 \text{ GeV}$ (in B rest frame), but can have any bilinear current. $B \rightarrow D^* \ell \nu$ **may** be comparable to $B \rightarrow D \ell \nu$.

MILC2 error estimates

| quantity | statist. | scale | light q | | heavy Q | pert. th. | |
|------------------------------|----------|-------|------------------|---------------|-----------|-----------|--------|
| | | | no $S_{\chi PT}$ | $S_{\chi PT}$ | | 1-loop | 2-loop |
| f_B | 1 | 1 | 2 | 1 | 1.5 | 7.5 | 2 |
| $f_B \sqrt{B_B}$ | 1.3 | 2 | 2.5 | 1 | 1.6 | 8.5 | 2.5 |
| f_{B_s}/f_B | 0.5 | – | 2.5 | 1 | 0.5 | – | – |
| ξ | 1 | – | 3 | 1 | 0.6 | – | – |
| $B \rightarrow \pi \ell \nu$ | 1.5 | 0.5 | 2.7 | 1.3 | 1.5 | 7.5 | 2 |
| $B \rightarrow D \ell \nu$ | 0.3 | 0.3 | 1.4 | 0.7 | 0.5 | 2.5 | 0.7 |

Estimated per cent errors. “Light q ” includes light quark chiral and discretization errors. “Heavy Q ” means heavy quark discretization errors. Semileptonic form factors are for restricted range $0.35 \text{ GeV} \lesssim \vec{p}_\pi \lesssim 1 \text{ GeV}$ (in B rest frame), but can have any bilinear current. $B \rightarrow D^* \ell \nu$ **may** be comparable to $B \rightarrow D \ell \nu$.

Summary error estimates (%)

| quantity | now | 1-2 yrs. MILC0 | 3-5 yrs. MILC1 | 5-8 yrs. MILC2 |
|------------------------------|-------|-------------------|-------------------|-------------------|
| f_B | 15 | 10, 9, 7, 6 | 9, 8, 5, 4 | 8, 8, 4, 3 |
| $f_B \sqrt{B_B}$ | 15-20 | 12, 11, 8, 7 | 10, 10, 6, 5 | 9, 9, 5, 4 |
| f_{B_s}/f_B | 6 | 5, 3 | 3, 2 | 3, 1 |
| ξ | 7 | 6, 4 | 4, 3 | 3, 1.5 |
| $B \rightarrow \pi \ell \nu$ | 15 | 11, 10, 8, 7 | 9, 9, 6, 5 | 8, 8, 4, 3 |
| $B \rightarrow D \ell \nu$ | 6 | 4, 3, 3, 2 | 3, 3, 2, 1.6 | 3, 3, 2, 1.2 |

Key: **red**= NO $S_{\chi PT}$; 1-loop pert. th.

purple= $S_{\chi PT}$; 1-loop pert. th.

navy= NO $S_{\chi PT}$; 2-loop pert. th. (or pert. th. not needed)

blue= $S_{\chi PT}$; 2-loop pert. th. (or pert. th. not needed)

Summary error estimates (%)

| quantity | now | 1-2 yrs. MILC0 | 3-5 yrs. MILC1 | 5-8 yrs. MILC2 |
|------------------------------|-------|-------------------|-------------------|-------------------|
| f_B | 15 | 10, 9, 7, 6 | 9, 8, 5, 4 | 8, 8, 4, 3 |
| $f_B \sqrt{B_B}$ | 15-20 | 12, 11, 8, 7 | 10, 10, 6, 5 | 9, 9, 5, 4 |
| f_{B_s}/f_B | 6 | 5, 3 | 3, 2 | 3, 1 |
| ξ | 7 | 6, 4 | 4, 3 | 3, 1.5 |
| $B \rightarrow \pi \ell \nu$ | 15 | 11, 10, 8, 7 | 9, 9, 6, 5 | 8, 8, 4, 3 |
| $B \rightarrow D \ell \nu$ | 6 | 4, 3, 3, 2 | 3, 3, 2, 1.6 | 3, 3, 2, 1.2 |

Key: **red**= NO $S_{\chi PT}$; 1-loop pert. th.

purple= $S_{\chi PT}$; 1-loop pert. th.

navy= NO $S_{\chi PT}$; 2-loop pert. th. (or pert. th. not needed)

blue= $S_{\chi PT}$; 2-loop pert. th. (or pert. th. not needed)

= my best guess of what will be available

Hope for the future

- 5 or more years from now, it **may** become feasible to do “no-excuses” calculations on unstable particles
- $B \rightarrow \rho \ell \nu$, $B \rightarrow K^* \gamma$ **may** have errors comparable to $B \rightarrow \pi \ell \nu$ on **MILCO**

Disclaimers

- Errors on predicted errors are at least 30% (and that's “ 1σ ”).
- Errors on predicted errors increase as we move further into the future.